# Reactor to Rheology: Coupled models for polymers

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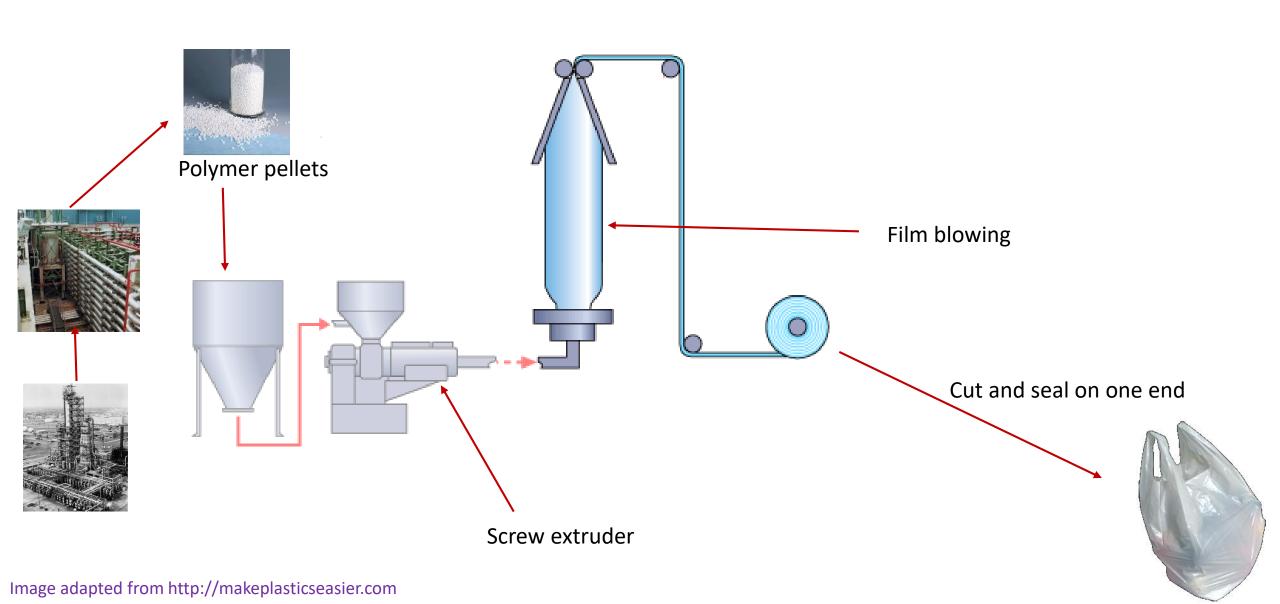
**Dietmar Auhl** Petra Bacova **Taihyun Chang** Jaap den Doelder **John Embery Christine Fernyhough David Hoyle Lian Hutchings Nat Inkson Michael Kapnistos** Mark Kelmanson **Helen Lentzakis Angel Moreno Pradeep Shirodkar Johannes Soulages lakovos Vittorias Dimitris Vlassopoulos** and others...

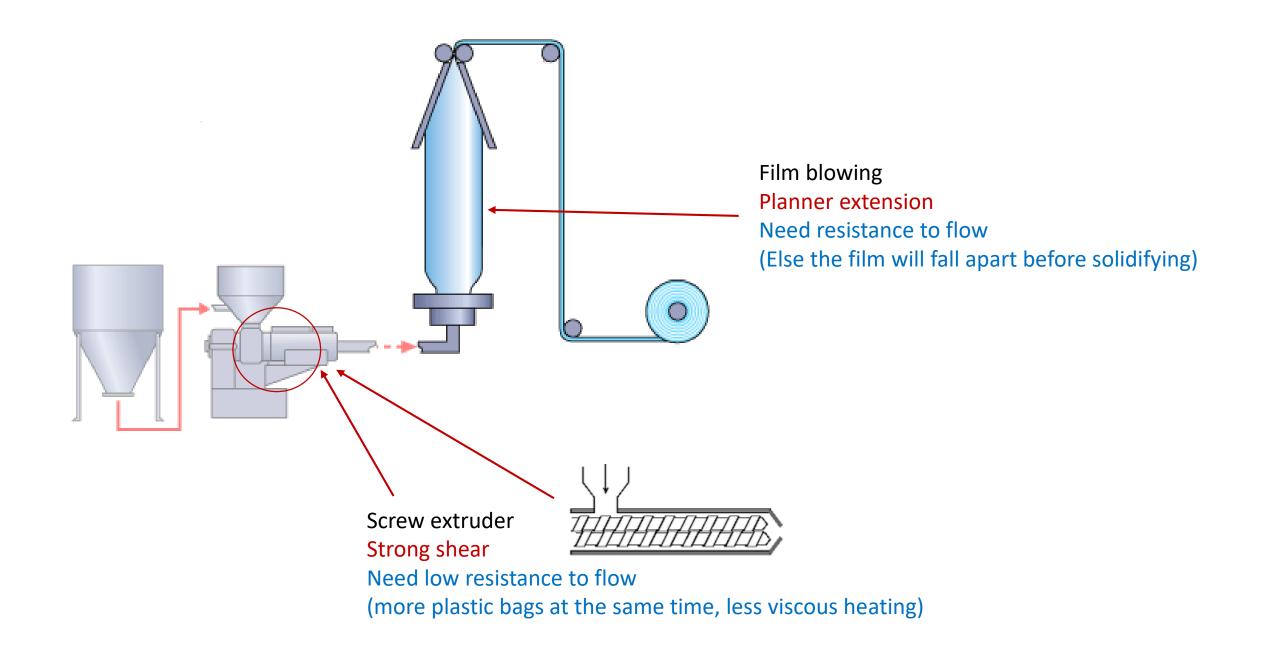


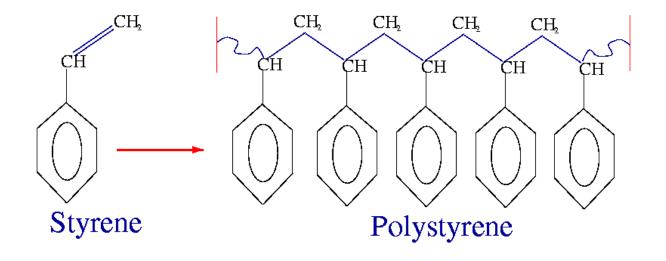


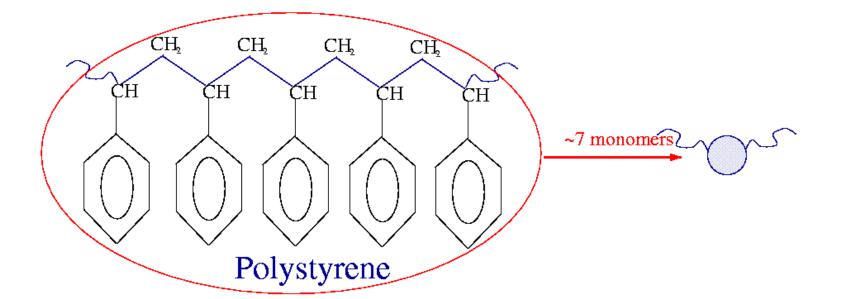
## Rough outline

- > Introduction to polymer molecules and their dynamics in the molten state
- Virtual molecules in imaginary flow
- $\triangleright$  Gelation in  $\alpha$ -olefin/diene copolymers
- > Outlook



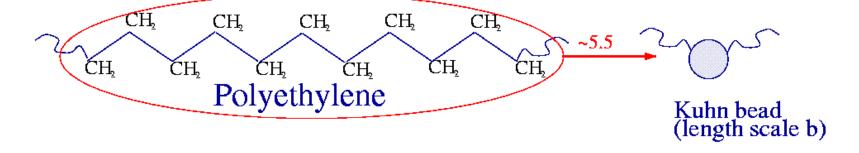






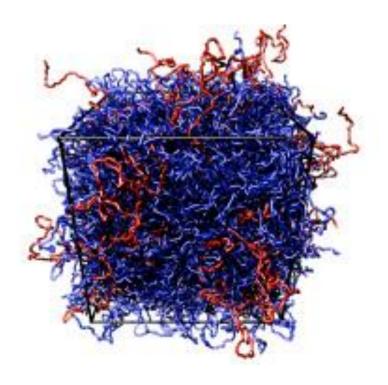


Random walk conformation in melt Entropy dominated



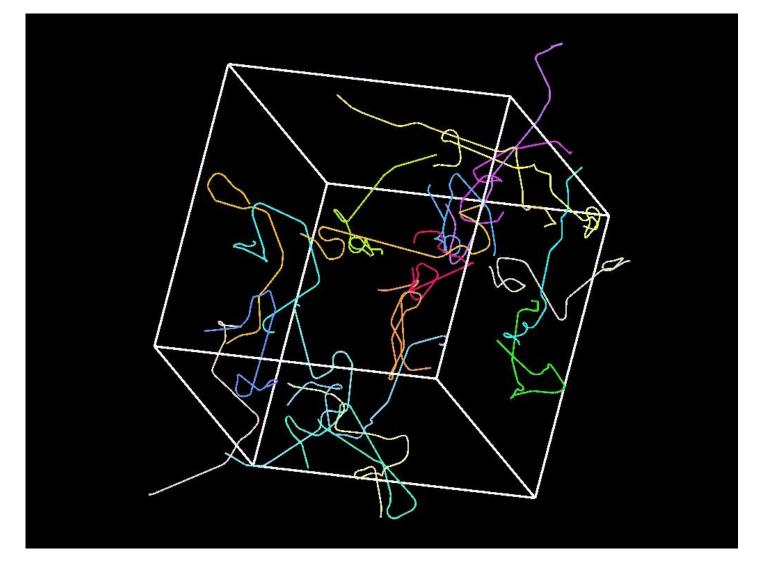
**Energy scale k**<sub>B</sub>**T** 

## Tube theory



Long time and length scales: Chemical details are hidden in a small number of parameters

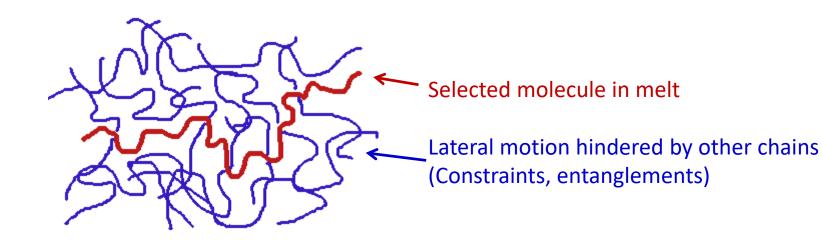
Bead-spring model (soft potential, but stiff enough to ensure that chains don't cross each other).

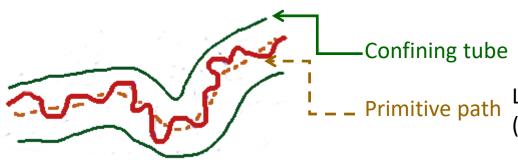


Unfold molecules from periodic box

Long linear molecules:

In certain timescale, the chains are approximately confined in a tube like region.





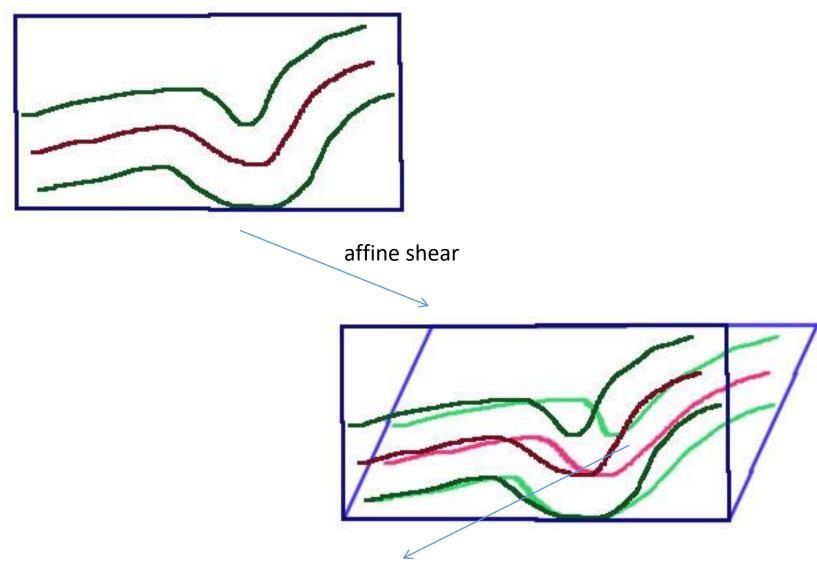
Length scale: **tube diameter, a** (persistence length of primitive path, contains chain of one entanglement)

Entanglement molar mass, M<sub>e</sub>

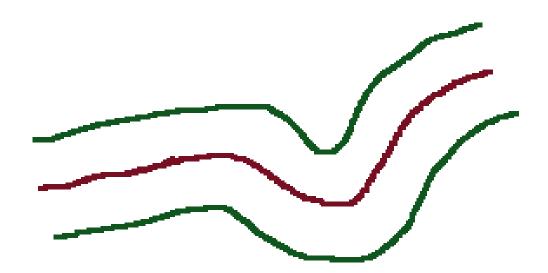
Polyethylene: ~7 Kuhn beads Polystyrene: ~24 Kuhn beads

Time scale: Entanglement time,  $\tau_e$  (Rouse Relaxation time of one entanglement; Time scale at which chains become aware of the tube)

# Stress decay



Deformation changes conformation
Chains can revert back to equilibrium conformation by going out of old tube



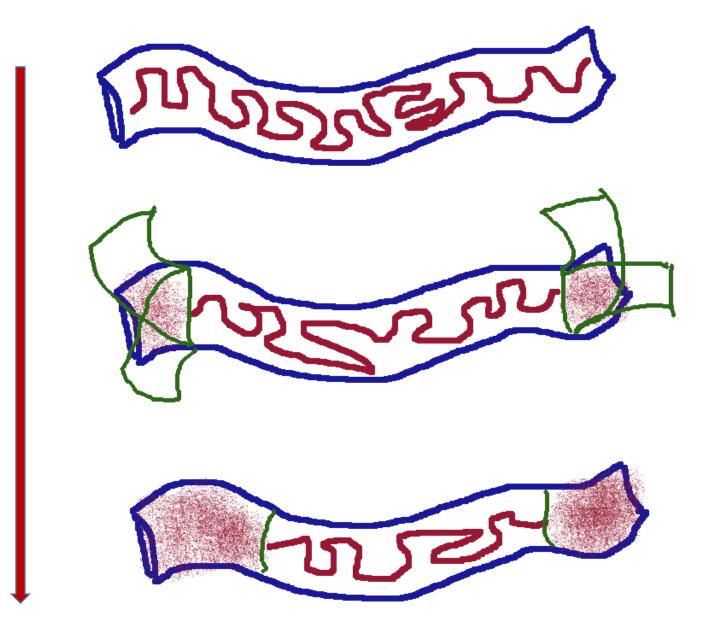
# **Reptation:**

1-d diffusion of kinks. Length to diffuse  $\approx$  Z (Chain ends are under entropic tension) Friction  $\approx$  Z  $\tau_d \approx Z^3$ 

# Contour length fluctuation

Experiments:  $\tau_d \sim Z^{3.4}$ 

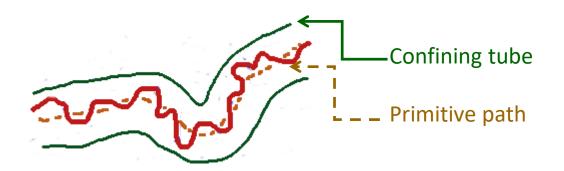
time

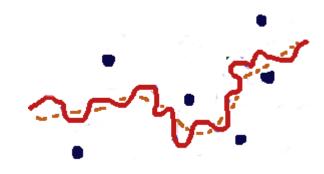


Random walk of end bead does not cost much.

Doi, Edwards

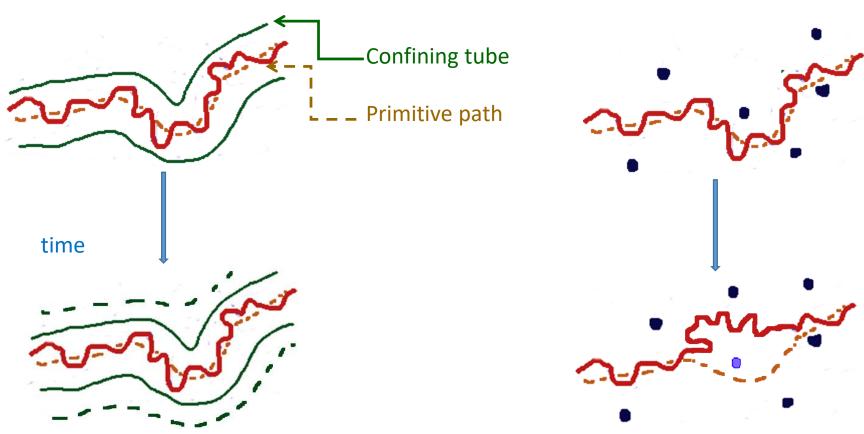
## Constraint release





Discrete view

#### Constraint release



Tube dilation (Softening of confining potential with relaxation)

 $\varphi_{\text{ST}}$  : Density of entanglements

Experiments on branched polymer relaxation:

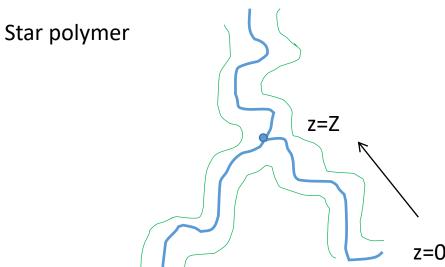
Effective number of entanglement on the chain: Z  $\rightarrow$  Z  $\varphi_{ST}^{\alpha}$ 

Stress decay  $G(t) \sim G0 \phi(t) \phi_{ST}(t)$ 

 $\phi(t) \rightarrow$  fraction of surviving old tube

Marucci

#### Deep retraction



Retraction potential: 
$$\beta \frac{dU}{dz} = \frac{3z}{Z} \phi^{\alpha}$$

Assume all friction located at chain end.

$$\tau_K(z) = \left(\frac{3\pi^5 Z^3}{2\phi_0^{\alpha}}\right)^{1/2} \tau_e \int_0^z dz' e^{\beta U(z')}$$

Softening due to constraint release goes in the exponential.

Dynamic tube dilation is extremely important for branched polymers.

Kuzzu, Doi; Pearson, Helfand; Ball, McLeish

#### Side branches

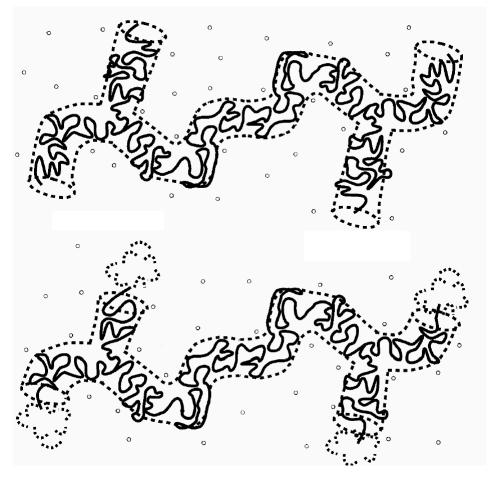


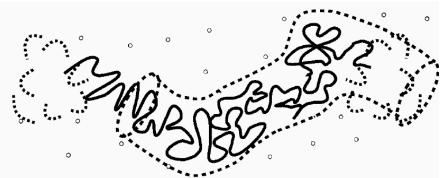
Short art retracts at time  $t_a$ At this timescale, branchpoint can take random hop of order a

Associated friction : 
$$\zeta_a \sim \frac{k_B T}{a^2} t_a$$

"backbone" reptation slows down from this extra friction.

# H-polymer





"backbone" can only escape once side-arms have had chance to relax.



Low density Polyethylene (LDPE)

Molecules made up of 100 to a million carbons per molecule (all of them in the same sample). Number of branch points per molecule between zero (linear molecule) and a few hundreds.

#### **Problems:**

We do not know the actual molecules.

Even if someone told us the detailed composition and the tube picture is correct, it will be humanly impossible to figure out the stress response.

(Number of different relaxation models connected by if-then-else statements.)

Even for a two or three component blend, it is difficult to write down a closed form for stress decay. Each new branching topology needs lengthy calculations.

### Ron Larson 2001 Combinatorial rheology of branched polymer melts

Take differential approach and let the computer do the hard work.

After a step strain,

1. Figure out what can relax in a small time interval.

2. Update tube diameter accounting for the relaxed portions.

Loop till everything has relaxed.

Can have different molecules in the computation having different mass/branching (Larson ignored multiple levels of branching)

Each affecting relaxation of all other molecules by affecting the dynamics of the tube diameter.

- ➤ Have some way to generate an ensemble of representative molecules on computer
- ➤ Incorporate branch-on-branch structures
- ➤ Figure out ways to predict stress in non-linear flows (Polymer processing almost always involve highly non-linear flow)



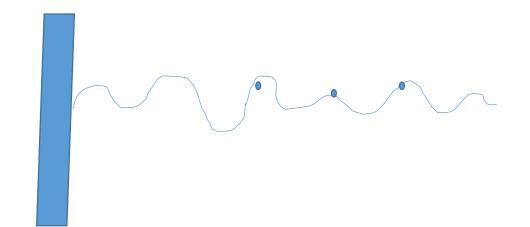
BoB

http://sourceforge.net/projects/bob-rheology

**J Rheol., 50,** 207 (2006); **58,** 737 (2014). **Science 333,** 1871 (2011)

#### Multiple friction along the backbone

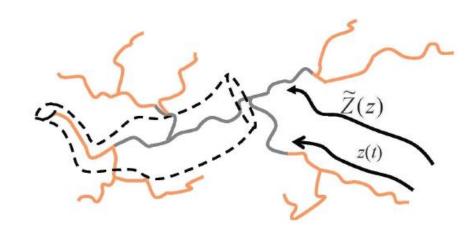
Find equivalent one dimensional problem by insisting that dissipation in retracting a certain length remain unchanged



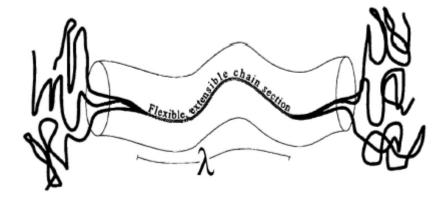
Distance to effective arm retraction pivot :  $\widetilde{Z}(t)$ 

Modify retraction potential as: 
$$\beta \frac{dU}{dz} = \frac{3z\phi^{\alpha}}{\tilde{Z}}$$

Added dynamics for the pivot point  $\widetilde{Z}(t)$ 



## Nonlinear flow Pom-pom model



Flow aligns and stretches confining tube
Stretch can relax by withdrawing the branches
Maximum stretch grows with number of side-branches

#### Backbone:

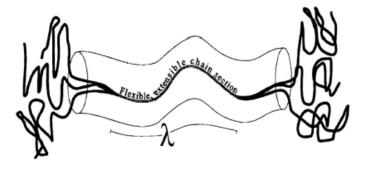
Stretch relaxation time → side arm retraction time

Orientation relaxation time → reptation time

Pretend each bits of any molecule are something like a pompom molecule

#### Pompom model:

- Orientation relaxation time
- Stretch relaxation time
- Number of side-arms



Pretend each bits of any molecule are something like a pompom molecule

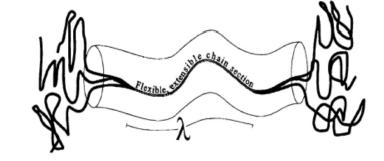
#### Pompom model:

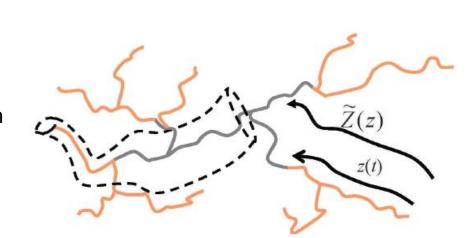
- Orientation relaxation time
- Stretch relaxation time
- Number of side-arms

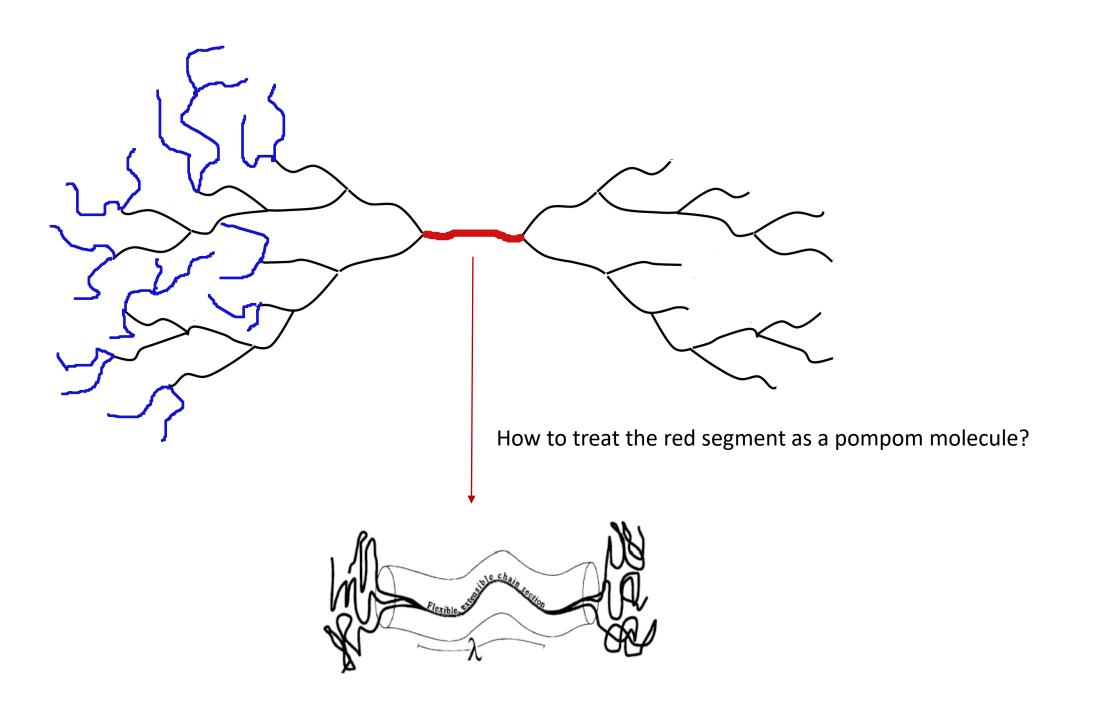
Calculated while finding linear relaxation

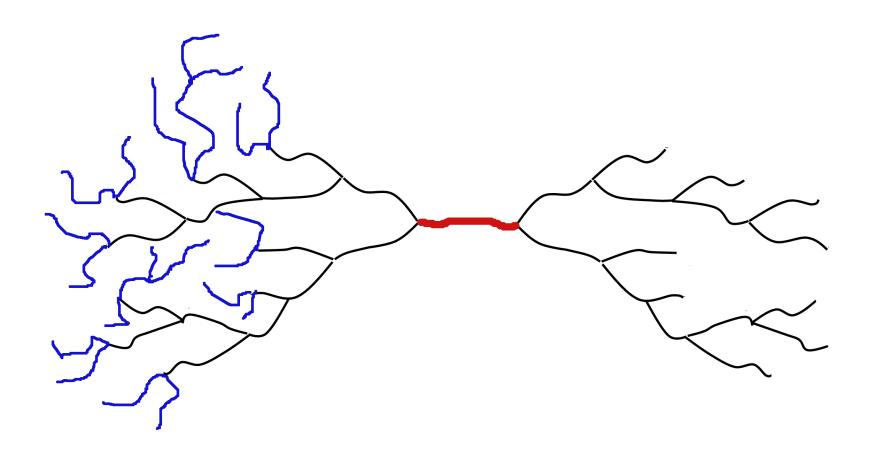
Invert z(t)

Invert Z(t)





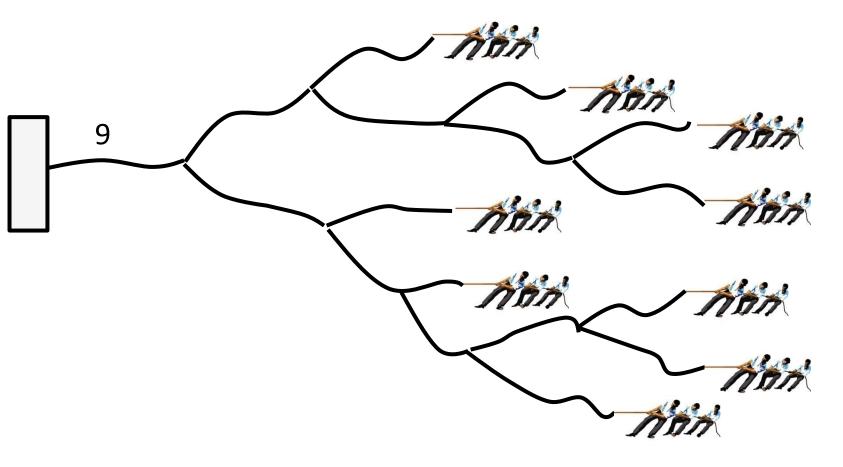




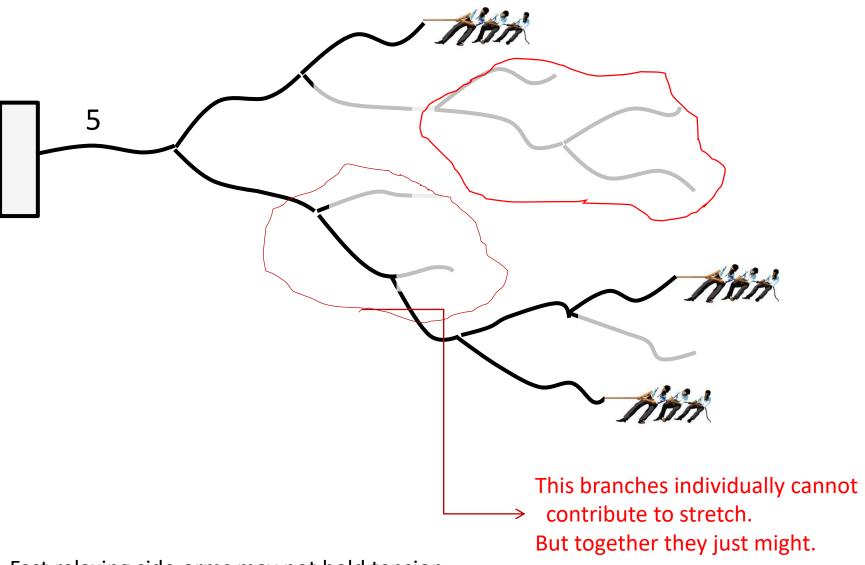
When stretched, first the right hand branch-point will withdraw in the red tube. (On left, there are many more branches)

Concentrate on the right end of the red segment.

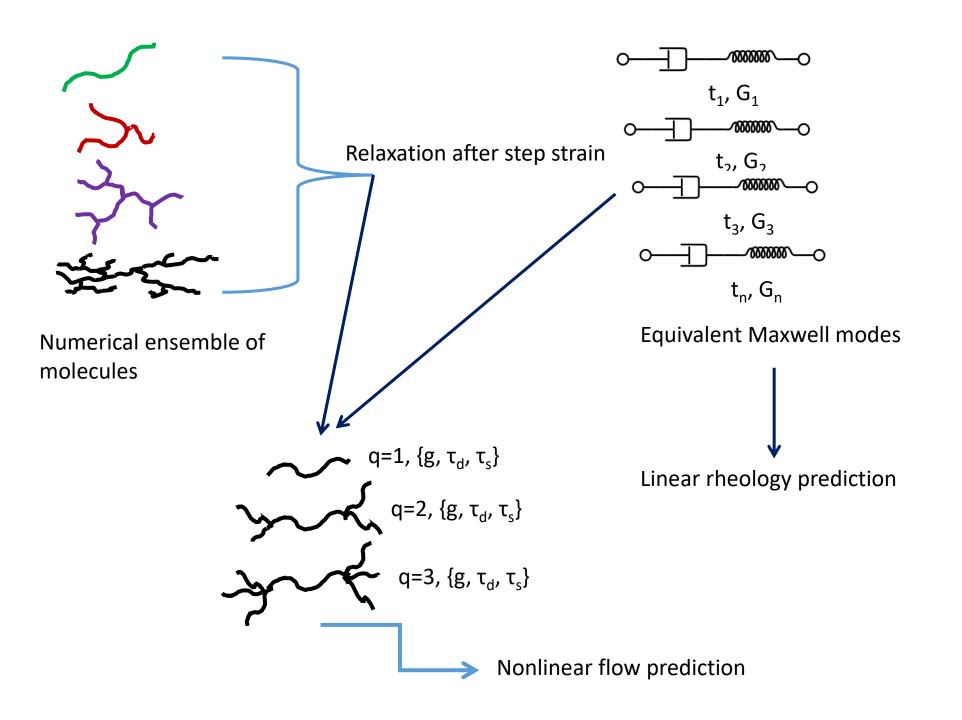
#### Maximum stretch



Each chain end provide 3 kT/a tension

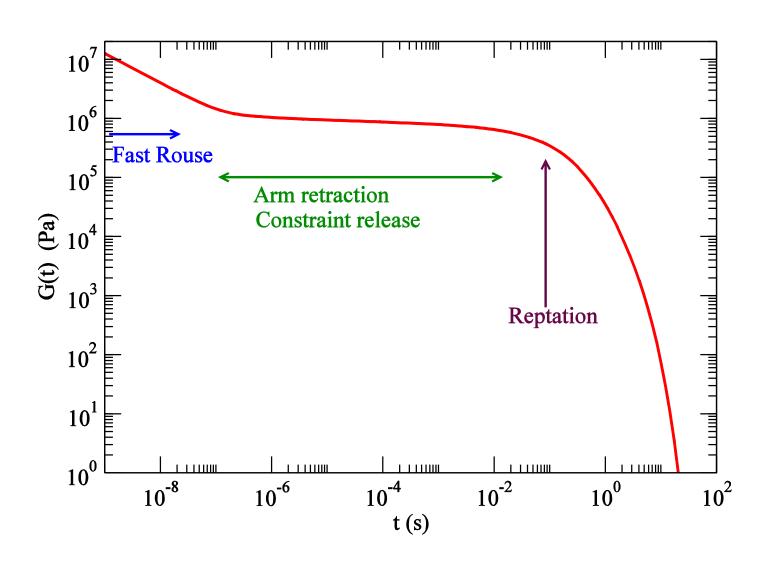


Fast relaxing side-arms may not hold tension Fast is relative concept – relative to flow-rate.



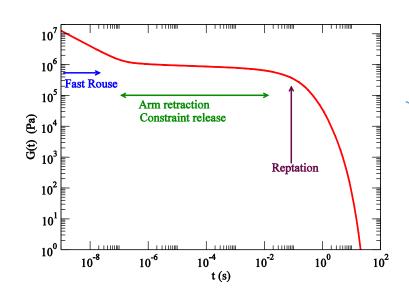
# Experiments:

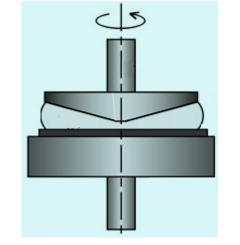
Linear Rheology



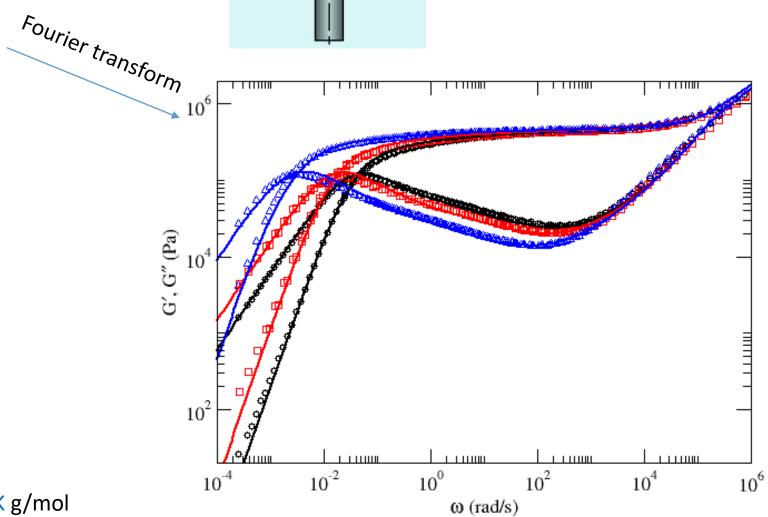
#### **Experiments:**

**Linear Rheology** 



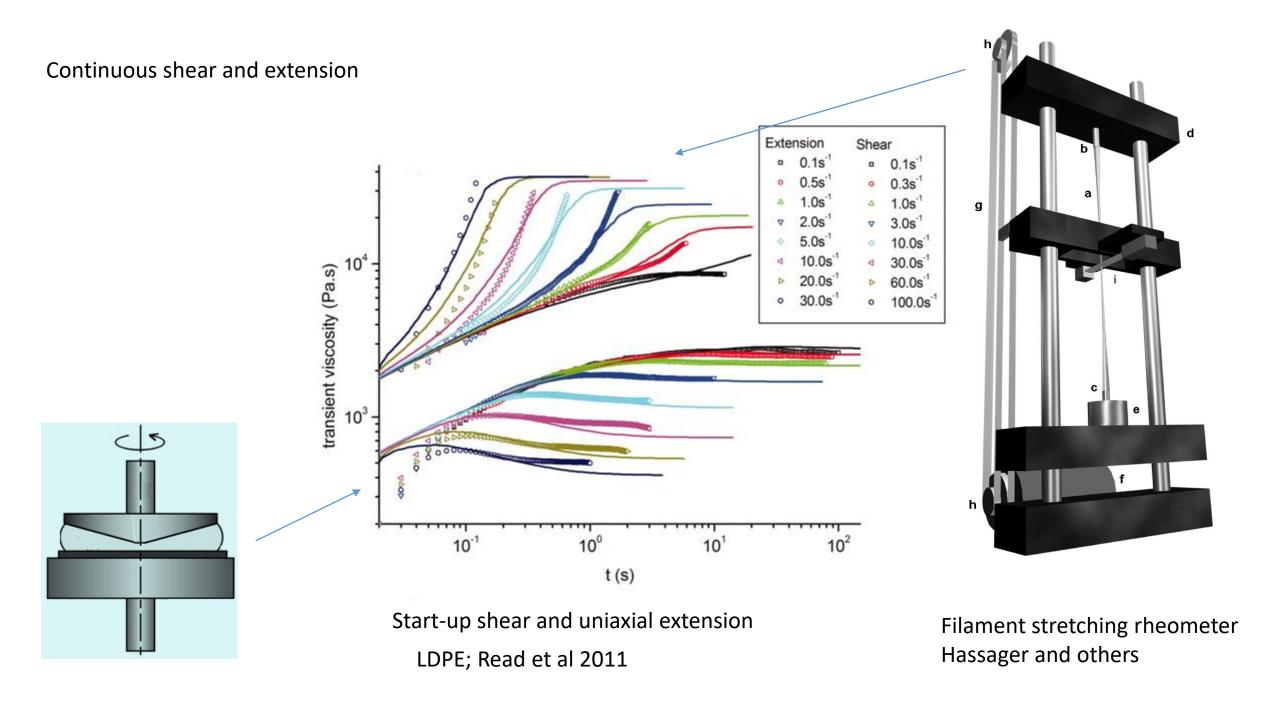


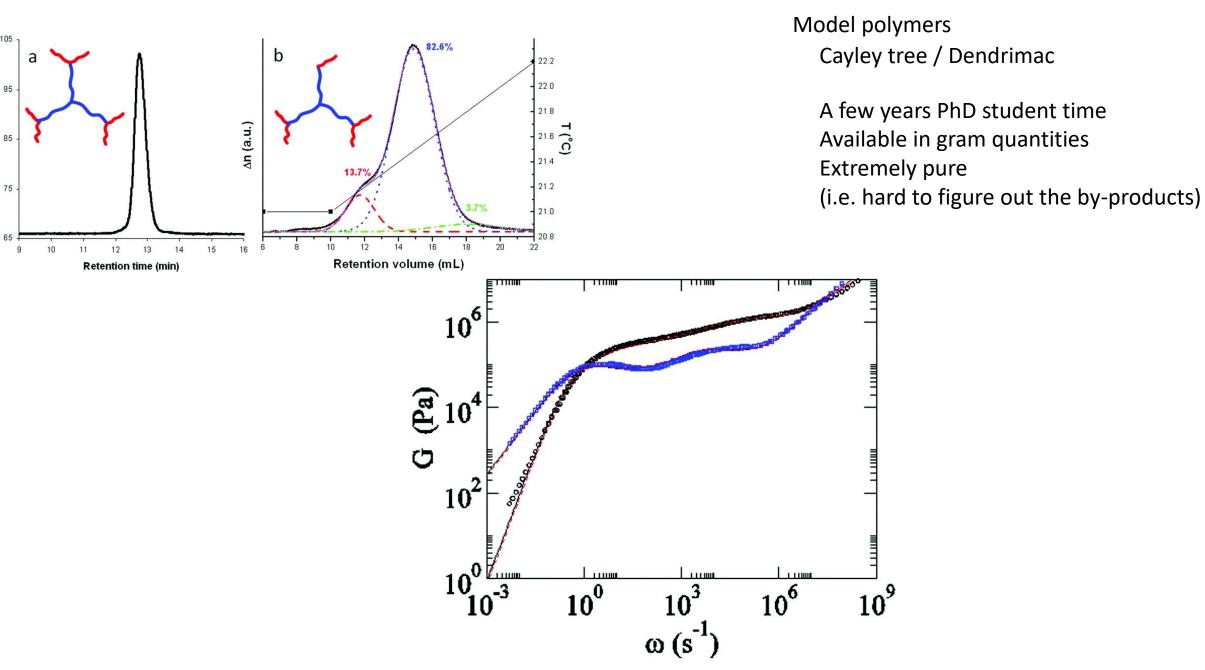
Cone and plate rheometer Small amplitude oscillatory shear



Experimental data: Auhl et al 2008

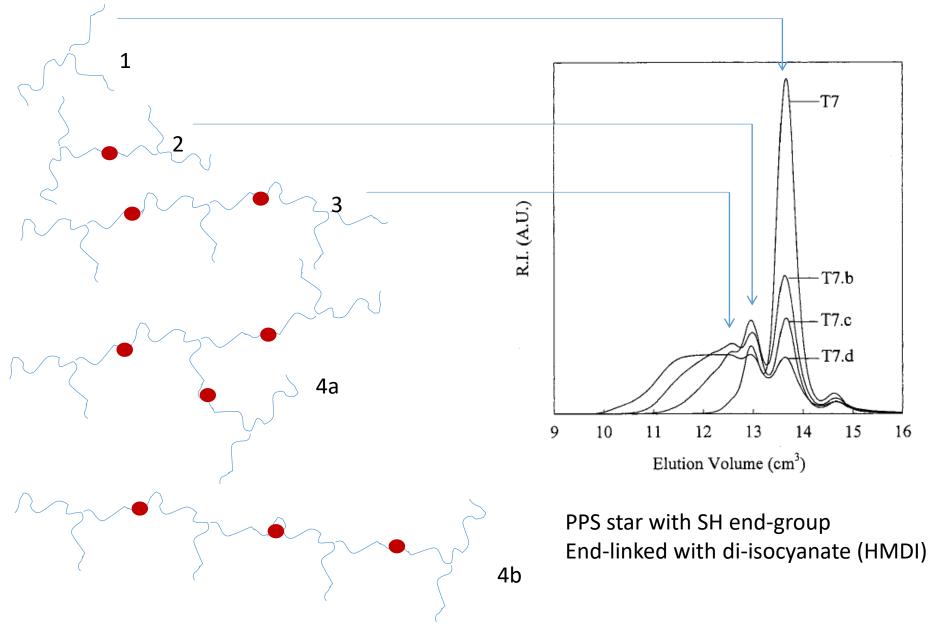
Polyisoprene linear: 483K, 634K, 1131K g/mol



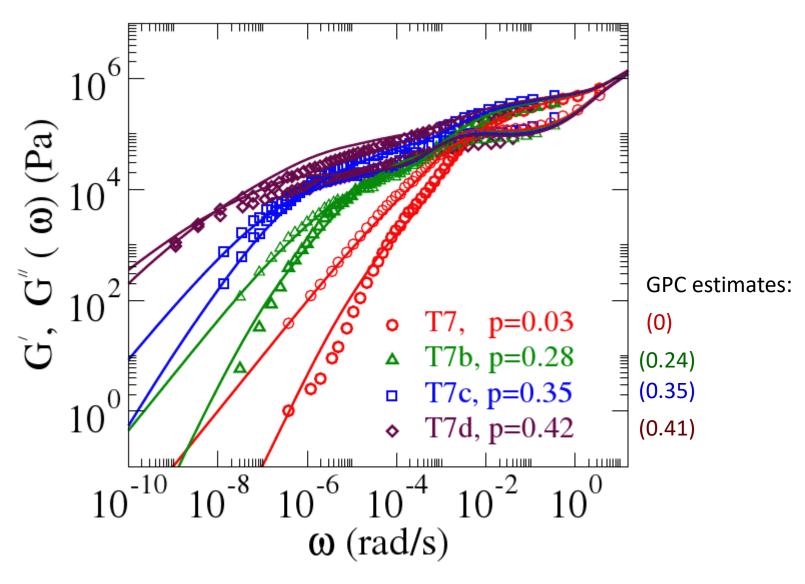


Hutching et al.; ACS Macro Lett.1, 404 (2012)

### End-linked star polymers

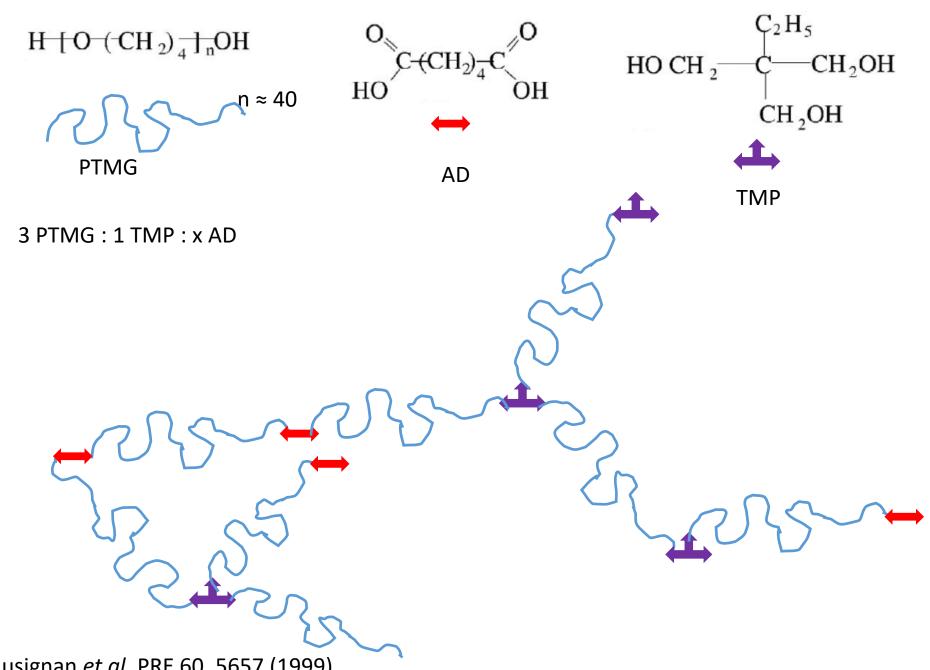


Nicol, Nicolai, Durand; Macromolecules 34, 5205 (2001)



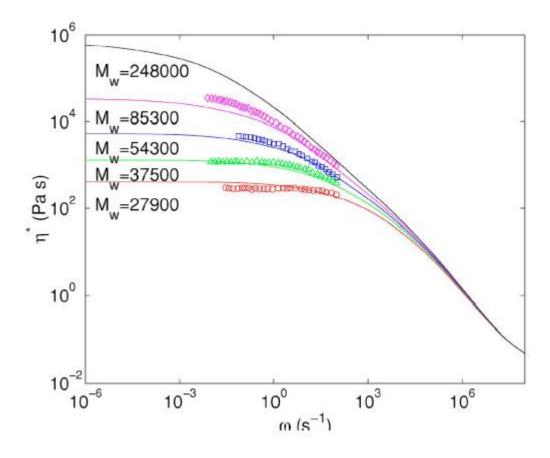
Trace oxygen can link some of the stars in parent resin. den Doelder, Das, Read; Rheo Acta **50**, 469 (2011)

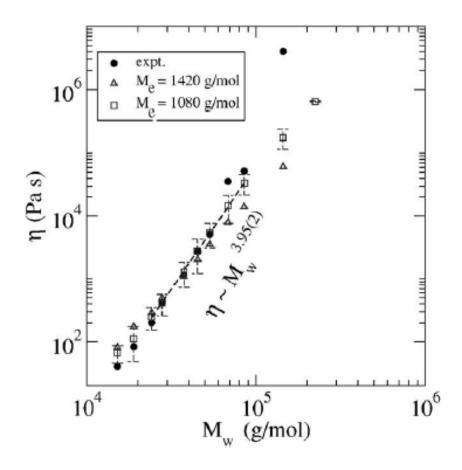
#### Randomly branched polymer from vulcanization class



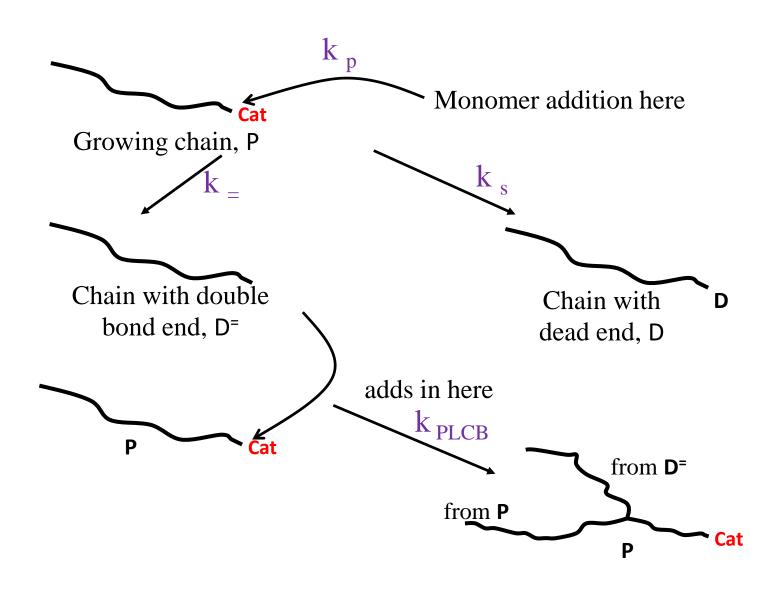
Lusignan et al. PRE 60, 5657 (1999)

Fit structural data to find rate constants.

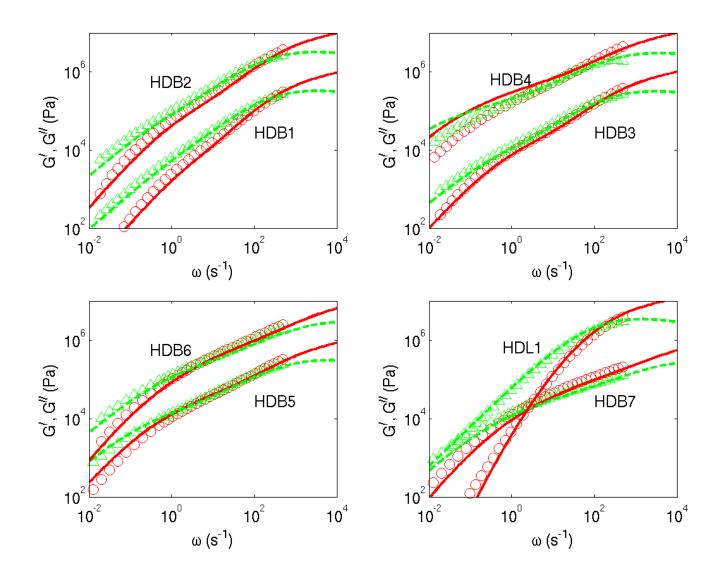


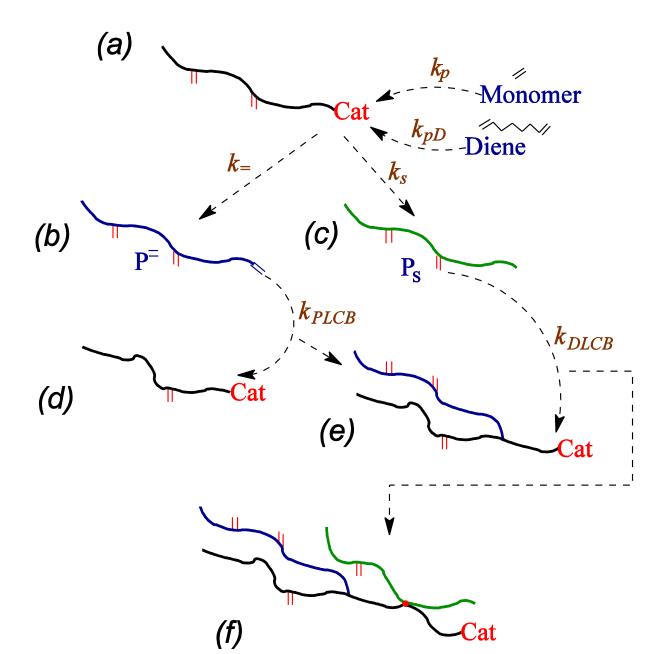


Das et al. PRE 74, 011404 (2006)



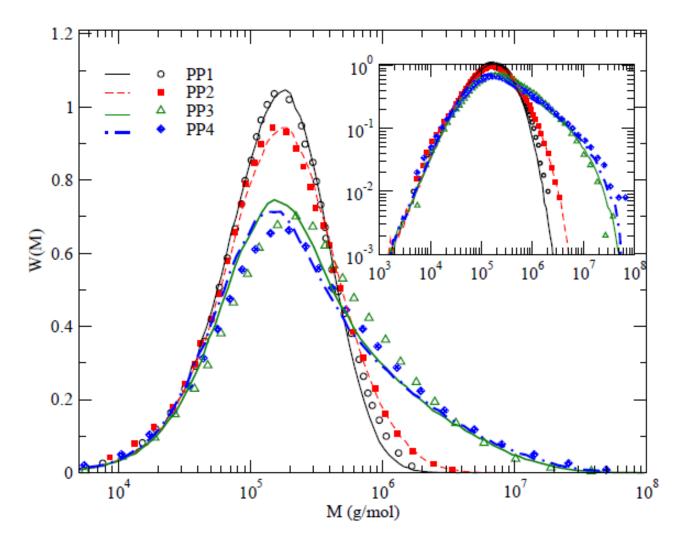
Convert rate constants to probabilities. Use recursive Monte Carlo to generate representative molecules.



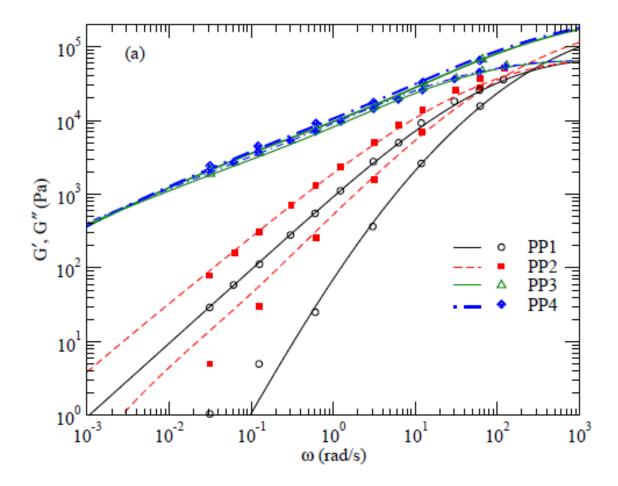


Chains with terminal double bonds are rare. Copolymerize with diene to have lots of pendant double bonds.

## Propylene-diene

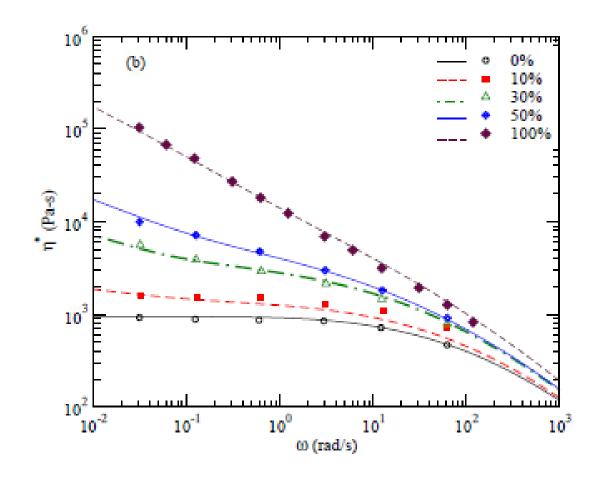


data from Ye et al., Ind. Eng. Chem. Res. 43, 2860 (2004)



Since we have the molecules, we can calculate SAOS response.

Blend our numerical resins numerically and compare with experimental blends...



Das, Read, Soulages, Shirodkar; Macromolecules, 47, 5860 (2014)

How to define gel-point?

Reduced distance to gel-point: 
$$\varepsilon = \frac{[D]_{c} - [D]}{[D]_{c}}$$

Molar mass develops a power-law tail: 
$$\phi_N(M) \sim M^{-\tau} f\left(\frac{M}{M_{\text{char}}}\right)$$

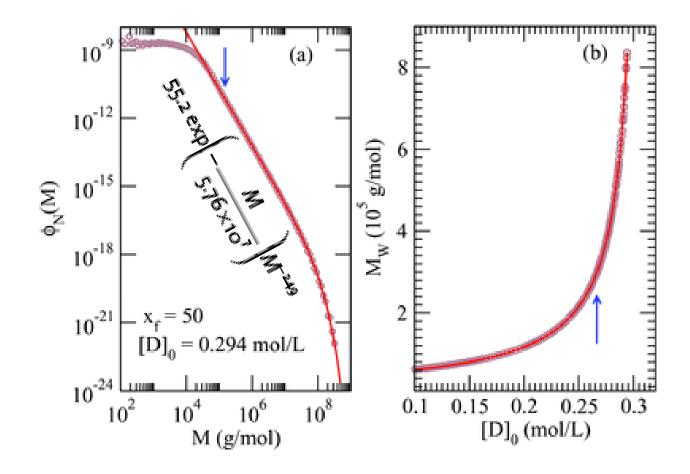
Largest molar mass where the power-law is valid diverges:  $M_{\rm chur} \sim \varepsilon^{-1/\sigma}$ 

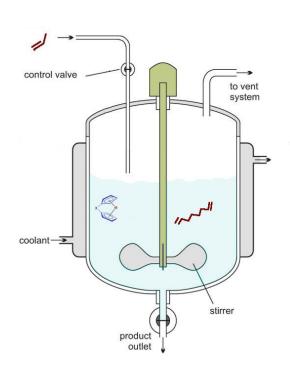
Low-order molar mass moments may remain finite depending on the exponent  $\tau$ Polymer community often accepts that gelation means diverging  $M_w$  (2<sup>nd</sup> moment)

#### Semibatch reactor:

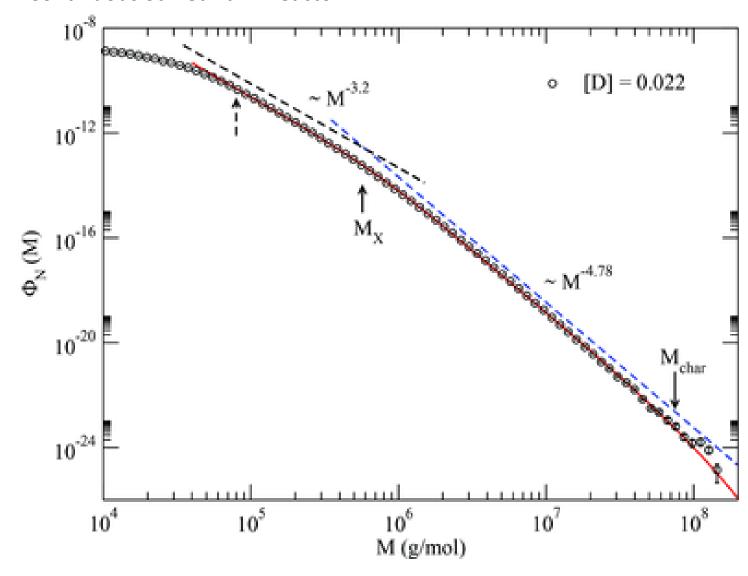
Mean-field gelation (percolation on Bethe lattice) remains valid  $\tau$ =5/2.

M<sub>W</sub> diverges (but M<sub>N</sub> remains finite) at the gel-point

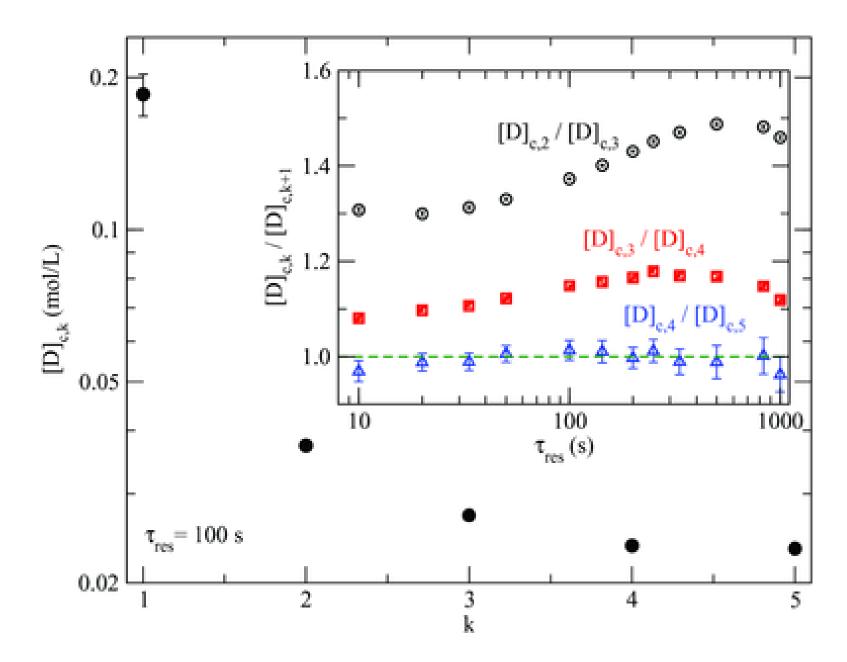




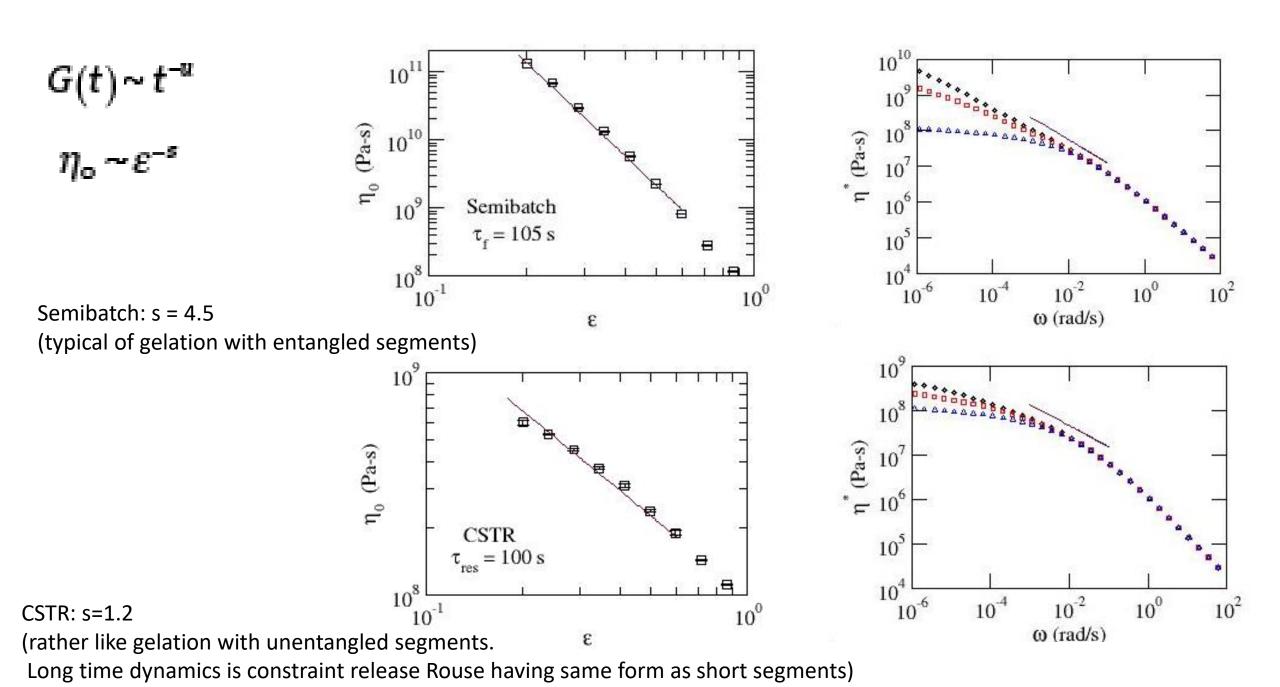
### Continuous Stirred Tank Reactor

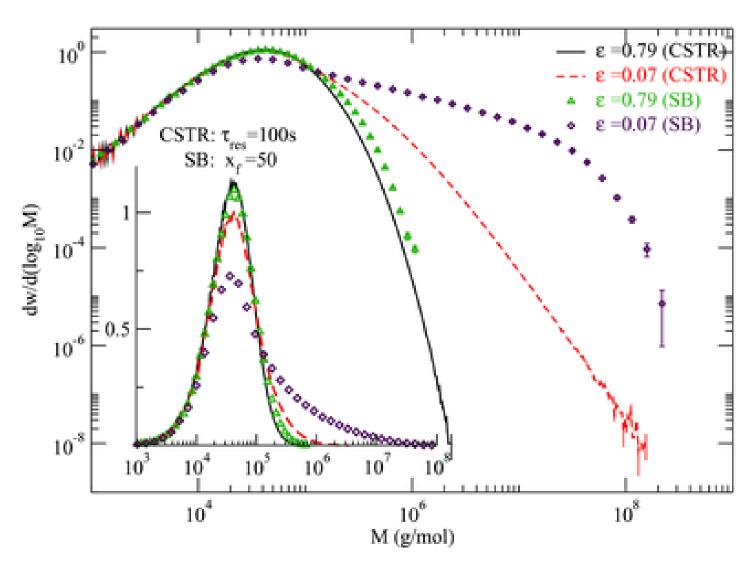


 $\tau$  = 4.5 Valid only very close to gelation First moment to diverge is the 4<sup>th</sup> moment



Apparent gel-point from extrapolating from kth moment





Experimentally should be hard to detect gelation in CSTR Exponential distribution of residence time.

More time a rare molecule remains in the reactor, more branched it becomes.

LDPE: the stuff plastic bags are made of

Free radical polymerization at high pressure

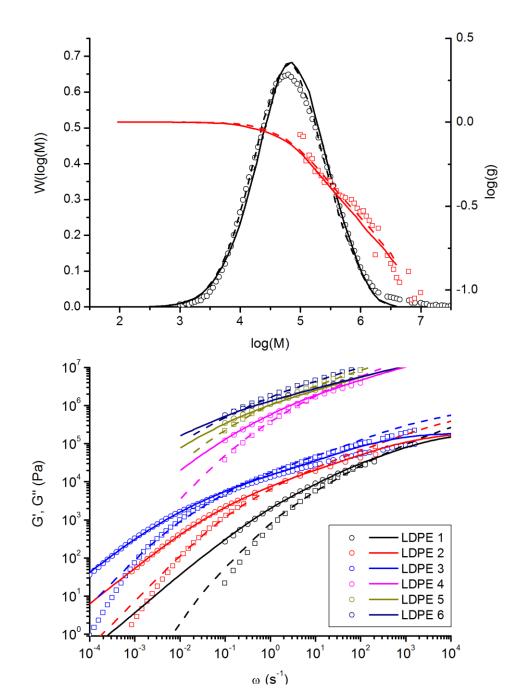
Initiator → free radical

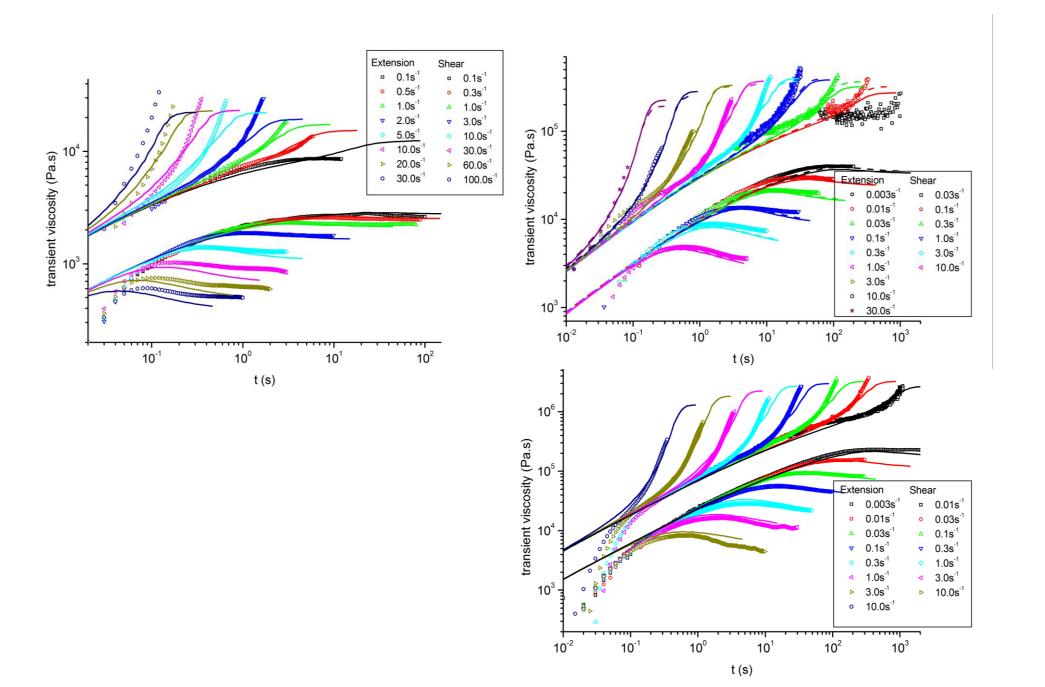
Propagation :: add monomers

Termination by combination or disproportionation

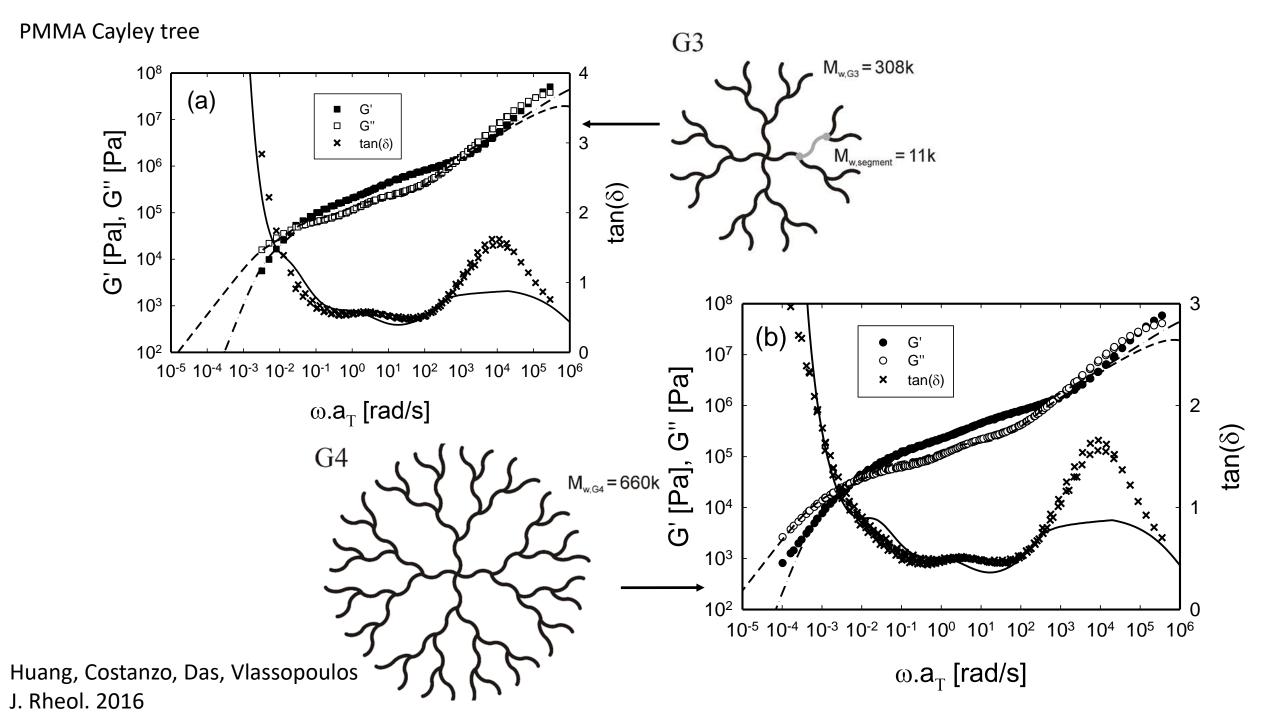
Chain transfer leading to side branches or scission

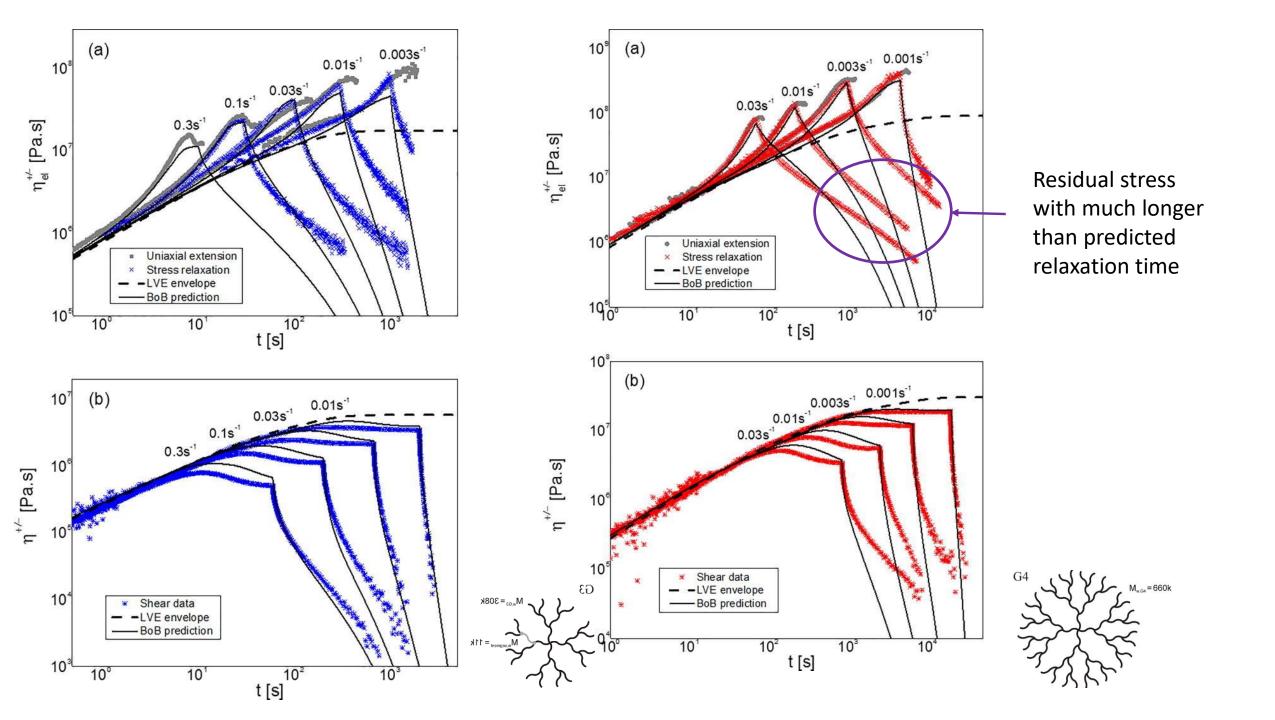
Rate constants by fitting molar mass distribution and the g-factor

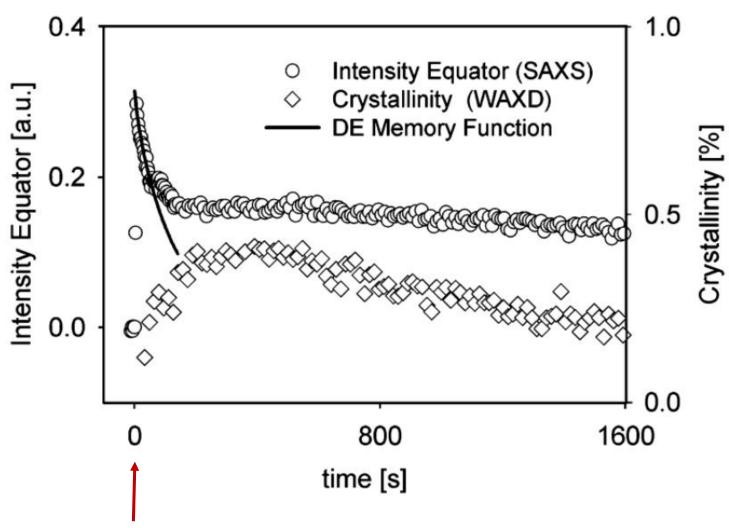












Step shear above crystallization temp.

Bimodal linear PE.

Shear rate such that high molar mass fraction is stretched.

# Flow induced crystallization in polymers: from molecules to processing

