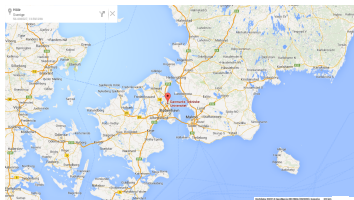
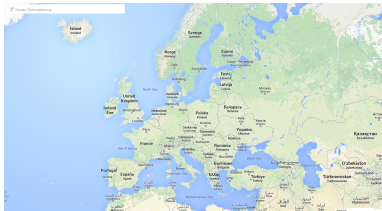


# Numerical methods for EIT

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Advanced Instructional School of Theoretical and Numerical Aspects of  
Inverse Problems  
Bangalore  
June 16-27, 2014

# About me...



**HD TOMO**  
HIGH - DEFINITION TOMOGRAPHY

$$f(x+\Delta x) = \sum_{i=0}^{\infty} \frac{(\Delta x)^i}{i!} f^{(i)}(x)$$

A collection of mathematical symbols including  $\Delta$ ,  $\int_a^b$ ,  $\epsilon$ ,  $\Theta$ ,  $\sqrt{17}$ ,  $\Omega$ ,  $\int$ ,  $\delta e^{in}$ ,  $\infty$ ,  $\chi^2$ ,  $\Sigma$ , and  $!$ .

# Outline

1. The Calderón problem and linearization
2. The CGO-method for reconstruction in 2D
3. The CGO-method for reconstruction in 3D
4. The Calderón problem with partial data

# 1. The Calderón problem and linearization

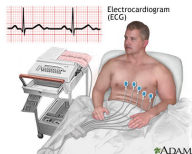


# Motivation

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The use of electricity in treatment and monitoring of patients:

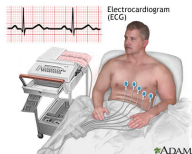
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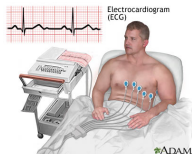
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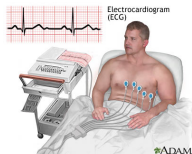
- Electrocardiography (EKG)
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- Defibrillator



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The use of electricity in treatment and monitoring of patients:

- Electrocardiography (EKG)
- Electro stimulation
- Defibrillator
- Electrical Impedance Tomography



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- Fricke and Morse reported in the article "The electric capacity of tumours of the breast" (1926) that the electrical properties of breast tumours differ from healthy tissue

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- Barber-Brown (1989):

Tissue	Conductivity (mS/cm)
Blood	6.7
Liver	2.8
Skeletal muscle	8.0 (long.), 0.6 (trans.)
Cardiac muscle	6.3(long.), 2.3 (trans.)
Lung (expiration-inspiration)	1.0 - 0.4
Fat	0.36
Bone	0.06

# Measurement setup EIT





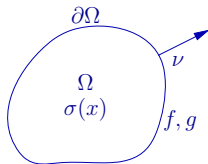
## Mathematical model for EIT

Smooth bounded domain  $\Omega \subset \mathbb{R}^d$ ,  $d = 2, 3$ , conductivity coefficient

$$0 < c \leq \sigma \leq C < \infty; \quad \sigma \equiv 1 \text{ near } \partial\Omega.$$

Voltage potential  $u$  in  $\Omega$  generated by boundary voltage potential  $f$

$$\begin{aligned} \nabla \cdot \sigma \nabla u &= 0 \text{ in } \Omega, \\ u|_{\partial\Omega} &= f. \end{aligned}$$



Current field:  $J = \sigma \nabla u$ .

Normal component of current field at  $\partial\Omega$ :

$$g = \nu \cdot J = \sigma \partial_\nu u|_{\partial\Omega}.$$

Dirichlet to Neumann (voltage to current) map

$$\Lambda_\sigma : f \mapsto g.$$

## Dirichlet to Neumann map

Dirichlet to Neumann (voltage to current) map

$$\Lambda_\sigma: H^{1/2}(\partial\Omega) \rightarrow H^{-1/2}(\partial\Omega)$$
$$f \mapsto g.$$

Weakly defined for  $h \in H^{1/2}(\partial\Omega)$  by

$$\langle \Lambda_\sigma f, h \rangle = \int_{\partial\Omega} (\Lambda_\sigma f) \bar{h} \, ds(x) = \int_{\Omega} \sigma \nabla u \cdot \overline{\nabla v} \, dx, \quad v \in H^1(\Omega) : v|_{\partial\Omega} = h.$$

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Note:

- $\Lambda_\sigma$  bounded operator  $H^{1/2}(\partial\Omega) \rightarrow H^{-1/2}(\partial\Omega)$
- $\Lambda_\sigma$  is unbounded, selfadjoint in  $L^2(\partial\Omega)$
- $\Lambda_\sigma - \Lambda_1$  is compact in  $L^2(\partial\Omega)$

# The Calderón problem

Forward problem:

$$\Lambda: \sigma \mapsto \Lambda_\sigma$$

From interior conductivity to boundary fields.

# The Calderón problem

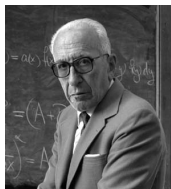
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- **Uniqueness:** is  $\Lambda$  injective?



A P Calderón

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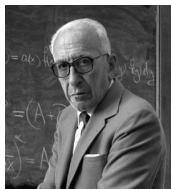
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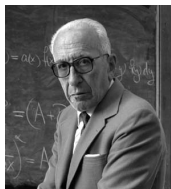
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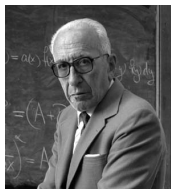
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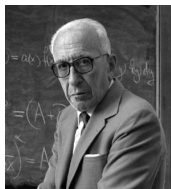


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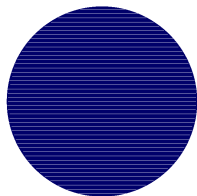
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Calderón problem is ill-posed:

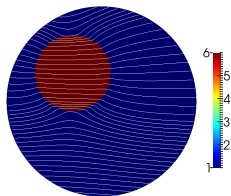
- Large change in conductivity yields small change in data
- Small change in data yields large change in reconstruction

## Example with $f(\theta) = \cos(\theta)$

Current flow  $J = \sigma \nabla u$  :



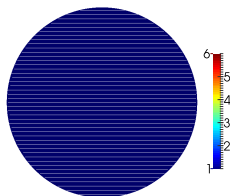
Homogeneous conductivity



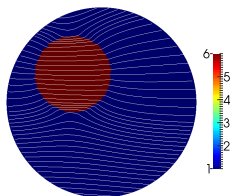
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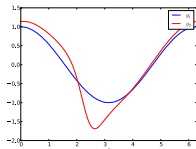


Homogeneous conductivity

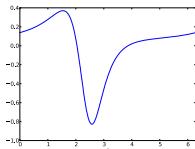


Perturbed conductivity

Boundary normal current  $g = \sigma \partial_\nu u$  :



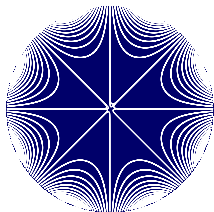
Boundary currents



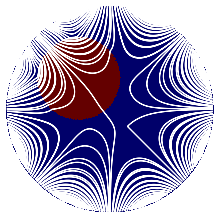
Difference

## Example with $f(\theta) = \cos(4\theta)$

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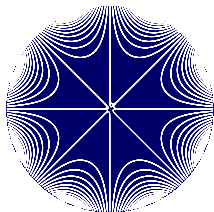
Homogeneous conductivity



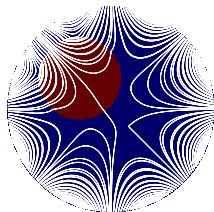
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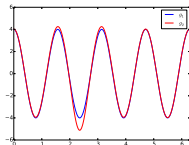


Homogeneous conductivity

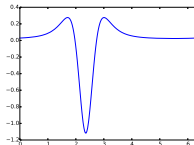


Perturbed conductivity

Boundary normal current  $g = \sigma \partial_\nu u$ :



Boundary currents



Difference

# Short and incomplete history

1980 Calderón: Problem posed, uniqueness for linearized problem, and linear, approximate reconstruction algorithm

## 3D

- 1987 Sylvester and Uhlmann: Uniqueness for smooth conductivities. Implicit reconstruction algorithm  
1987-88 Novikov, Nachman-Sylvester-Uhlmann, Nachman: Uniqueness for conductivities with 2 derivatives and explicit high frequency reconstruction algorithm. Multidimensional D-bar equation.  
1990 Alessandrini: Stability  
2003 Brown-Torres, Päivärinta-Panchenko-Uhlmann: Uniqueness for conductivities with 3/2 derivatives.  
2006 Cornean-Knudsen-Siltanen: Low frequency reconstruction algorithm  
2010 Bikowski-Knudsen-Mueller: Numerical implementation of simplified reconstruction algorithm  
2011-14 Debary - Hansen- Knudsen: Implementation of more accurate numerical reconstruction method  
2012 Haberman - Tataru: Uniqueness for Lipschitz conductivities

## 2D

- 1996 Nachman: Uniqueness and reconstruction for  $W^{2,p}(\Omega)$  conductivities.  
1997 Liu: Stability for  $W^{2,p}(\Omega)$  conductivities  
1997 Brown-Torres: Uniqueness for  $W^{1,p}(\Omega)$  conductivities  
2001 Barceló-Barceló-Ruiz: Stability for  $C^{1+\epsilon}$  conductivities  
2001 Knudsen-Tamasan: Reconstruction for  $C^{1+\epsilon}$  conductivities  
2005 Astala-Päivärinta: Uniqueness and reconstruction for  $L^\infty(\Omega)$   
2009 Knudsen-Lassas-Mueller-Siltanen: Regularized  $\bar{\partial}$ -method  
2010 Clop-Faraco-Ruiz: Stability for discontinuous conductivities

+ many more

## Reconstruction algorithms

- **Linearization:**

Replace non-linear operator  $\Lambda : \sigma \mapsto \Lambda_\sigma$  by linear operator and solve linear inverse problem.

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Minimize over  $\gamma$

$$\Phi(\gamma) = \|\Lambda_\gamma - \Lambda_\sigma\|_Y + \alpha \|\gamma\|_X$$

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- **Bayesian inversion:**

$$\pi(x|y) = \frac{\pi(y|x)\pi_{\text{pr}}(x)}{\pi(y)}$$

for  $x = \sigma$  and  $y = \Lambda_\sigma$ .

## Linearization

Green's formula yields

$$\langle (\Lambda_{\sigma_1} - \Lambda_{\sigma_0})f_1, f_0 \rangle = \int_{\Omega} (\sigma_1 - \sigma_0) \nabla u_1 \cdot \overline{\nabla u_0} dx,$$
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For  $\sigma_0 = 1$

$$\langle (\Lambda_{1+\delta\sigma} - \Lambda_1)f_1, f_0 \rangle \approx \langle (d\Lambda[1]\delta\sigma)f_1, f_0 \rangle = \int_{\Omega} \delta\sigma \nabla v_1 \cdot \overline{\nabla v_0} dx$$

with  $\Delta v_j = 0$ ,  $v_j|_{\partial\Omega} = f_j$ .

## The linearized problem

Concerns the inversion of the mapping

$$\delta\sigma \mapsto (d\Lambda[1])(\delta\sigma).$$

We want to compute  $\delta\sigma$  by knowing:

$$\int_{\Omega} \delta\sigma \nabla v_1 \cdot \overline{\nabla v_0} dx$$

for all harmonic functions  $v_0, v_1$ .

Questions:

- Uniqueness?
- Stable reconstruction?

## Comparison

Problem: find  $\sigma$  from knowing

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2. Linearized Calderón problem: known diffuse kernel

$$k((i, j), y) = \nabla v_i \cdot \overline{\nabla v_j}, \quad \Delta v = 0.$$

3. Computerized Tomography: known localized kernel

$$k(\theta, r, y) = \delta(y \cdot \theta - r).$$

## Exponentially growing harmonics

Fix real vector  $\xi \in \mathbb{R}^n$  and complex vector  $\zeta \in \mathbb{C}^n$  such that

$$\zeta \cdot \zeta = (\xi + \zeta) \cdot (\xi + \zeta) = 0.$$

Harmonic functions in  $\mathbb{R}^n$ .

$$v_1(x, \zeta) = e^{ix \cdot \zeta}, \quad v_0(x, \zeta) = e^{ix \cdot (\xi + \bar{\zeta})}$$

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and

$$\begin{aligned} \langle d\Lambda_1[\delta\sigma] e^{ix \cdot \zeta}, e^{ix \cdot (\xi + \bar{\zeta})} \rangle &= -\frac{1}{2} |\xi|^2 \int_{\Omega} \delta\sigma e^{-ix \cdot \xi} dx \\ &= -\frac{1}{2} (2\pi)^{n/2} |\xi|^2 \widehat{\delta\sigma}(\xi) \end{aligned}$$

$$\Leftrightarrow \widehat{\delta\sigma}(\xi) = -\frac{2}{(2\pi)^{n/2} |\xi|^2} \langle d\Lambda_1[\delta\sigma] e^{ix \cdot \zeta}, e^{ix \cdot (\xi + \bar{\zeta})} \rangle.$$

## Uniqueness and stable reconstruction

$$\widehat{\delta\sigma}(\xi) = -\frac{2}{(2\pi)^{n/2}|\xi|^2} \langle d\Lambda_1[\delta\sigma]e^{ix\cdot\xi}, e^{ix\cdot(\xi+\bar{\zeta})} \rangle.$$

- **Uniqueness:** Injectivity of Fourier transform.

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**Stabilization:** Noise amplified exponential harmonics.  
Remedy: Avoid high frequencies by low-pass filtering:

$$\delta\sigma_R(x) = -\frac{2}{(2\pi)^n} \int_{\mathbb{R}^n} \frac{\langle d\Lambda_1[\delta\sigma] e^{ix\cdot\xi}, e^{ix\cdot(\xi+\bar{\zeta})} \rangle}{|\xi|^2} e^{ix\cdot\xi} \chi_R(\xi) d\xi.$$

**Regularization scheme** if  $R$  is chosen correctly.



# Calderón's reconstruction method

Treat non-linear data

$$\mathbf{t}^{\text{exp}}(\xi, \zeta) := \langle (\Lambda_\sigma - \Lambda_1) e^{ix \cdot \zeta}, e^{ix \cdot (\xi + \bar{\zeta})} \rangle = \int_{\Omega} \delta\sigma \nabla u \cdot \nabla e^{-ix \cdot (\xi + \zeta)} dx$$

as linear data in previous formula

$$\sigma^{\text{Cal}}(x) = 1 - \frac{2}{(2\pi)^n} \int_{\mathbb{R}^n} \frac{\mathbf{t}^{\text{exp}}(\xi, \zeta)}{|\xi|^2} e^{ix \cdot \xi} \chi_R(\xi) d\xi.$$

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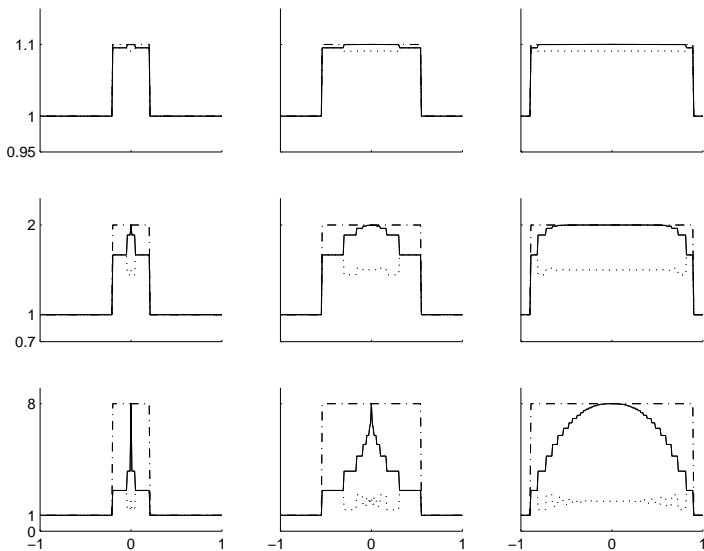
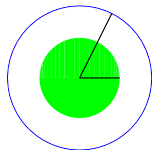
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Estimate:

$$\begin{aligned} \|\sigma^{\text{Cal}}(x) - \eta_\gamma * \sigma\|_{L^\infty} &\leq \|R(k) \hat{\eta}(k/\gamma)\|_{L^1(\mathbb{R}^2)} \\ &\leq C \|\sigma - 1\|_{L^\infty(\Omega)}^{1+\alpha} (\log(\|\sigma - 1\|_{L^\infty(\Omega)}))^2. \end{aligned}$$

# Concentric reconstruction (linear)



## Non concentric reconstruction (linear)

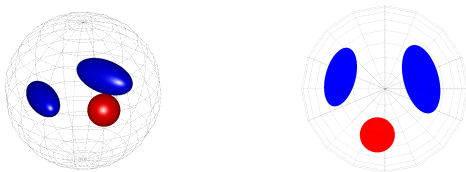
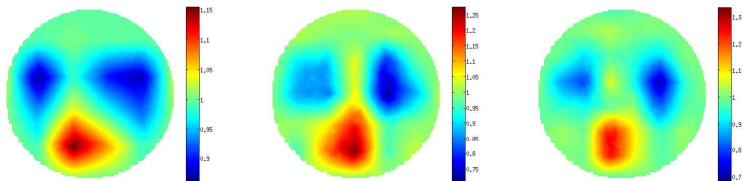


Figure: Increasing Fourier truncation for Calderón's method



## Conclusion, linearized reconstruction

- Easy to implement and fast: based on FFT
- Recovers (with no regularization) the support of perturbation  $\delta\sigma$  (cf [von Harrach - Seo, 2010] and [Knudsen - Lassas - Mueller - Siltanen, 2007])
- Recovers well low contrast perturbations
- Regularized algorithm recovers smooth approximation
- Can rigorous mathematics allow better reconstructions?