

Framework

Hubble rate $H \equiv \frac{a}{a}$

<u>Matter component: perfect fluid</u> $T_{\mu\nu} = pg_{\mu\nu} + (\rho + p) u_{\mu}u_{\nu}$

equation of state $p = w \rho$

+ cosmological constant = Einstein equations

$$\begin{cases} H^2 + \frac{\mathcal{K}}{a^2} = \frac{1}{3} (8\pi) \\ \frac{\ddot{a}}{a} = \frac{1}{3} [\Lambda - 4\pi G_1] \end{cases}$$



$$\int \frac{w}{w} = \frac{1}{3}$$
 dust
$$\frac{w}{3} = \frac{1}{3}$$
 radiation

 $\pi G_{\rm N} \rho + \Lambda$

 $[(\rho + 3p)]$

Particular solution: dust and radiation

integrate conservation equation



$$w(a)] d \ln a \bigg\} \underset{w \to \text{cst}}{=} \rho_{\text{ini}} \left(\frac{a}{a_{\text{ini}}}\right)^{-3(1+w)}$$

Phenomenologically valid description for 14 Gyrs!!!!



r < 0.08

Numerical simulation for large scale structure formation...



Standard Failures and inflationary solutions

Singularity Horizon $d_{\rm H} \equiv a(t) \int_{t_{\rm i}}^{t} \frac{d\tau}{a(\tau)}$ can be made as big as one wishesFlatness $\frac{\rm d}{{\rm d}t} |\Omega - 1| = -2 \frac{\ddot{a}}{\dot{a}^3}$ $\ddot{a} > 0$ & $\dot{a} > 0$

Homogeneity & Isotropy

Initial Universe = very small patch Accelerated expansion drives the shear to zero...

Perturbations **Others** dark matter/energy, baryogenesis, ...

- Not solved... actually not addressed!

accelerated expansion (inflation)



vacuum state!

- + attractor
- Bonus of the theory: predictions!!!

• solves cosmological puzzles

makes falsifiable predictions ... •• • ... consistent with all known observations



Inflation

•••

uses GR + scalar fields [(semi-)classical] • can be implemented in high energy theories string implementation (brane inflation, ...)

y bother with alternatives?

From R. Brandenberger, *in* M. Lemoine, J. Martin & PP (Eds.), "Inflationary cosmology", Lect. Notes Phys. **738** (Springer, Berlin, 2007).



A brief history of bouncing cosmology

R. C. Tolman, "On the Theoretical Requirements for a Periodic Behaviour of the Universe", PRD 38, 1758 (1931) G. Lemaître, "L'Univers en expansion", Ann. Soc. Sci. Bruxelles (1933)

-> A. A. Starobinsky, "On one non-singular isotropic cosmological model", Sov. Astron. Lett. 4, 82 (1978) V. N. Melnikov, S.V. Orlov, Phys. Lett. A 70, 263 (1979). R. Durrer & J. Laukerman, "The oscillating Universe: an alternative to inflation", Class. Quantum Grav. 13, 1069 (1996)

...

Many new ideas, models...

-> M. Novello & S.E. Perez Bergliaffa, "Bouncing cosmologies", Phys. Rep. 463, 127 (2008) -> D. Battefeld & PP, "A Critical Review of Classical Bouncing Cosmologies", Phys. Rep. 571, 1 (2015) -> R. Brandenberger & PP, "Bouncing cosmologies: Progress and problems", Found. Phys. (2017)



The issue we are interested in: the singularity



The issue we are interested in: the singularity





Model listing:

Quantum gravity

LQG & LQC

Canonical quantum gravity (WdW) String theory

Non relativistic quantum gravity W)

Singularity problem...



a quantum effect?





Mod.Phys.Lett. A31, 1640006 (2016).

Model listing:

Quantum gravity

Ekpyrotic & cyclic Branes

LQG & LQC

Canonical quantum gravity (WdW) String theory



Non relativistic quantum gravity W)

Ekpyrotic scenario:



 $\mathcal{S}_{_5} \propto$

$$\int_{\mathcal{M}_5} \mathrm{d}^5 x \sqrt{-g_5} \left[R_{(5)} - \frac{1}{2} \left(\partial \varphi \right)^2 - \frac{3}{2} \frac{\mathrm{e}^{2\varphi} \mathcal{F}^2}{5!} \right],$$

$$= \int_{\mathcal{M}_4} \mathrm{d}^4 x \sqrt{-g_4} \left[\frac{R_{(4)}}{2\kappa} - \frac{1}{2} \left(\partial \phi \right)^2 - V(\phi) \right]_{\pm}$$

$$V(\varphi) = -V_{\rm i} \exp\left[-\frac{4\sqrt{\pi\gamma}}{m_{_{\rm Pl}}}(\varphi - \varphi_{\rm i})\right]_{\rm f}$$

Singular...

Cyclic extension



Model listing:

Quantum gravity

Ekpyrotic & cyclic Branes

String gas cosmology Antigravity Galileon Massive gravity

Multiverse models Strings & AdS/CFT



Non relativistic quantum gravity (dW)

Horava-Lifshitz Lee-Wick & Quintom F(R), f(T), Gauss-Bonnet

Mimetic matter Non-linear electromagnetic action Spinors & torsion

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Standard Failures and bouncing solutions



PP & N. Pinto-Neto, Phys. Rev. D78, 063506 (2008)



 $\ddot{a} < 0 \ \& \ \dot{a} < 0$

- accelerated expansion (inflation) or decelerated contraction (bounce)



Standard Failures and bouncing solutions



PP & N. Pinto-Neto, Phys. Rev. D78, 063506 (2008)



 $\ddot{a} < 0 \ \& \ \dot{a} < 0$

- accelerated expansion (inflation) or decelerated contraction (bounce)





$t_{\rm ini} \rightarrow -\infty$

Standard Failures and bouncing solutions



PP & N. Pinto-Neto, Phys. Rev. D78, 063506 (2008)



 $\ddot{a} < 0 \ \& \ \dot{a} < 0$

- accelerated expansion (inflation) or decelerated contraction (bounce)



Standard Failures and bouncing solutions



PP & N. Pinto-Neto, Phys. Rev. D78, 063506 (2008)



 $\ddot{a} < 0 \& \dot{a} < 0$

- accelerated expansion (inflation) or decelerated contraction (bounce)

The problem with contraction: BKL/shear instability



Mean Hubble parameter

Ekpyrotic solution:



ultra stiff eq. of state



Cosmic no-hair for ekpyrosis

- Bianchi I-VIII curvature always negative
- Impose $\rho \geq 0$
- Ultra stiff e.o.s., i.e. $\gamma = 1 + w > 2$

All initially contracting, spatially homogeneous, orthogonal Bianchi Type I-VIII cosmologies and all Bianchi type IX universes sourced by an ultra-stiff fluid with an equation of state such that $(\gamma - 2)$ is positive definite collapse into an isotropic singularity, where the sink is a spatially flat and isotropic FRW universe.

J. Lidsey, CQG23 (2006) 3517

Anisotropic pressures and ekpyrosis

- Add stress-energy tensor anisotropies (& flat anisotropic Universe):
- $\sigma_{ab} + 3H\sigma_{ab} = \pi_{ab}$
- Energy density anisotropies $\sigma^2 \sim \sigma_{ab} \sigma^{ab}$ can grow faster than V^{-2} Ekpyrotic fluid must grow faster... fine-tuned w.r.t. the anisotropic stress
- Expansion + 3-curvature anisotropy (Mixmaster) doesn't re-expand

Mixmaster



- Curvature + expansion anisotropies
- Chaotic oscillations infinite in number on approach to singularity
- Chaotic behaviour = attractor
- Anisotropy growth --> collapse



Isotropisation by shear viscosity

 $\sigma_{ab} + 3H\sigma_{ab} = \pi_{ab} \qquad \&$

 $\pi_{ab} = -\kappa \rho^{1/2} \sigma_{ab}$

C. Ganguly & M. Bruni, PRL123 (2019) 201301

Cosmic no-hair reloaded

- Bianchi I-VIII curvature always negative
- impose $\rho \geq 0$
- Matter obeys Strong Energy Condition, i.e. $\rho + 3p > 0$
- Viscosity coefficient κ obeys $\kappa > 3(2 \gamma)$

All initially contracting, spatially homogeneous, orthogonal Bianchi Type I-VIII cosmologies and all Bianchi type IX universes sourced by fluid obeying the Strong Energy Condition and with a shear viscous term such that the coefficient of viscosity κ obeys the condition that $\kappa > 3(2 - \gamma)$ collapse into an isotropic singularity, where the sink is a spatially flat and isotropic FRW universe.

C. Ganguly, arXiv:2008.02286

<u>Quantum cosmology</u>

 $K_{ij} = -\nabla_j^{(h)} n_i = \frac{1}{2N} \left(\nabla_j^{(h)} N_i + \nabla_i^{(h)} N_j - \frac{\partial h_{ij}}{\partial t} \right)$

Action (Einstein-Hilbert, compact space):

$$S = \frac{1}{16\pi G_{\rm N}} \left[\int_{\mathcal{M}} \sqrt{-g} \left(R - 2\Lambda \right) \mathrm{d}^4 x \right]$$

$$\longrightarrow S = \int L dt = \frac{1}{16\pi G_{\rm N}} \int dt \left[\int d^3 x \, N \, v \right]$$

Canonical momenta

 $x + 2 \int_{\partial \mathcal{M}} \sqrt{hK^{i}}_{i} \mathrm{d}^{3}x + \mathcal{S}_{\mathrm{matter}} \left[\Phi\left(x\right)\right]$

$\sqrt{h}\left(K_{ij}K^{ij} - K^2 + {}^3R - 2\Lambda\right) + L_{\text{matter}}$

 $\iota^{ij}K$

aints

Hamiltonian

$$H \equiv \int d^3x \left(\pi^0 \dot{N} + \pi^i \dot{N}_i + \pi^{ij} \dot{N}_i \right)$$
$$= \int d^3x \left(\pi^0 \dot{N} + \pi^i \dot{N}_i + N \mathcal{H}_i \right)$$
$$\mathcal{H} = -\frac{1}{\sqrt{2}}$$

variation wrt lapse: $\mathcal{H} = 0 \rightarrow$ Hamiltonian constraint variation wrt shift: $\mathcal{H}^i = 0 \rightarrow$ momentum constraint

 $\dot{h}_{ij} + \pi^{\Phi} \dot{\Phi} - L$ $\mathcal{H} + N_i \mathcal{H}^i$ $\frac{1}{\sqrt{h}} \left(h_{ik} h_{jl} - \frac{1}{2} h_{ij} h_{kl} \right) \pi^{ij} \pi^{kl} - \sqrt{h} R$ $\mathcal{H}^{i} = -2\sqrt{h}\nabla_{j}\left(\frac{\pi^{ij}}{\sqrt{h}}\right)$

\implies classical description complete

Superspace & canonical quantization relevant configuration space Riem $(\Sigma) \equiv \{h_{ij}(x^{\mu}), \Phi(x^{\mu}) | x \in \Sigma\}$ $GR \Longrightarrow invariance/diffeomorphisms \Longrightarrow Conf = \frac{Riem(\Sigma)}{Diff(\Sigma)}$: superspace Wave functional $\Psi[h_{ij}(x), \Phi(x)] = \langle h_{ij}, \Phi | \Psi \rangle$ + Dirac canonical quantization procedure

$$\pi^{ij} \to -i \frac{\delta}{\delta h_{ij}} \qquad \pi^{\Phi} \to -i \frac{\delta}{\delta \Phi} \qquad \pi^{0}$$

parameters

primary constraints
$$\begin{cases} \hat{\pi}^0 = -i\frac{\delta\Psi}{\delta N} = 0\\ \hat{\pi}^i = -i\frac{\delta\Psi}{\delta N_i} = 0 \end{cases}$$

momentum $\hat{\mathcal{H}}^{i}\Psi = 0 \implies i\nabla_{j}^{(h)}\left(\frac{\delta\Psi}{\delta h_{ij}}\right) = 8\pi G_{N}\hat{T}^{0i}\Psi$ same Ψ for configurations related by a coordinate transformation

Hamiltonian

$$\begin{aligned} \hat{\mathcal{H}}\Psi &= \begin{bmatrix} -16\pi G_{\mathrm{N}} \mathcal{G}_{ijkl} \frac{\delta^{2}}{\delta h_{ij} \delta h_{kl}} + \frac{\sqrt{h}}{16\pi G_{\mathrm{N}}} \left(-^{3}R + 2\Lambda + 16\pi G_{\mathrm{N}} \hat{T}^{00} \right) \end{bmatrix} \Psi = 0 \\ \mathcal{G}_{ijkl} &= \frac{1}{2\sqrt{h}} \left(h_{ik} h_{jl} + h_{il} h_{jk} - h_{ij} h_{kl} \right) \\ \text{De Witt metric} \end{aligned}$$

$$\begin{aligned} & \mathcal{W} heeler - De \text{ Witt equation} \end{aligned}$$

primary constraints
$$\begin{cases} \hat{\pi}^0 = -i \frac{\delta \Psi}{\delta N} = 0\\ \hat{\pi}^i = -i \frac{\delta \Psi}{\delta N_i} = 0 \end{cases}$$

momentum
$$\hat{\mathcal{H}}^{i}\Psi = 0 \implies i\nabla_{j}^{(h)} \left(\frac{\partial\Psi}{\delta h_{i}}\right)$$

same Ψ for configuration

Hamiltonian

 $\hat{\mathcal{H}}\Psi = 0$

$\left(\frac{\delta\Psi}{\delta h_{ij}}\right) = 8\pi G_{\rm N} \hat{T}^{0i} \Psi$

ons related by a coordinate transformation

TIMELESS Schrödinger equation

mini-superspace

restrict attention from an infinite dimensional configuration space to a 2 dimensional space *= mini-superspace*

$$h_{ij} \mathrm{d} x^i \mathrm{d} x^j \quad \mapsto \quad a^2(t) \left[\frac{\mathrm{d} r^2}{1 - \mathcal{K} r^2} + r^2 \left(\mathrm{d} \theta^2 + \sin^2 \theta \mathrm{d} \phi^2 \right) \right]$$

WDW equation becomes Schrödinger like for $\Psi[a(t), \phi(t)]$

Conceptual & technical issues

infinite # d.o.f. to a few: mathematical consistency? freeze momenta... Heisenberg uncertainties? [quantization, minisuperspace] $\neq 0$

The clock issue in quantum cosmology

GR = constrained system: lack of external time

arbitrary degree of freedom: internal clock

arbitrary non vanishing lapse function $d\tau = N d\tau' \implies \frac{d}{d\tau'} \mathcal{O}(q_i, p_i) = \{\mathcal{O}, NC\}_{P.B}$ hamiltonian H = NC

Constrained system $C(\lbrace q^k \rbrace, \lbrace p_k \rbrace) = H_{tot}(\lbrace q^k \rbrace, \lbrace p_k \rbrace) = 0$ Canonical transformation $(\{q^k\}, \{p_k\}) \mapsto (\{Q^a\}, \{P_a\})$ Quantum system $\hat{H}_{tot}\Psi \equiv \hat{C}\Psi(Q^a) = 0$ $\hat{C} \text{ linear in } \hat{P}_{\alpha} = -i \frac{\partial}{\partial Q^{\alpha}}$

$$\hat{C} = \hat{P}_{\alpha} + \hat{H}(P_{1}, \cdot)$$
$$\hat{C}\Psi(Q^{a}) = 0 \Longrightarrow t$$

$\cdots, P_{\alpha-1}, P_{\alpha+1}, \cdots, P_n, \{Q^a\})$

time-dependent Schrödinger equation

A simple example

 $\mathcal{S} = \frac{1}{2} \int \left(\frac{\dot{x}^2}{z} \right)$

First, redefine time: $d\tau = zdt$ \longrightarrow S =

Classical EOMs

$$\begin{cases} \frac{\mathrm{d}^2 x}{\mathrm{d}\tau^2} = -x \\ \frac{\mathrm{d}^2 y}{\mathrm{d}\tau^2} = -y \end{cases} \longrightarrow 2 \mathrm{i}$$

$$-zx^2 - \frac{\dot{y}^2}{z} + zy^2 \right) \mathrm{d}t$$

$$\frac{1}{2}\int \left[\left(\frac{\mathrm{d}x}{\mathrm{d}\tau}\right)^2 - x^2 - \left(\frac{\mathrm{d}y}{\tau}\right)^2 + y^2\right]\mathrm{d}\tau$$

independent harmonic oscillators $H_{tot} = H + H_y$

 $H_{\rm tot} = P_T + H$ on shell $H_{\rm tot} \approx 0$

Quantization: x only! $\hat{H}_{tot}\psi(x,T) = 0 \implies i\frac{\partial\psi}{\partial T} = \hat{H}\psi(x,T)$

 $\& \int |\psi(x,T)|^2 \mathrm{d}x = 1$ y remains classical (clock)

Scale factors $\begin{cases} a_1 = e^{\beta_0 + \beta_+ + \sqrt{3}\beta_-} \\ a_2 = e^{\beta_0 + \beta_+ - \sqrt{3}\beta_-} \\ a_3 = e^{\beta_0 - 2\beta_+} \end{cases}$

Action
$$S = \int d\tau \Big(\underbrace{p_0 \dot{\beta}_0 + p_+ \dot{\beta}_+ + p_- \dot{\beta}_-}_{d\theta} - I \Big)$$

 $\frac{d\theta}{d\tau}$ canonical one-form
constraint $C = \frac{e^{-3\beta_0}}{24} \Big(-p_0^2 + p_+^2 \Big)$

Volume $V \equiv a_1 a_2 a_3 = e^{3\beta_0}$ $\dot{\beta}_{-} - NC$ $d\beta_{0} = \frac{1}{3} e^{-3\beta_{0}} dV$ nical

 $+ p_{\perp}^2 + p_{\perp}^2$)

ensure canonical one-form remains canonical $p_V \equiv \frac{e^{-3p_0}}{2}p_0$ $\mathrm{d}\theta = p_V \mathrm{d}V + p_+ \mathrm{d}\beta_+ + p_- \mathrm{d}\beta_-$

 $C = \frac{3V}{8} \left(-p_V^2 + \frac{p_+^2 + p_+^2}{9V} \right)$ constraint

cyclic variable $\dot{p}_{+} = 0$ set $p_{+} = k \cos \alpha$ and $p_{-} = k \sin \alpha$ $d\theta = p_V dV + p_k dk + p_\alpha d\alpha + d \left(k \cos \alpha \beta_+ + k \sin \alpha \beta_- \right)$

 $p_{k} \equiv -(\cos \alpha \beta_{+} + \sin \alpha \beta_{-}),$ $p_{\alpha} \equiv (k \sin \alpha \beta_{+} - k \cos \alpha \beta_{-})$

$$\left(\frac{p_{-}^{2}}{2}\right)$$

Hamilton equations

$$\dot{k} = 0$$

$$\dot{p}_k = -N\frac{k}{12V}$$

$$\dot{V} = -N \frac{3V p_V}{4}$$
$$\dot{p}_V = -N \left[\frac{3}{8} \left(-p_V^2 + \frac{k^2}{9V^2} \right) - \frac{k^2}{12V^2} \right]$$

closed for V and p_V

+ constraint $\frac{3V}{8}\left(-p_V^2 + \frac{k^2}{9V^2}\right) = 0$

Solving directly in the action $S = \int d\theta = \int d\tau \left(p_V \dot{V} - \frac{V^2 p_V^2}{2} \right)$

classical unconstrained one dimensional system $\frac{\mathrm{d}}{\mathrm{d}\tau} \left(V p_V \right) = 0 \qquad \Longrightarrow \qquad V p_V = V_0 p_{V0}$

$$V = V_0 \mathrm{e}^{(V p_V) \tau} \qquad \text{and} \qquad$$

<u>Choosing a time</u> $\frac{d}{d\tau} \left(9\frac{p_k}{k}\right) = -\frac{3}{4}\frac{N}{V}$ monotonically increasing function

$$p_V = p_{V0} \mathrm{e}^{-(V p_V) \cdot \tau}$$

symmetric ordering choice

coordinate transformation $V \mapsto Z = \ln V$

 $U\hat{H}U^{-1} = -\partial_Z^2$, and $Z \in \mathbb{R}$

 $H = V^2 p_V^2 \quad \mapsto \quad \hat{H} = \sqrt{V} \frac{1}{i} \partial_V \sqrt{V} \cdot \sqrt{V} \frac{1}{i} \partial_V \sqrt{V}$

slow-gauge time

$$\mathcal{S} = \int \mathrm{d}\theta = \int \mathrm{d}\eta \left(V \right)$$

new time variable

$$\eta \equiv \frac{9p_k}{k} + \frac{V - \ln V}{Vp_V}$$

Quantization: a gaussian wave packet

$$u(V,\eta) = \frac{e^{-k^2/4}}{\sqrt{1+4i\eta}} \exp\left[-\frac{(V-ik/2)^2}{1+4i\eta}\right]$$

implement boundary conditions to ensure self-adjointness

$$\psi(V,\eta) = \frac{u(V+V_0,\eta) - u(-V+V_0,\eta)}{\left[\sqrt{\pi/2} \left(1 - e^{-V_0^2 - k^2/2}\right)\right]^{1/2}}$$

— solves the Schrödinger equation

$$i\frac{\partial}{\partial\eta}\psi = -$$

Operator ordering ambiguity $\pi_V^2 \mapsto \hat{V}^s \hat{\pi}_V \hat{V}^{-2s} \hat{\pi}_V \hat{V}^s$

self-adjoint hamiltonian on the half-line s > 3/4

V

Heisenberg equations of motion

> $\frac{\mathrm{d}}{\mathrm{d}\eta}\hat{V}^2 = -i[\hat{V}^2, \hat{H}] = 4\hat{D}$ $\frac{\mathrm{d}}{\mathrm{d}n}\hat{D} = -i[\hat{D},\hat{H}] = 2\hat{H}$

Closed algebra of operators $\begin{cases} [\hat{V}^2, \hat{H}] &= 4iD, \\ [\hat{D}, \hat{H}] &= 2i\hat{H}, \\ [\hat{V}^2, \hat{D}] &= 2i\hat{V}^2. \end{cases}$

 $\hat{D} \equiv \frac{1}{2} \left(\hat{V} \hat{\pi}_V + \hat{\pi}_V \hat{V} \right)$

Heisenberg equations of motion

$$\frac{\mathrm{d}}{\mathrm{d}\eta}\hat{V}^2 = -i[\hat{V}^2, \hat{H}]$$
$$\frac{\mathrm{d}}{\mathrm{d}\eta}\hat{D} = -i[\hat{D}, \hat{H}]$$

solution as time-dependent operators

$$\hat{D}(\eta) = 2\hat{H}\eta + \hat{D}(0)$$

expectation values follows similar equations...

-> semi-classical variables

$$\check{V}(t) = \sqrt{\langle \hat{V}^2(t) \rangle}$$

$$\check{\pi}_V(t) = \frac{\langle \hat{D}(t) \rangle}{\check{V}(t)}$$

$[] = 4\hat{D}$

$\hat{V}^2 = 4\hat{\eta}^2 + 4\hat{D}(0)\eta + \hat{V}^2(0)$

phase space solution

 η

Changing the time variable $\eta' = \eta'(\eta, V, \pi_V)$

redefining the dynamical variables in the process

change the canonical one-form

$$d\theta = \pi_V dV - \frac{\pi_V^2}{2} d\eta = \pi'_V dV' - \frac{\pi_V'^2}{2} d\eta' + d\left[(\eta - \eta') \frac{\pi_V'^2}{2} \right]$$

→ <u>same system!</u>

delay function $\Delta(V, \pi_V) = \eta' - \eta$ no dependency on time

$\pi'_V = \pi_V$ & $V' = V + \pi_V(\eta' - \eta)$ no change of range...

Delay function $\Delta = V e^{-2|\pi_V|/3} \sin(3V\pi_V)/(10\pi_V)$

Delay function $\Delta = V(\pi_V - 10^{-0.2}\pi_V^3 + \pi_V^5/10)$

Delay function $\Delta = 10^{-0.5} V \sin(2\pi_V) / \pi_V$

Delay function $\Delta = 10^{-0.5} (V+1) \cos(3\pi_V) / \pi_V$

Comparison between different delay functions

Same asymptotics

P. Małkiewicz, PP and S. Vitenti, *Phys. Rev.* **D101**, 046012 (2020) [arXiv:1911.09892]

- Bouncing alternative to inflation still alive

— Shear issue: ekpyrosis or shear viscosity

- Classical / quantum bounce

 Observational consequences & perturbations (talks by Shiv Sethi, Yi Fu Cai...)

Conclusions