



The bouncing scenario

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Framework

Homogeneous & isotropic metric (FLRW): $ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - \mathcal{K}r^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$

Hubble rate $H \equiv \frac{\dot{a}}{a}$

spatial curvature

Matter component: perfect fluid $T_{\mu\nu} = pg_{\mu\nu} + (\rho + p) u_\mu u_\nu$

equation of state

$$p = w\rho \rightarrow \begin{cases} w = 0 & \text{dust} \\ w = \frac{1}{3} & \text{radiation} \end{cases}$$

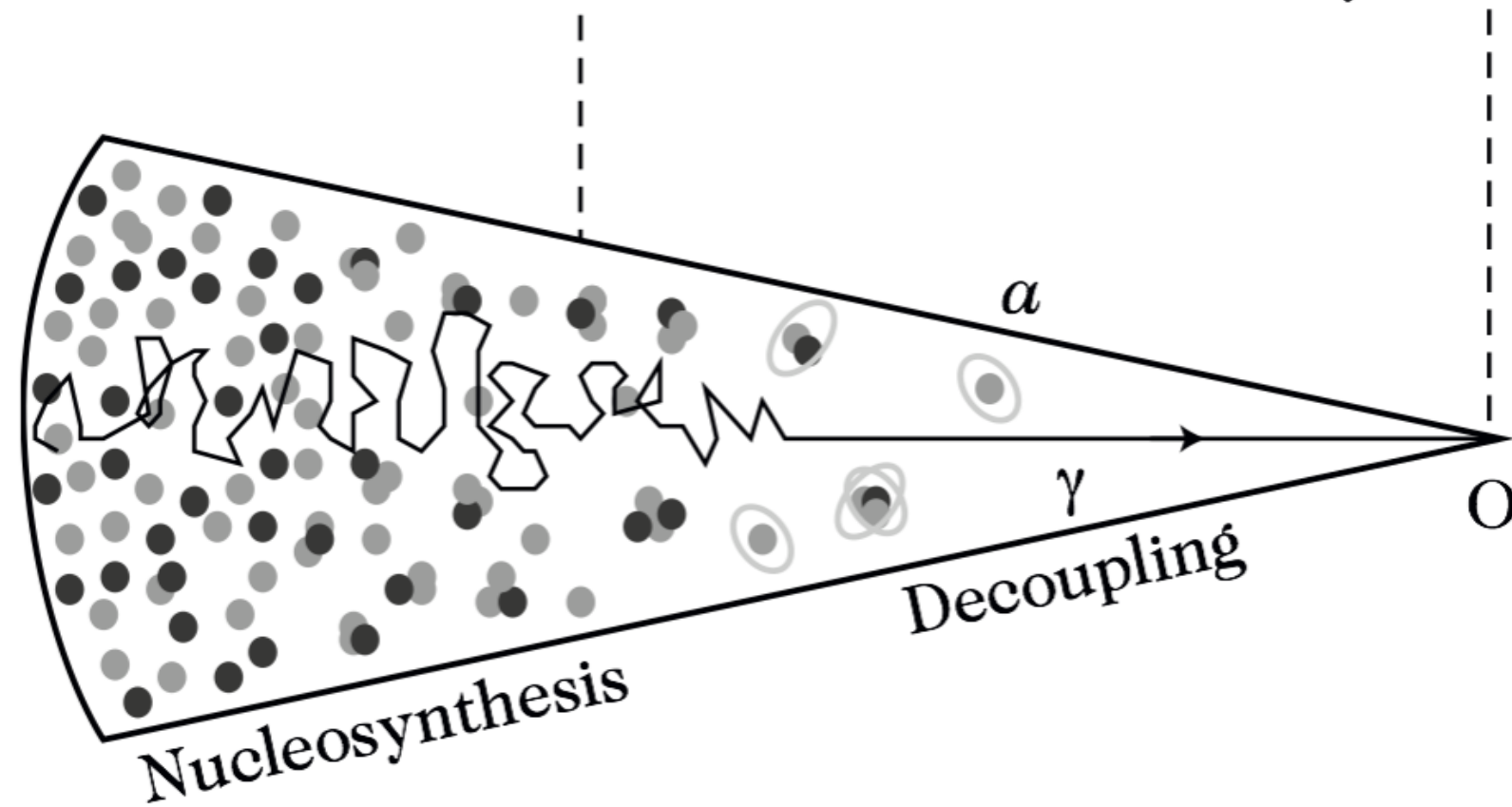
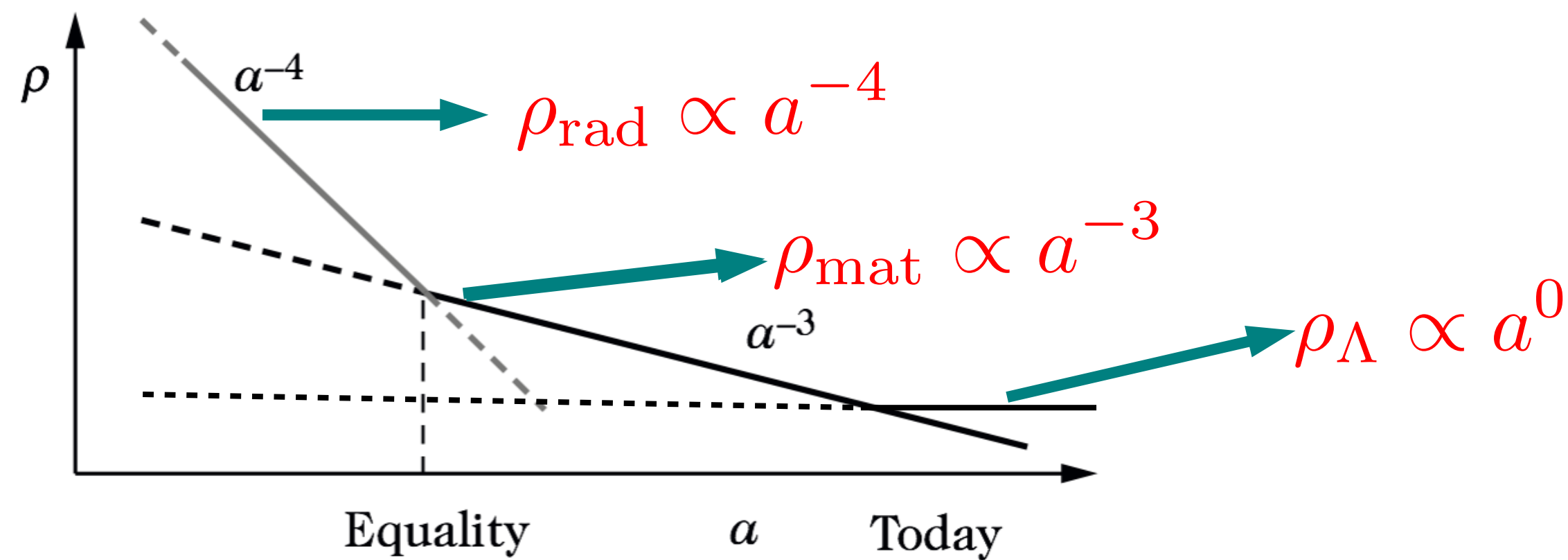
+ cosmological constant = Einstein equations

$$\begin{cases} H^2 + \frac{\mathcal{K}}{a^2} = \frac{1}{3} (8\pi G_N \rho + \Lambda) \\ \frac{\ddot{a}}{a} = \frac{1}{3} [\Lambda - 4\pi G_N (\rho + 3p)] \end{cases}$$

Particular solution: dust and radiation

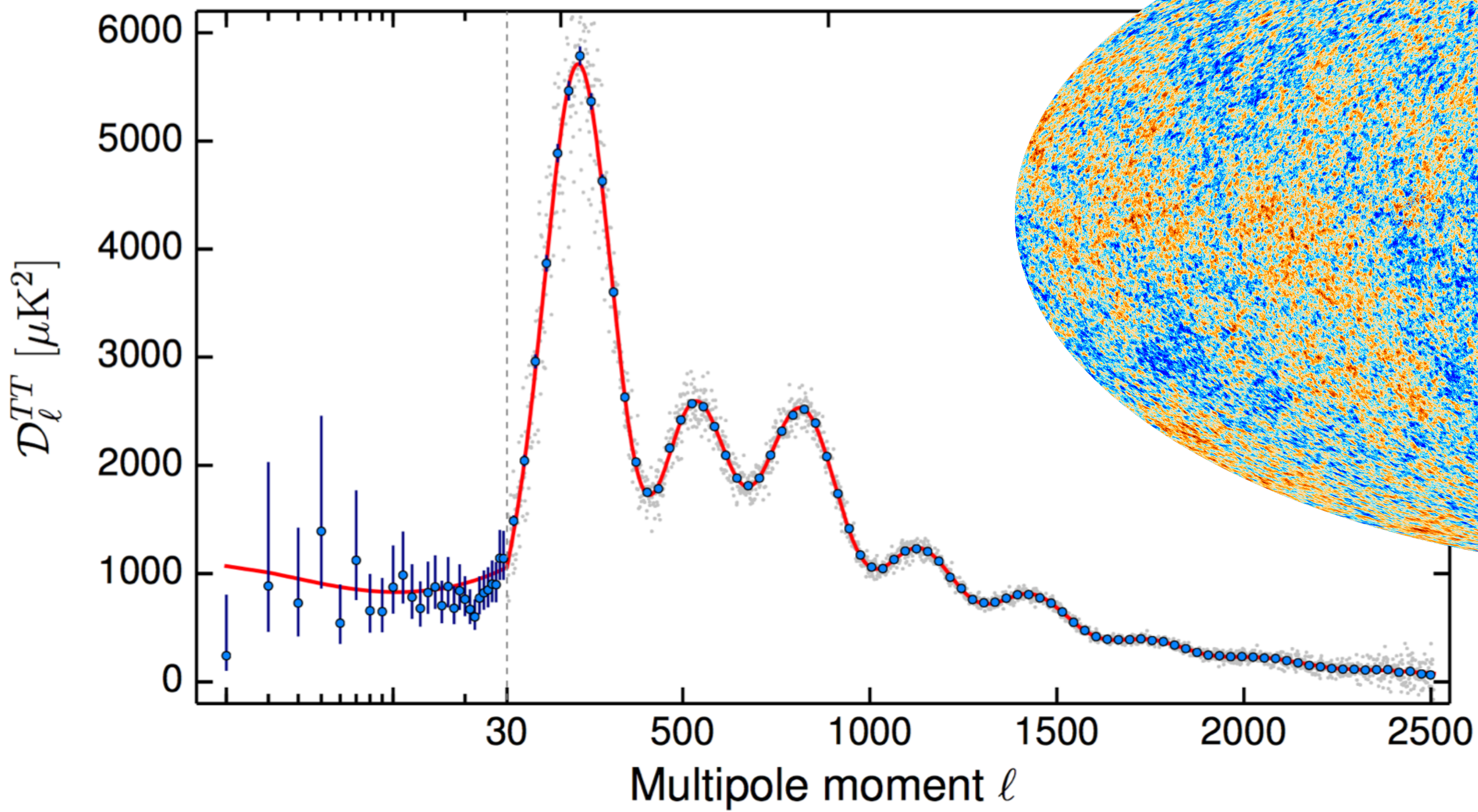
integrate conservation equation

$$\rho[a(t)] = \rho_{\text{ini}} \exp \left\{ -3 \int [1 + w(a)] d \ln a \right\} \stackrel{w \rightarrow \text{cst}}{=} \rho_{\text{ini}} \left(\frac{a}{a_{\text{ini}}} \right)^{-3(1+w)}$$



Phenomenologically valid description for 14 Gyrs!!!!

Planck 2015



$\Omega_{\kappa} = 0.000 \pm 0.005$

$n_s = 0.9639 \pm 0.0047$ almost scale invariant

$f_{NL}^{loc} = 0.8 \pm 5$
 $f_{NL}^{eq} = -4 \pm 43$
 $f_{NL}^{ort} = -26 \pm 21$

excluded
 gaussian signal

isocurvature $\lesssim 1\%$

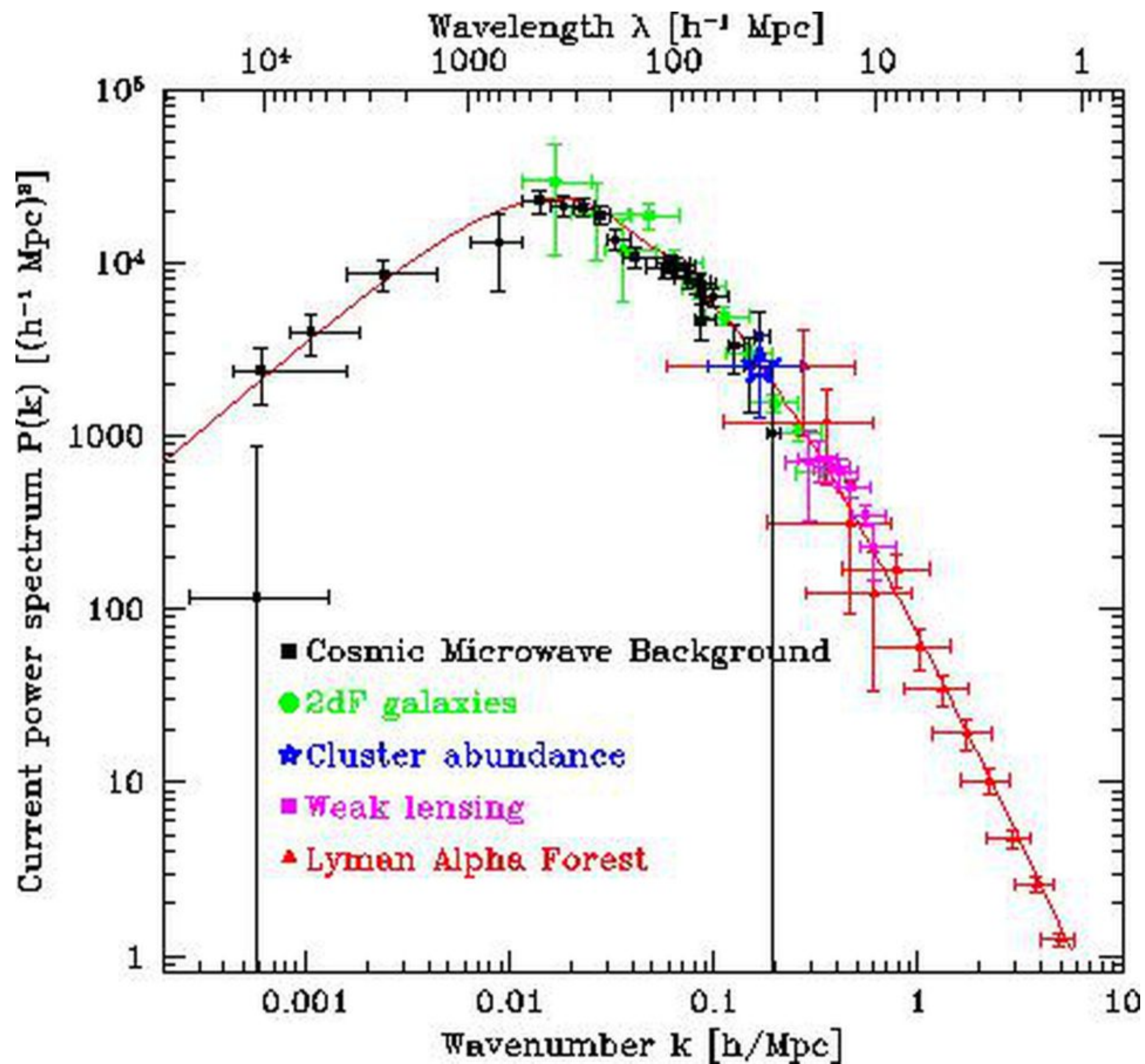
$r < 0.08$

quantum vacuum fluctuations of a single scalar d.o.f



compatible with
INFLATION

Numerical simulation for large scale structure formation...



Standard Failures and inflationary solutions

Singularity

Not solved... actually not addressed!

Horizon

$$d_H \equiv a(t) \int_{t_i}^t \frac{d\tau}{a(\tau)}$$

can be made as big as one wishes

Flatness

$$\frac{d}{dt} |\Omega - 1| = -2 \frac{\ddot{a}}{\dot{a}^3}$$

$$\ddot{a} > 0 \quad \& \quad \dot{a} > 0$$

accelerated expansion (**inflation**)

Homogeneity & Isotropy

Initial Universe = very small patch

Accelerated expansion drives the shear to zero...



vacuum state!

+ attractor

Perturbations

Bonus of the theory: predictions!!!

Others

dark matter/energy, baryogenesis, ...

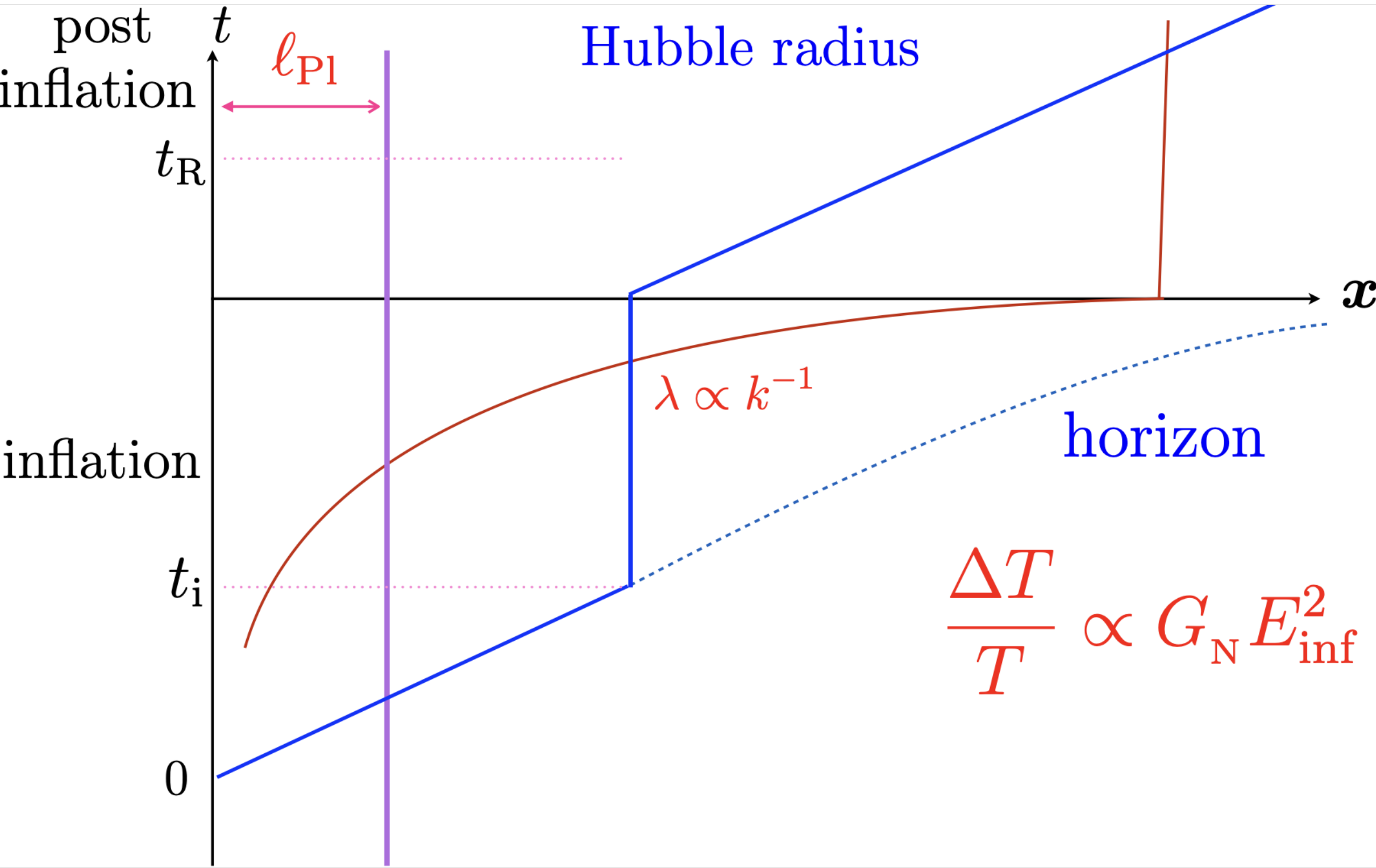
Inflation:

- ☺ **solves cosmological puzzles**
- ☺ **uses GR + scalar fields [(semi-)classical]**
- ☺ **can be implemented in high energy theories**
- ☺ **string implementation (brane inflation, ...)**

- ☺ **makes falsifiable predictions ...**
- ☺ **... consistent with all known observations**

why bother with alternatives?

From R. Brandenberger, in M. Lemoine, J. Martin & PP (Eds.), "Inflationary cosmology", Lect. Notes Phys. 738 (Springer, Berlin, 2007).



● Singularity $\exists t_{(\pm\infty)}; a(t) \rightarrow 0$

● Trans-Planckian

$$\exists t; l(t) = l_0 \frac{a(t)}{a_0} \leq l_{Pl}$$

● Hierarchy (amplitude)?

$$\frac{V(\varphi)}{\Delta\varphi^4} \leq 10^{-12}$$

● Classical GR?

$$\frac{\Delta T}{T} \propto G_N E_{inf}^2 \sim \left(\frac{E_{inf}}{M_{Pl}} \right)^2 \rightarrow E_{inf} \simeq 10^{-3} M_{Pl}$$

● η problem & Lyth bound

● Initial condition & entropy

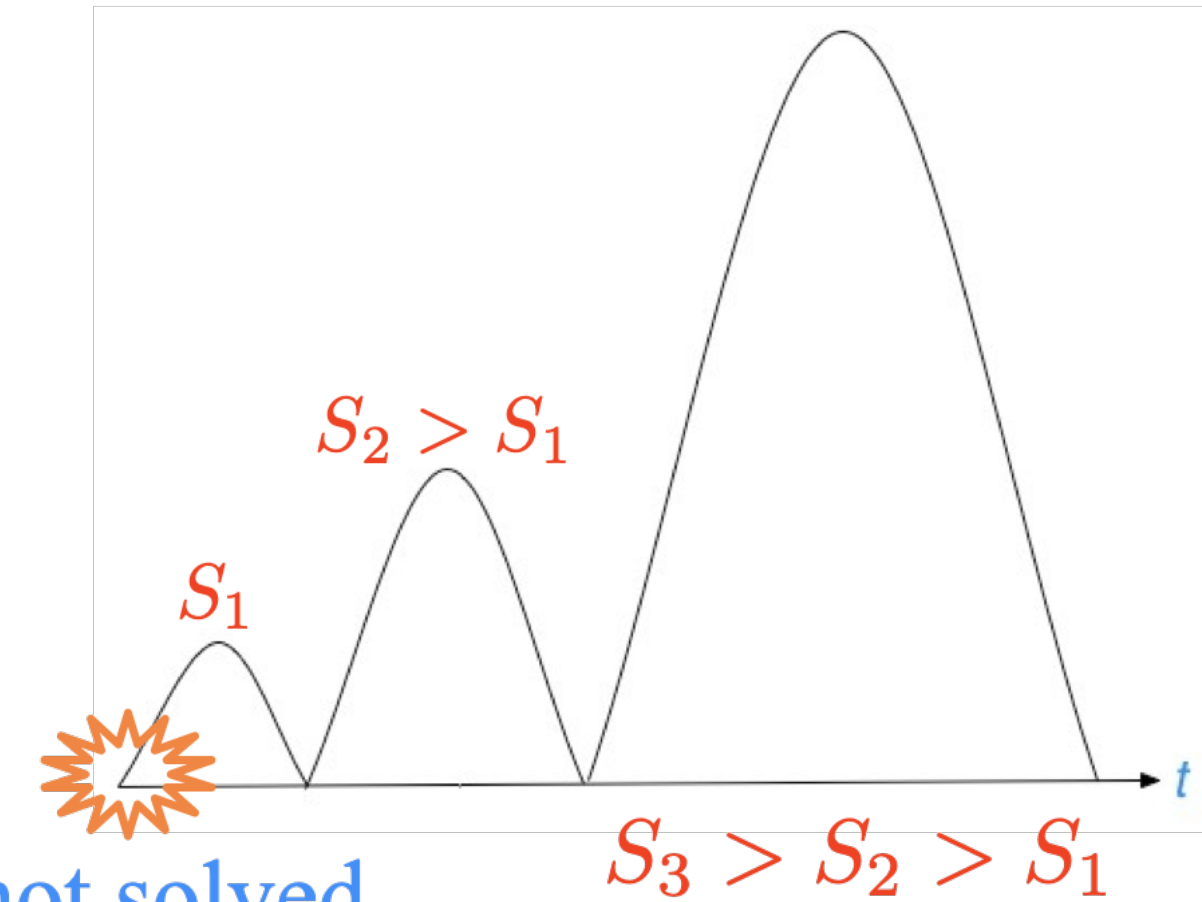
● Eternal inflation & measure (anthropic)

A brief history of bouncing cosmology

→ R. C. Tolman, “*On the Theoretical Requirements for a Periodic Behaviour of the Universe*”, PRD 38, 1758 (1931)

→ G. Lemaître, “*L’Univers en expansion*”, Ann. Soc. Sci. Bruxelles (1933)

...



...

Singularity pb not solved

→ A. A. Starobinsky, “*On one non-singular isotropic cosmological model*”, Sov. Astron. Lett. 4, 82 (1978)

→ V. N. Melnikov, S.V. Orlov, Phys. Lett. A 70, 263 (1979).

→ R. Durrer & J. Laukerman, “*The oscillating Universe: an alternative to inflation*”, Class. Quantum Grav. 13, 1069 (1996)

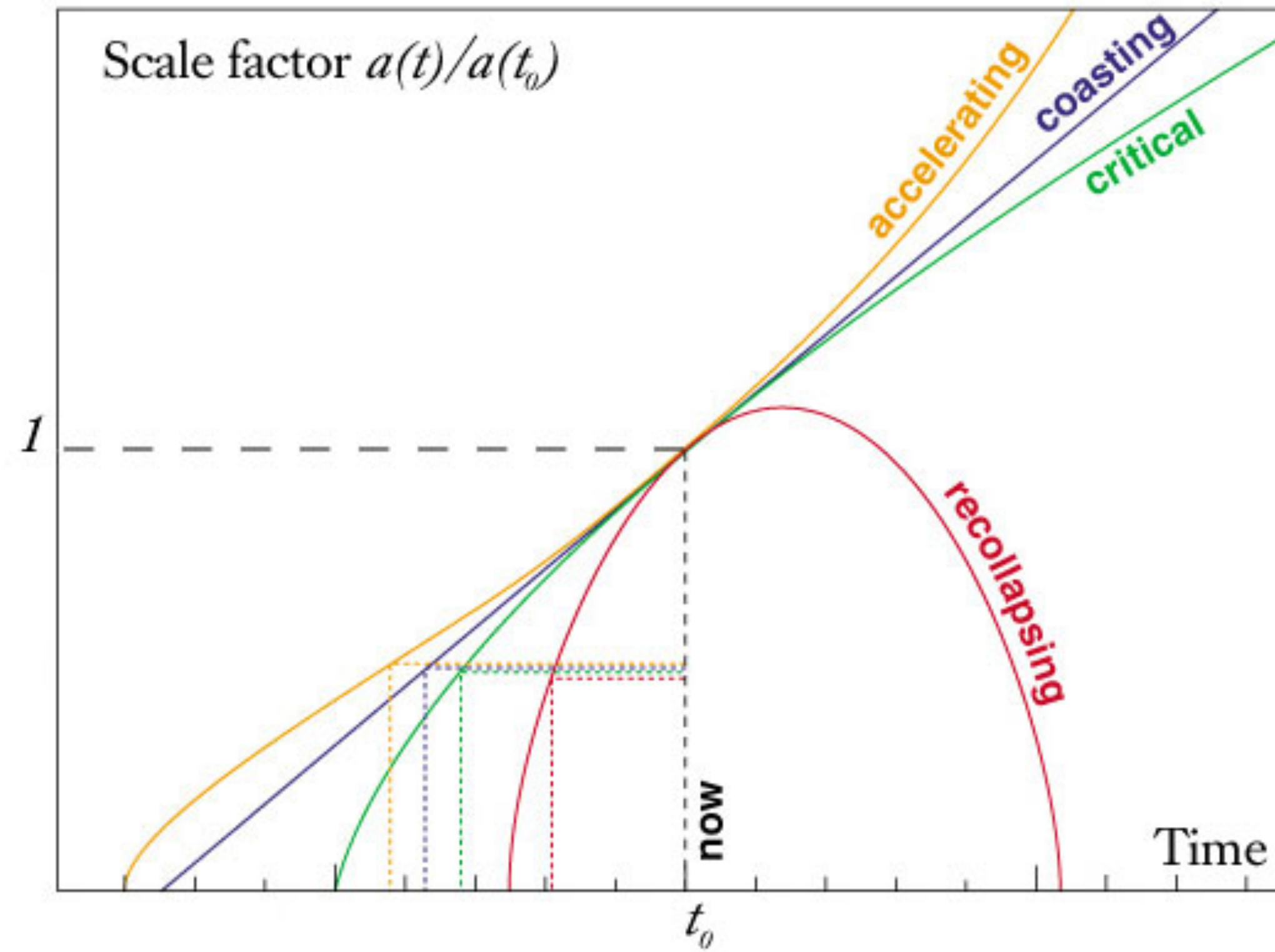
→ Many new ideas, models...

→ M. Novello & S.E. Perez Bergliaffa, “*Bouncing cosmologies*”, Phys. Rep. 463, 127 (2008)

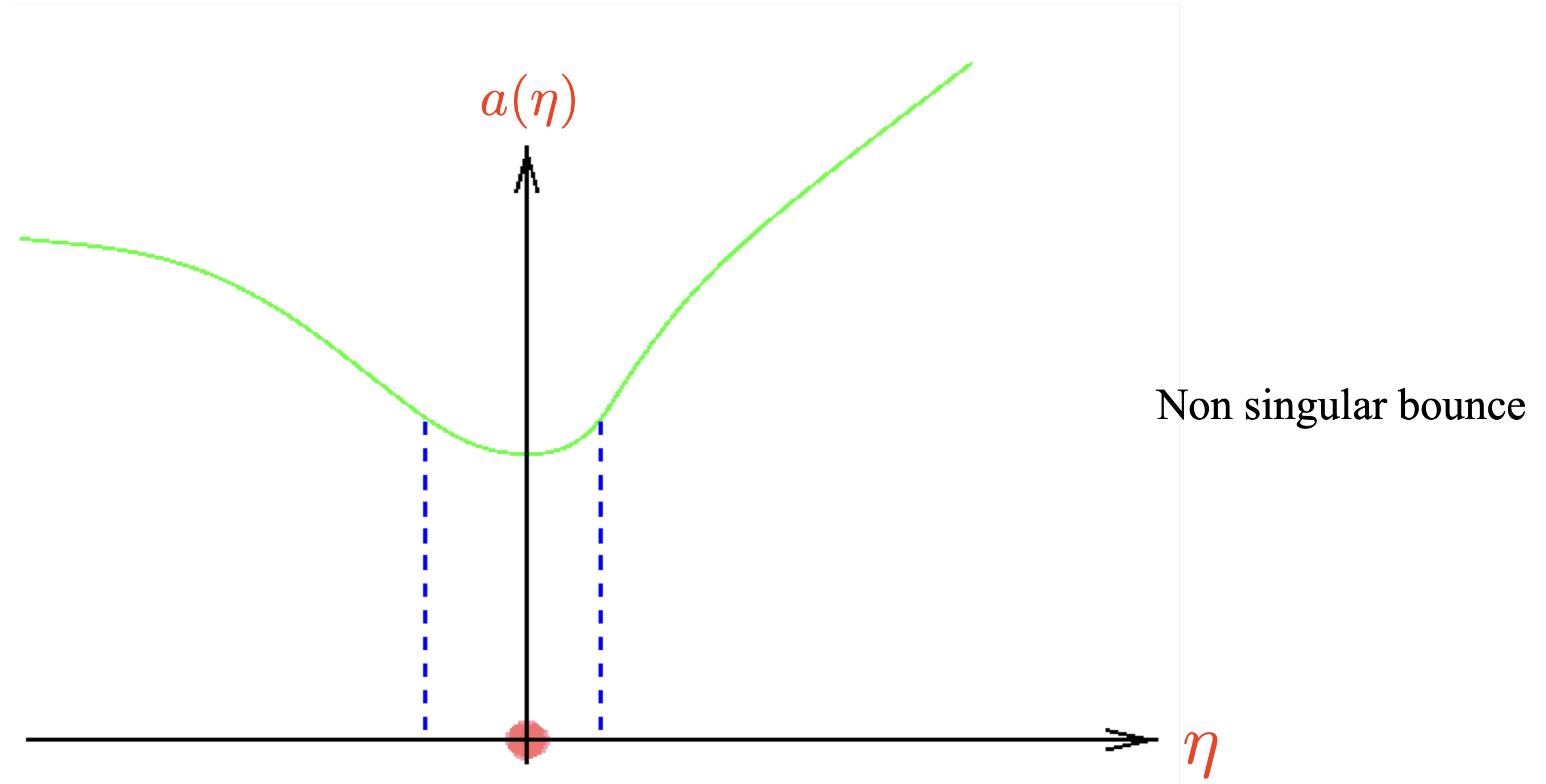
→ D. Battefeld & PP, “*A Critical Review of Classical Bouncing Cosmologies*”, Phys. Rep. 571, 1 (2015)

→ R. Brandenberger & PP, “*Bouncing cosmologies: Progress and problems*”, Found. Phys. (2017)

The issue we are interested in: the singularity



The issue we are interested in: the singularity



Model listing:

Quantum gravity

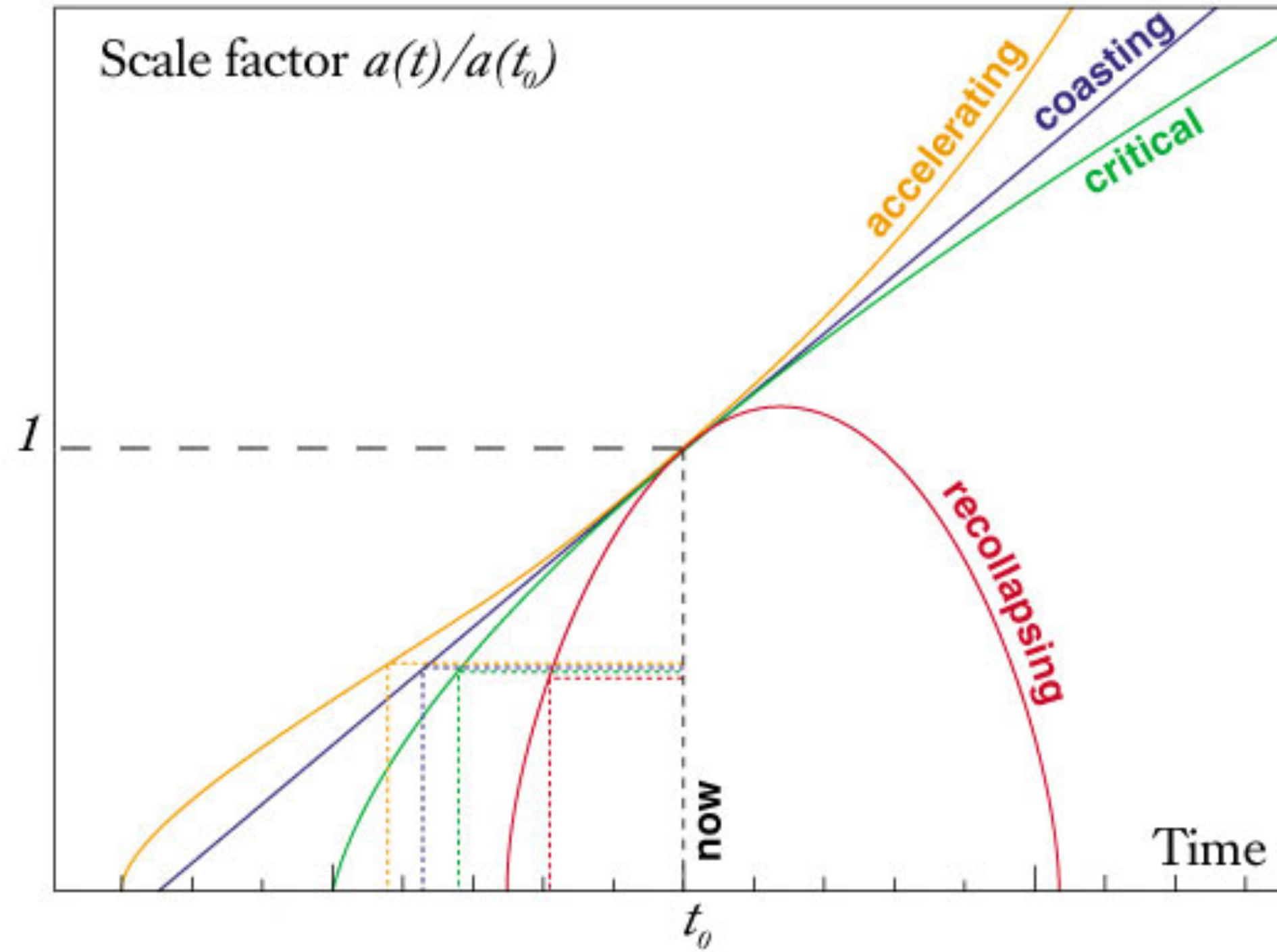
LQG & LQC

Canonical quantum gravity (WdW)

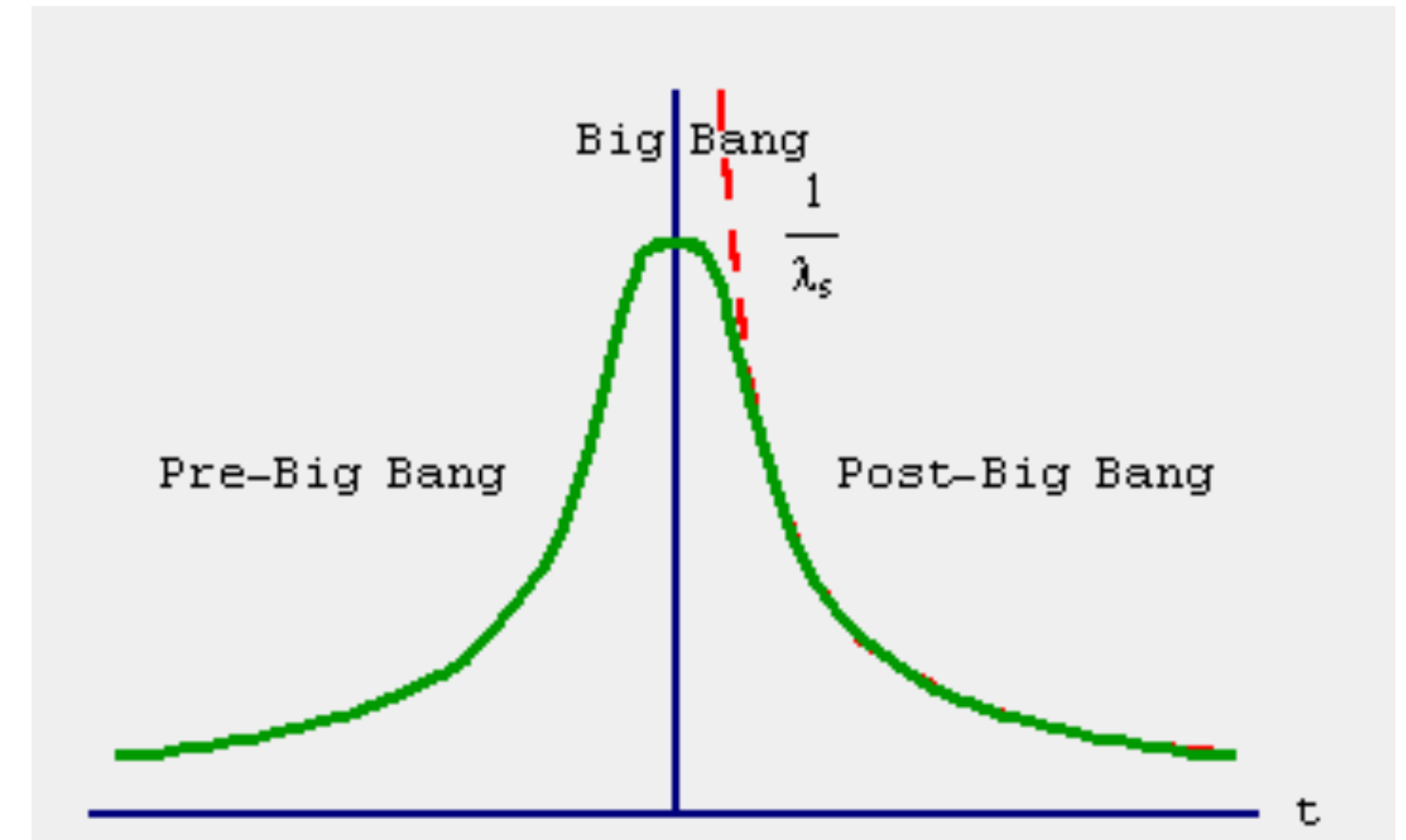
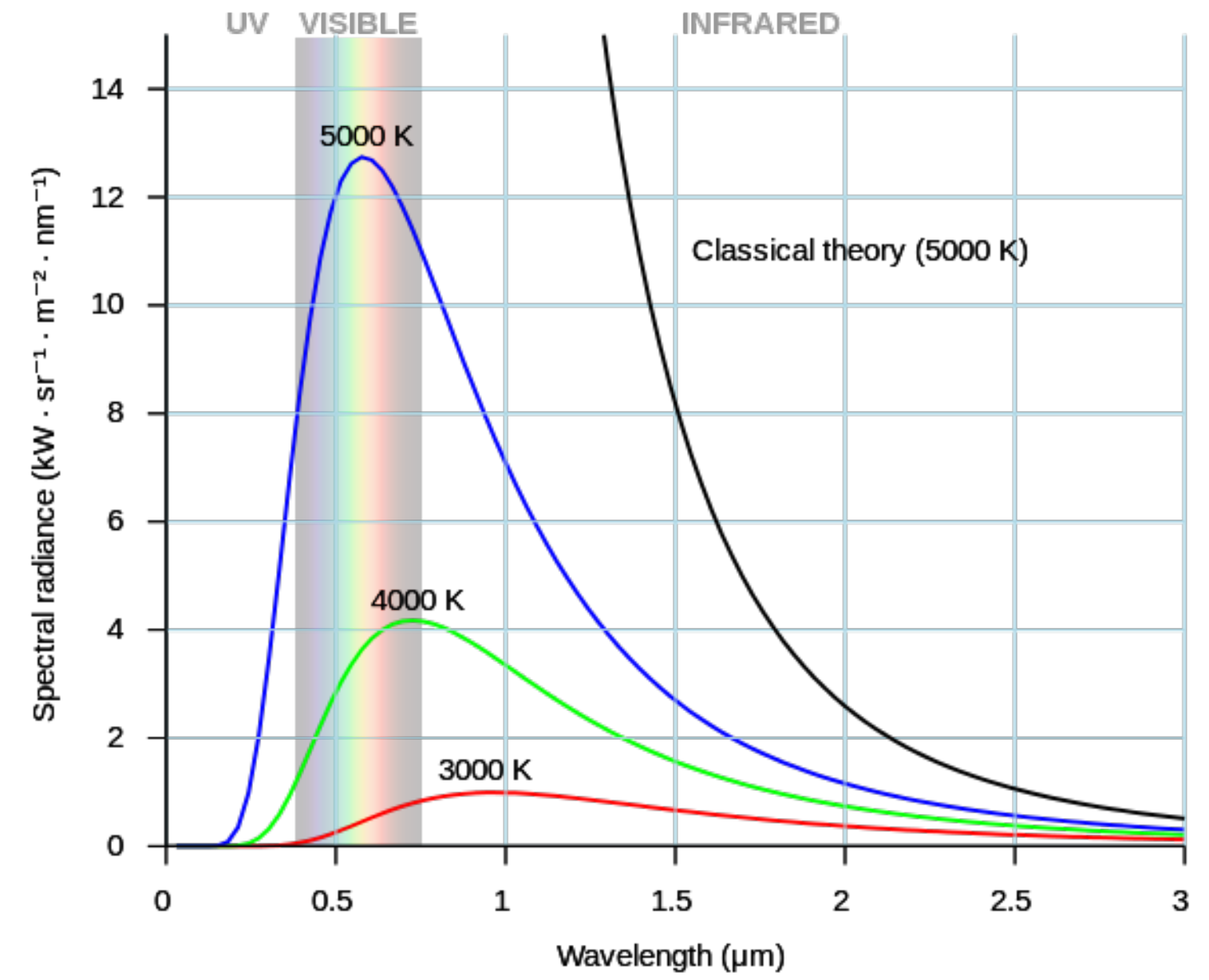
String theory

Non relativistic quantum gravity

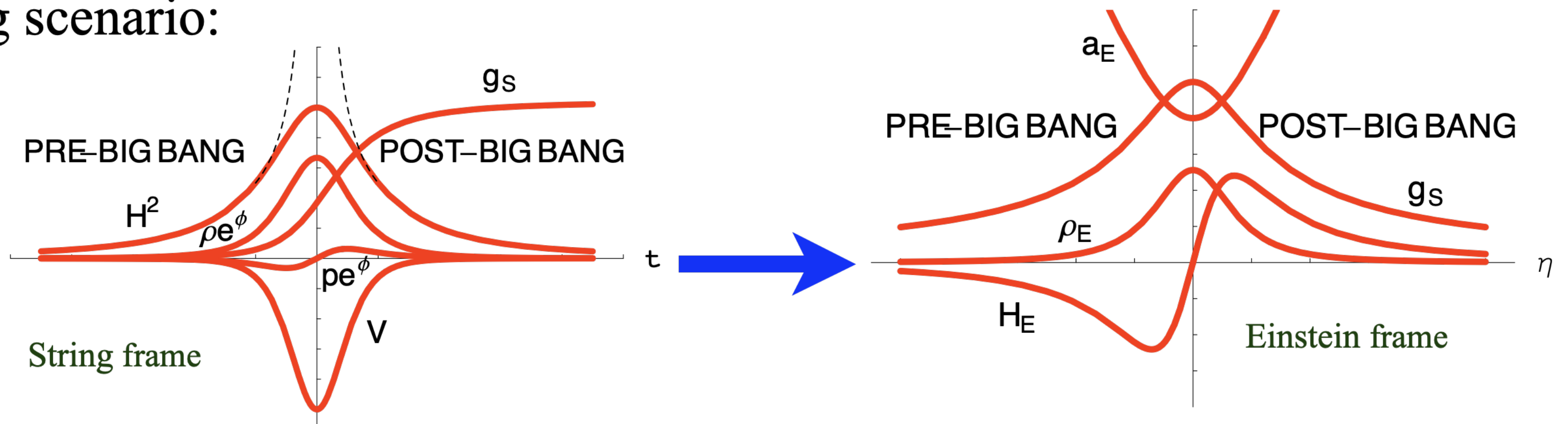
Singularity problem...



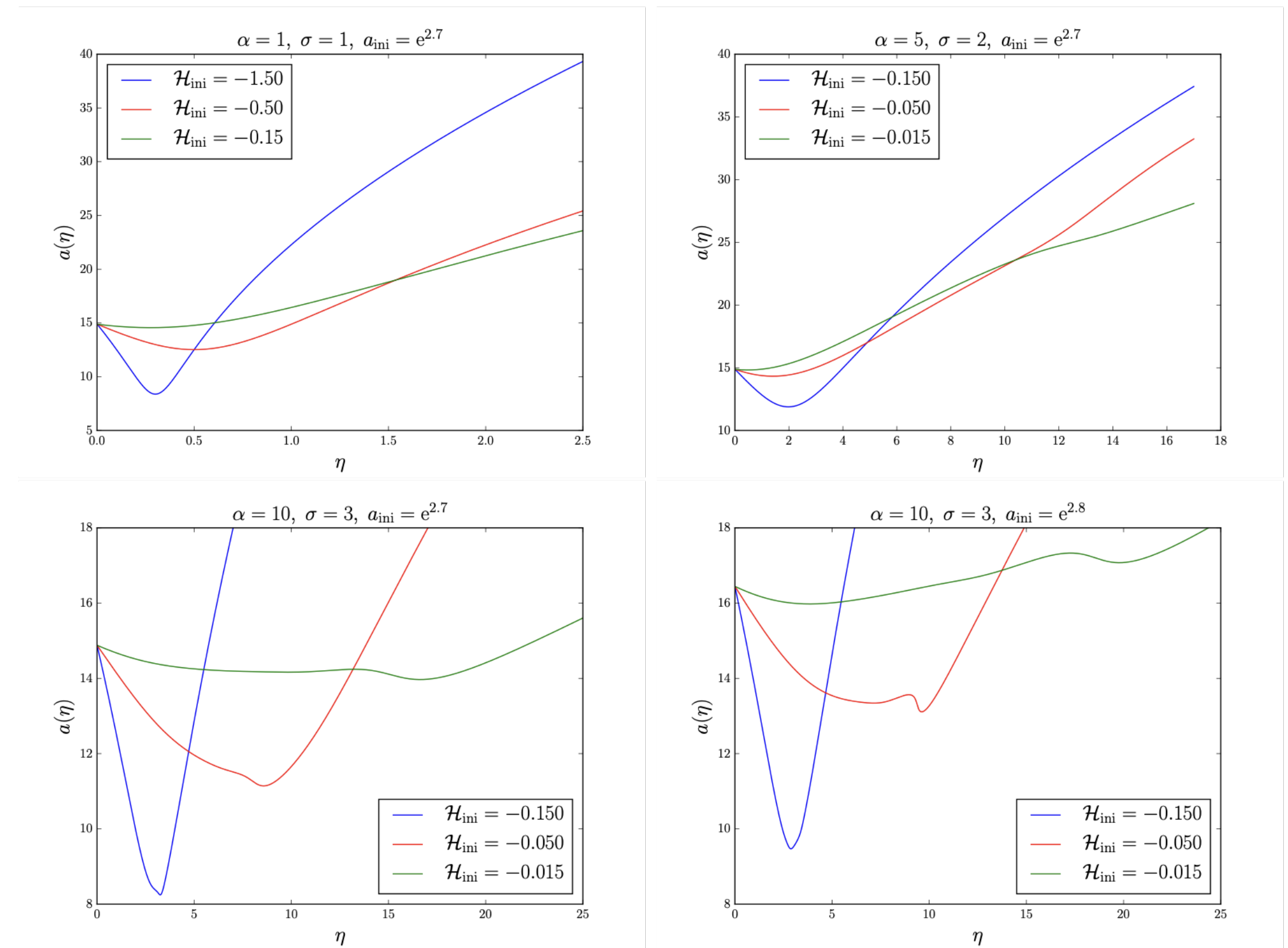
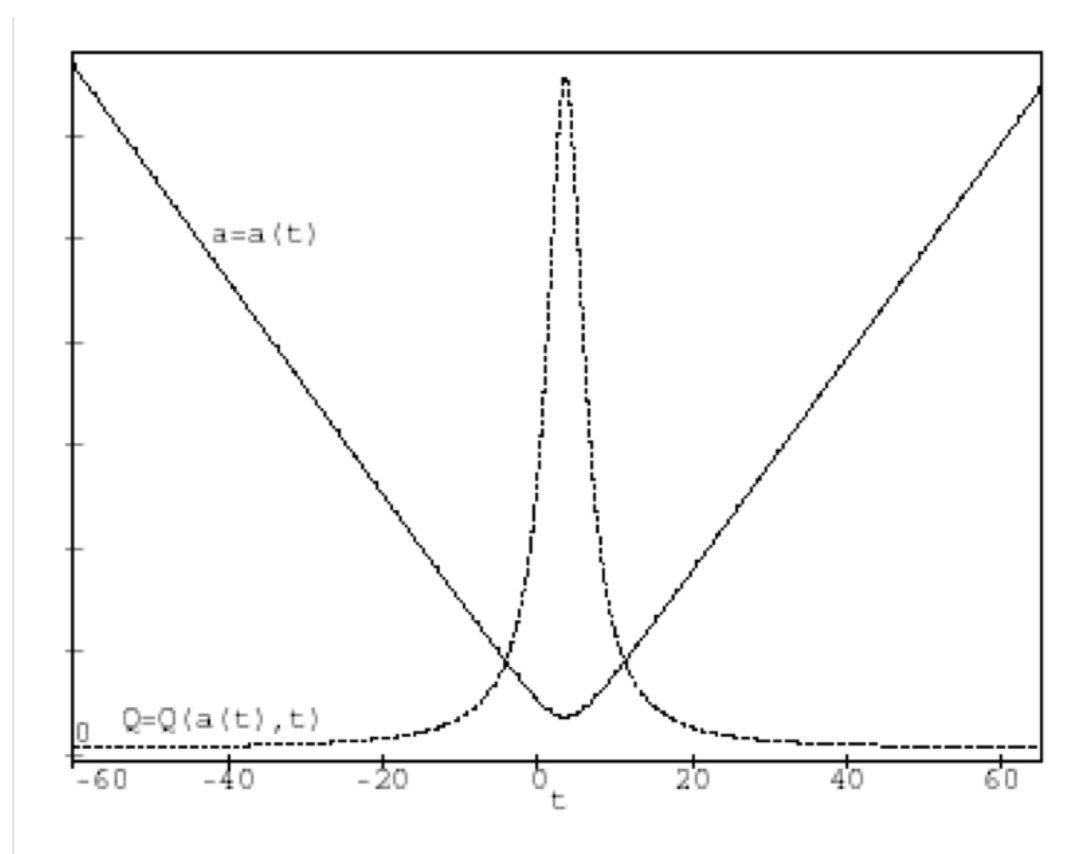
a quantum effect?



Pre Big Bang scenario:



quantum cosmology:



J. Acacio de Barros, N. Pinto-Neto & M. Sagorio-Leal
Phys. Lett. **A241**, 229 (1998)

S. Vitenti & PP
Mod.Phys.Lett. **A31**, 1640006 (2016).

Model listing:

Quantum gravity

LQG & LQC

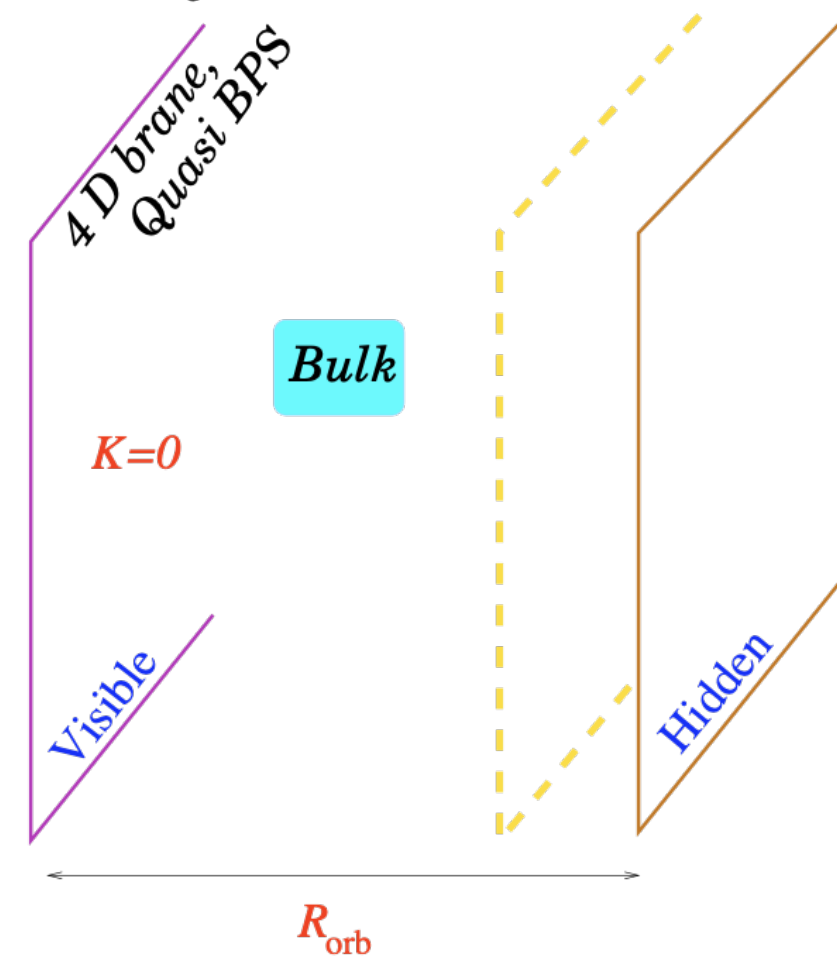
Non relativistic quantum gravity

Canonical quantum gravity (WdW)

Ekpyrotic & cyclic

String theory

Branes

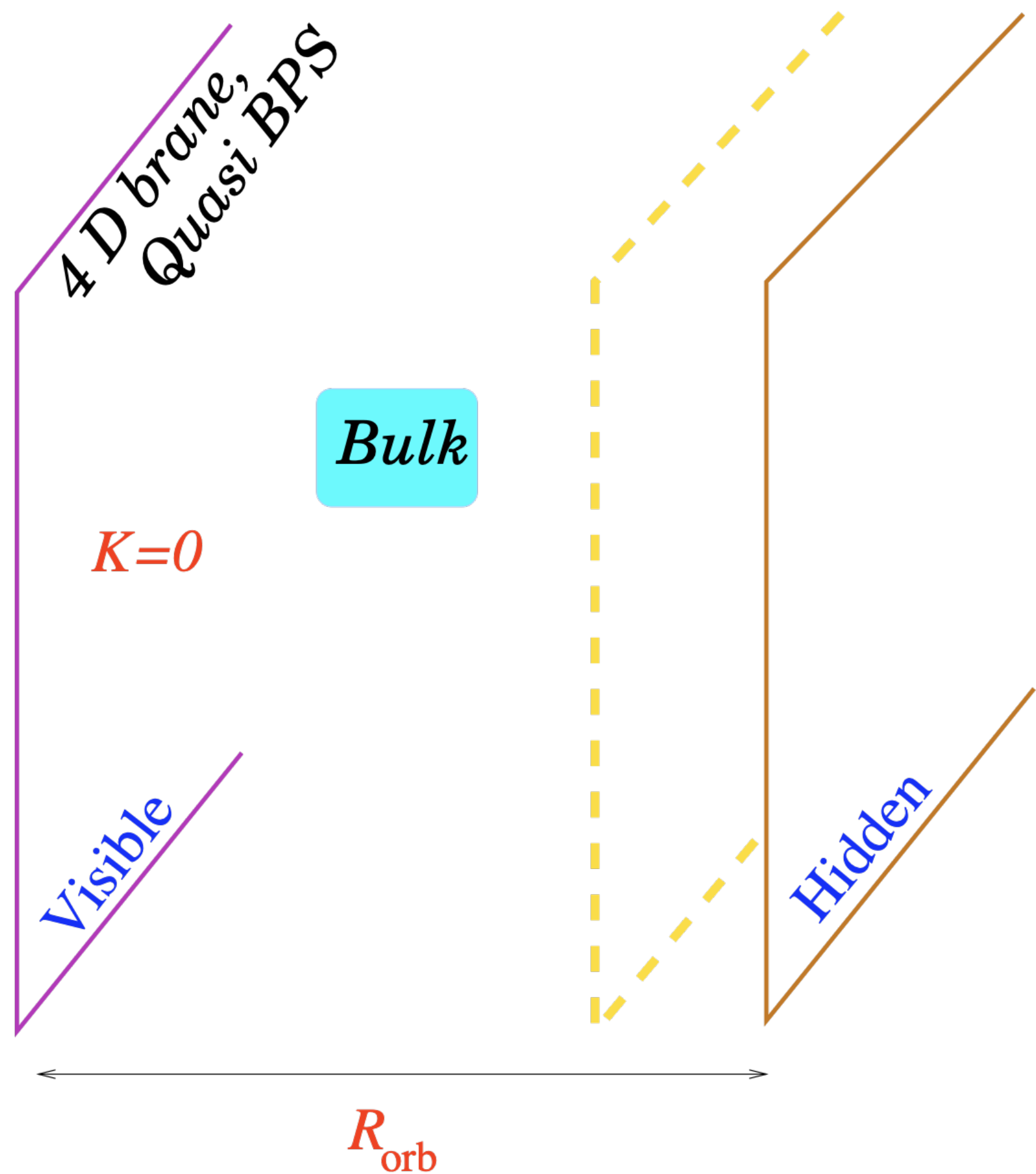
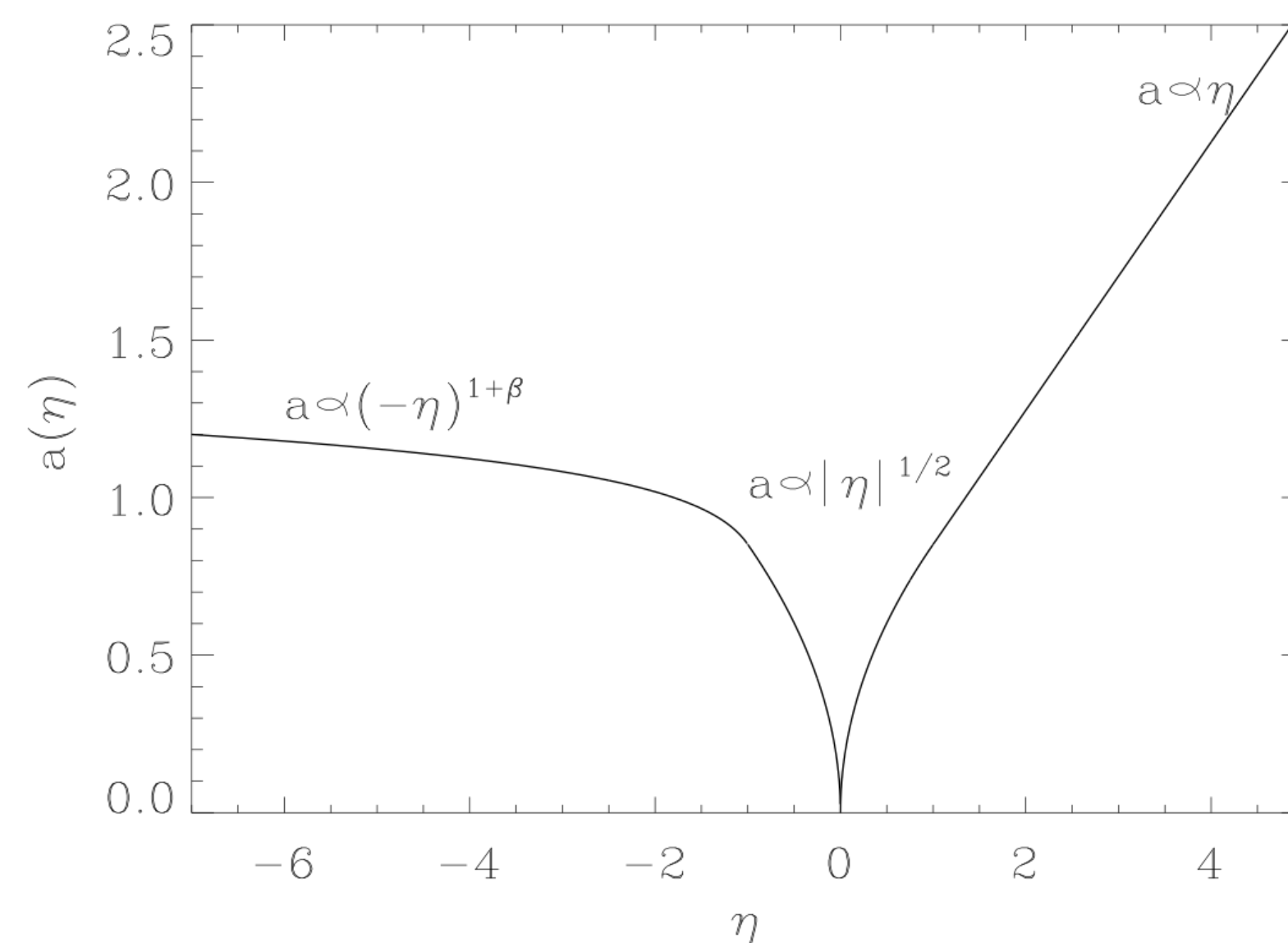


Ekpyrotic scenario:

$$\mathcal{S}_5 \propto \int_{\mathcal{M}_5} d^5x \sqrt{-g_5} \left[R_{(5)} - \frac{1}{2} (\partial\varphi)^2 - \frac{3}{2} \frac{e^{2\varphi} \mathcal{F}^2}{5!} \right],$$

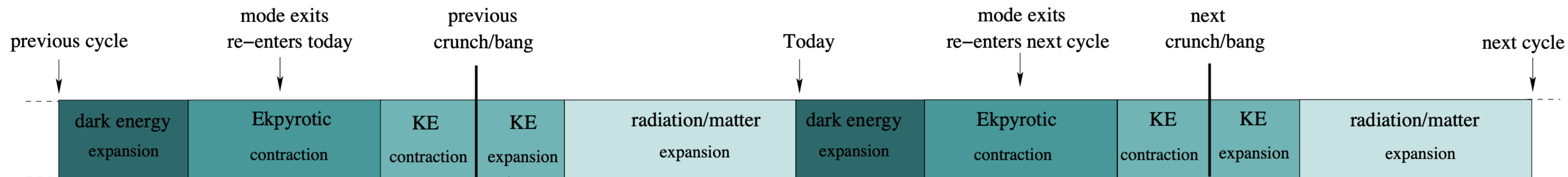
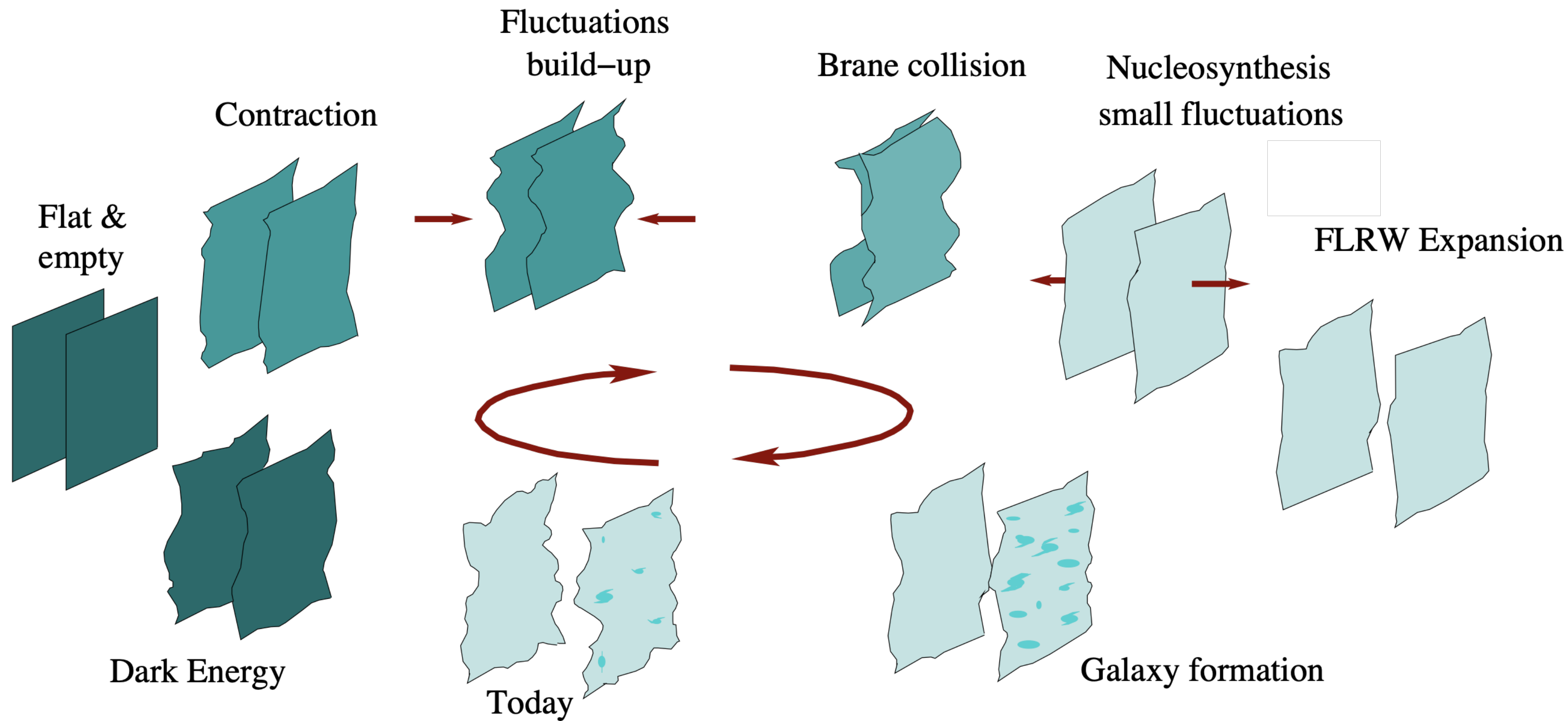
$$\mathcal{S}_4 = \int_{\mathcal{M}_4} d^4x \sqrt{-g_4} \left[\frac{R_{(4)}}{2\kappa} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right],$$

$$V(\varphi) = -V_i \exp \left[-\frac{4\sqrt{\pi\gamma}}{m_{\text{Pl}}} (\varphi - \varphi_i) \right],$$



Singular...

Cyclic extension



Model listing:

Quantum gravity

LQG & LQC

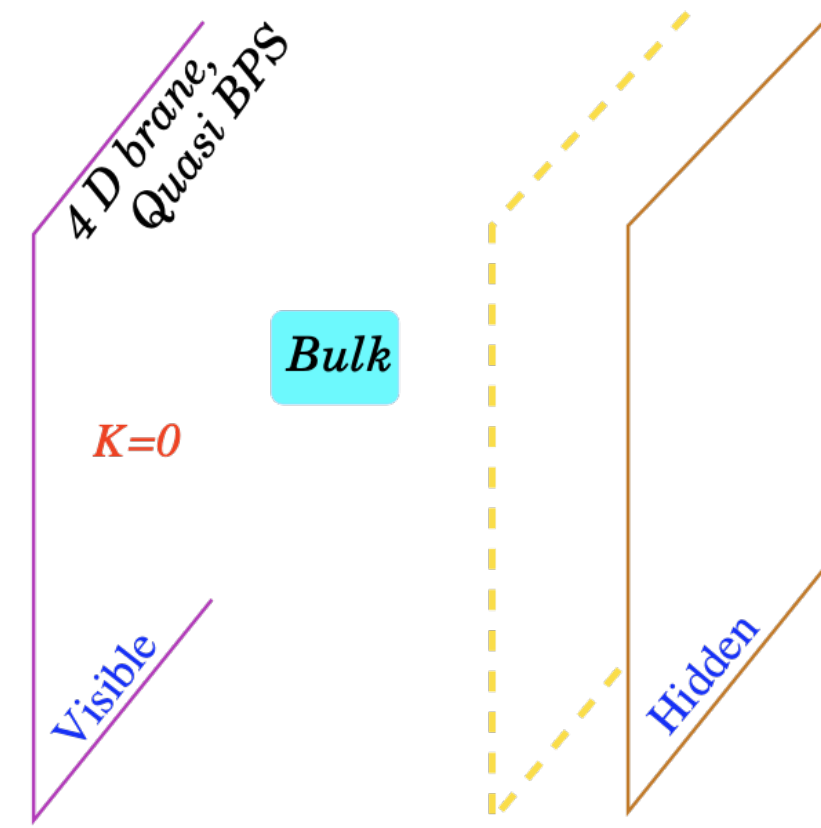
Non relativistic quantum gravity

Canonical quantum gravity (WdW)

Ekpyrotic & cyclic

String theory

Branes



Horava-Lifshitz

String gas cosmology

Lee-Wick & Quintom

Antigravity

$F(R)$, $f(T)$, Gauss-Bonnet

Galileon

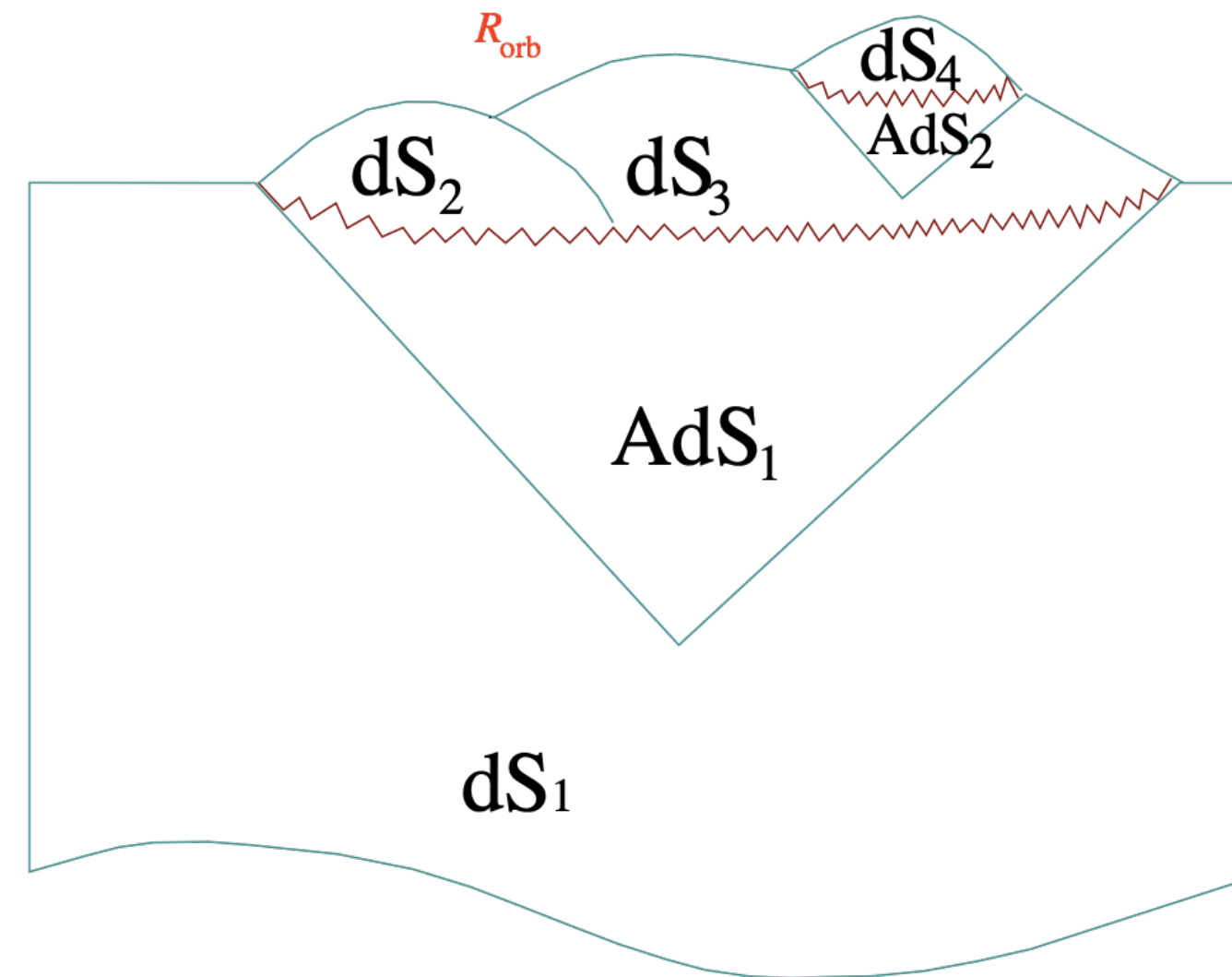
Massive gravity

Mimetic matter

Multiverse models

Non-linear electromagnetic action

Strings & AdS/CFT



Spinors & torsion

Standard Failures and bouncing solutions

Singularity

Merely a non issue in the bounce case!

Horizon

$$d_H \equiv a(t) \int_{t_i}^t \frac{d\tau}{a(\tau)}$$

can be made divergent easily if

$$t_i \rightarrow -\infty$$

Flatness

$$\frac{d}{dt} |\Omega - 1| = -2 \frac{\ddot{a}}{\dot{a}^3}$$

$$\ddot{a} < 0 \ \& \ \dot{a} < 0$$

accelerated expansion (**inflation**) or decelerated contraction (**bounce**)

Homogeneity

Large & flat Universe + low initial density + diffusion

$$\frac{t_{\text{dissipation}}}{t_{\text{Hubble}}} \propto \frac{\lambda}{R_H^{1/3}} \left(1 + \frac{\lambda}{AR_H^2} \right)$$

enough time to dissipate any wavelength



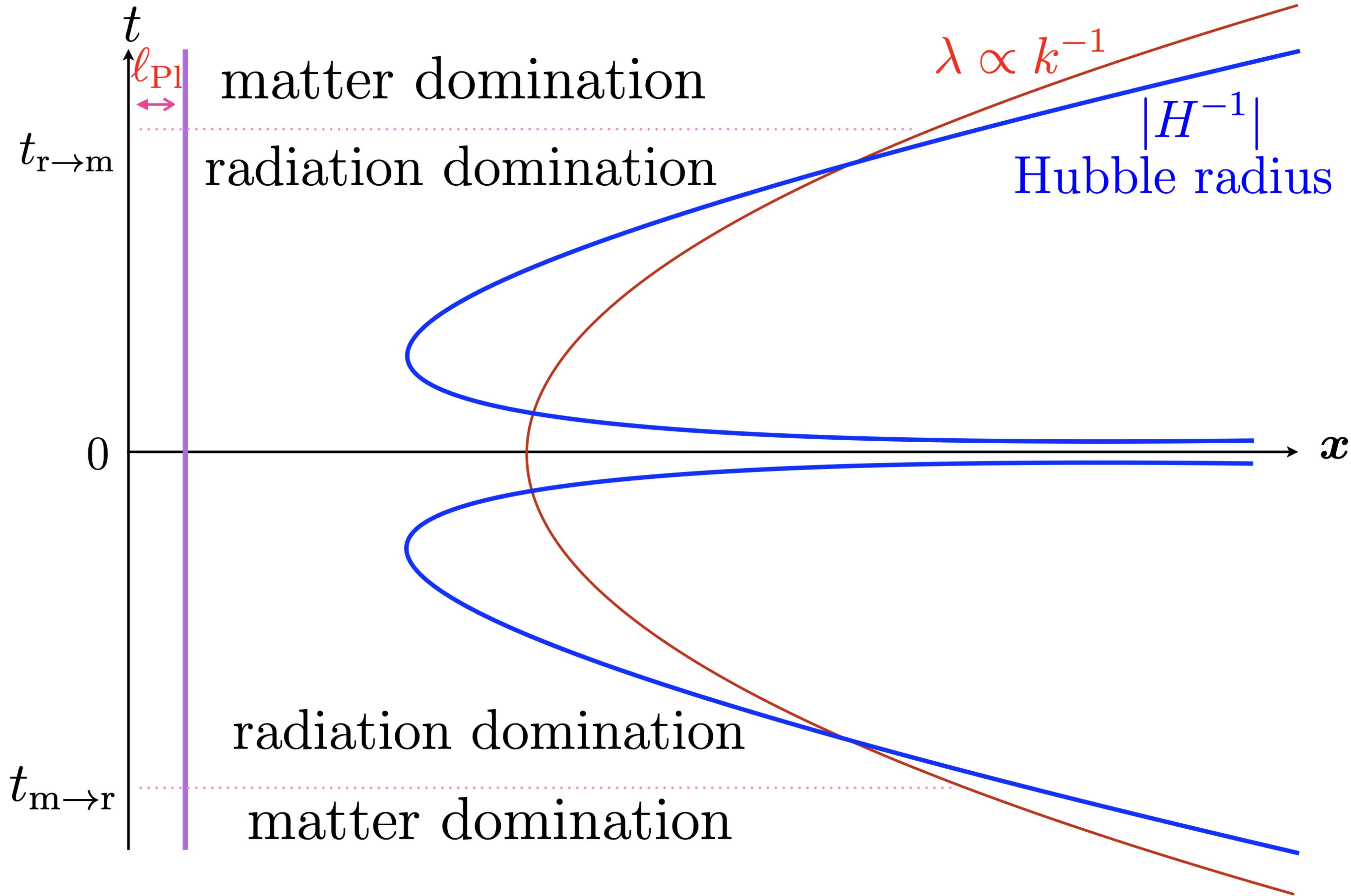
quantum vacuum fluctuations...

Isotropy

Potentially problematic: model dependent

Others

dark matter/energy, baryogenesis, ...



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$$d_{\text{H}}^{\text{cont}} = \frac{3(1+w)}{1+3w} t_{\text{end}} \left[1 - \left(\frac{t_{\text{ini}}}{t_{\text{end}}} \right)^{(1+3w)/[3(1+w)]} \right]$$

$$t_{\text{ini}} \rightarrow -\infty$$

Standard Failures and bouncing solutions

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quantum vacuum fluctuations...

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$$H^2 = \frac{1}{3} \left[-\frac{3\mathcal{K}}{a^2} + \frac{\rho_{m0}}{a^3} + \frac{\rho_{r0}}{a^4} + \frac{\rho_{\theta 0}}{a^6} + \dots + \frac{\rho_{\phi 0}}{a^{3(1+w_\phi)}} \right]$$

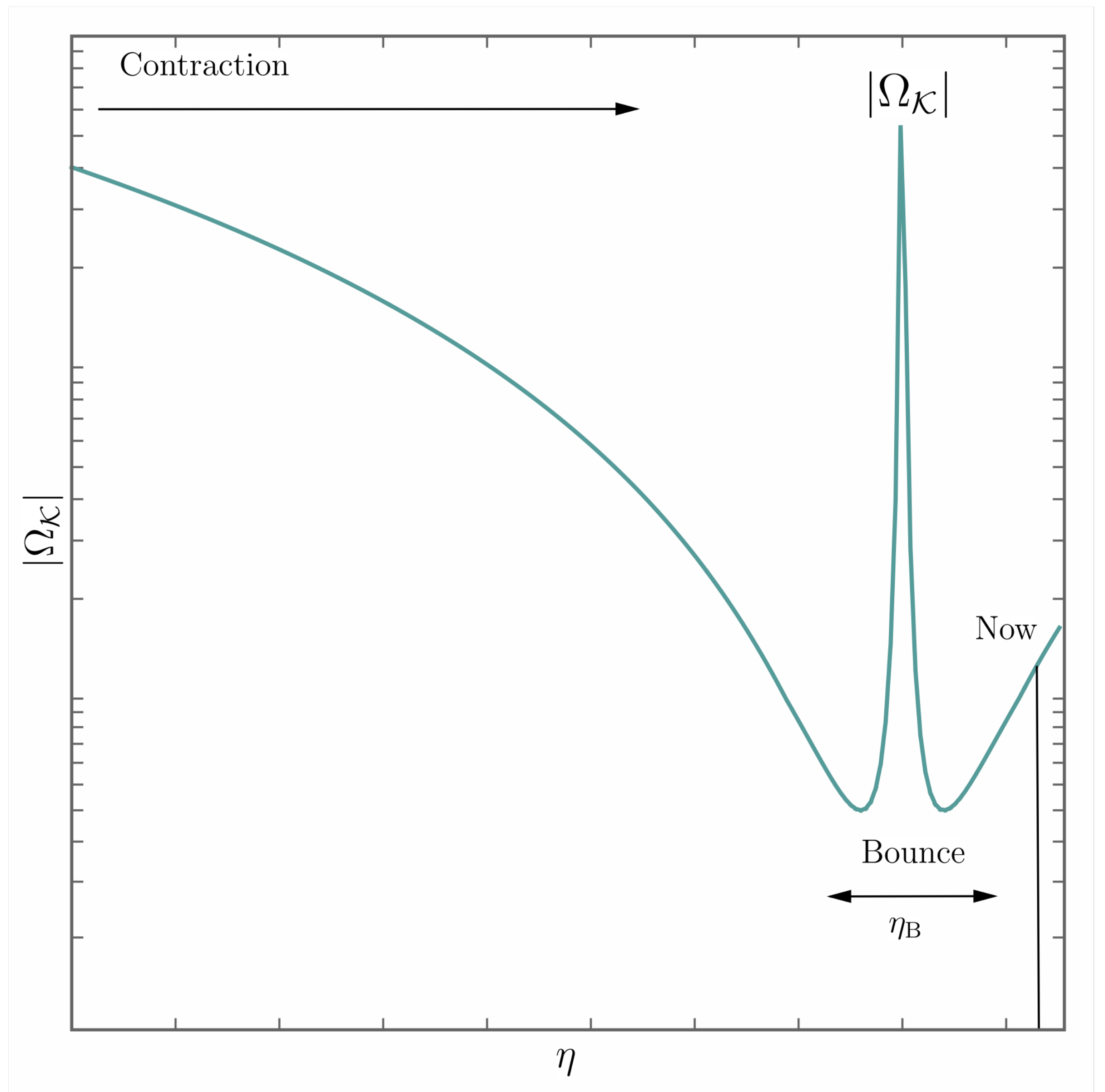
Critical density

$$\rho_c \equiv \frac{3H^2}{8\pi G_N}$$

\Rightarrow

$$\Omega \equiv \frac{\rho}{\rho_c}$$

Density parameter



Standard Failures and bouncing solutions

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Large & flat Universe + low initial density + diffusion

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enough time to dissipate any wavelength



quantum vacuum fluctuations...

Shear

Potentially problematic: model dependent

Others

dark matter/energy, baryogenesis, ...

The problem with contraction: BKL/shear instability

$$ds^2 = dt^2 - a^2(t) \sum_i e^{2\theta_i(t)} \sigma^i \sigma^i$$

Ricci flat:
 $\sigma^i = dx^i$

$$\sum_i \theta_i = 0$$

Average scale factor

$\frac{\dot{a}}{a}$ Mean Hubble parameter

$$H_i \equiv \frac{1}{ae^{\theta_i}} \frac{d}{dt} (ae^{\theta_i}) = H + \dot{\theta}_i$$

Friedman equations

$$H^2 = \frac{\rho_T}{3M_{Pl}^2} + \frac{1}{6} \sum_i \dot{\theta}_i^2$$

$$\dot{H} = -\frac{\rho_T + p_T}{2M_{Pl}^2} - \frac{1}{2} \sum_i \dot{\theta}_i^2$$

$$\ddot{\theta}_i + 3H\dot{\theta}_i = 0$$

$$\rho_{\text{shear}} \propto a^{-6}$$

Ekpyrotic solution:

$$w_{\text{ekp}} \gg 1 \implies \rho_{\text{ekp}} \propto a^{-3(1+w_{\text{ekp}})} \gg a^{-6} \text{ when } a \rightarrow 0$$

ultra stiff eq. of state

Problem: regular bounce \longrightarrow \exists phase with $w_{\text{bounce}} < -1$

So finally...

$$\rho_{\text{Shear}} \equiv \frac{M_{\text{Pl}}^2}{2} \sum_i \dot{\theta}_i^2 \propto a^{-6} \gg \rho_{\text{Fluid}}$$



Singularity!

Cosmic no-hair for ekpyrosis

- Bianchi I-VIII curvature always negative
- Impose $\rho \geq 0$
- Ultra stiff e.o.s., i.e. $\gamma = 1 + w > 2$



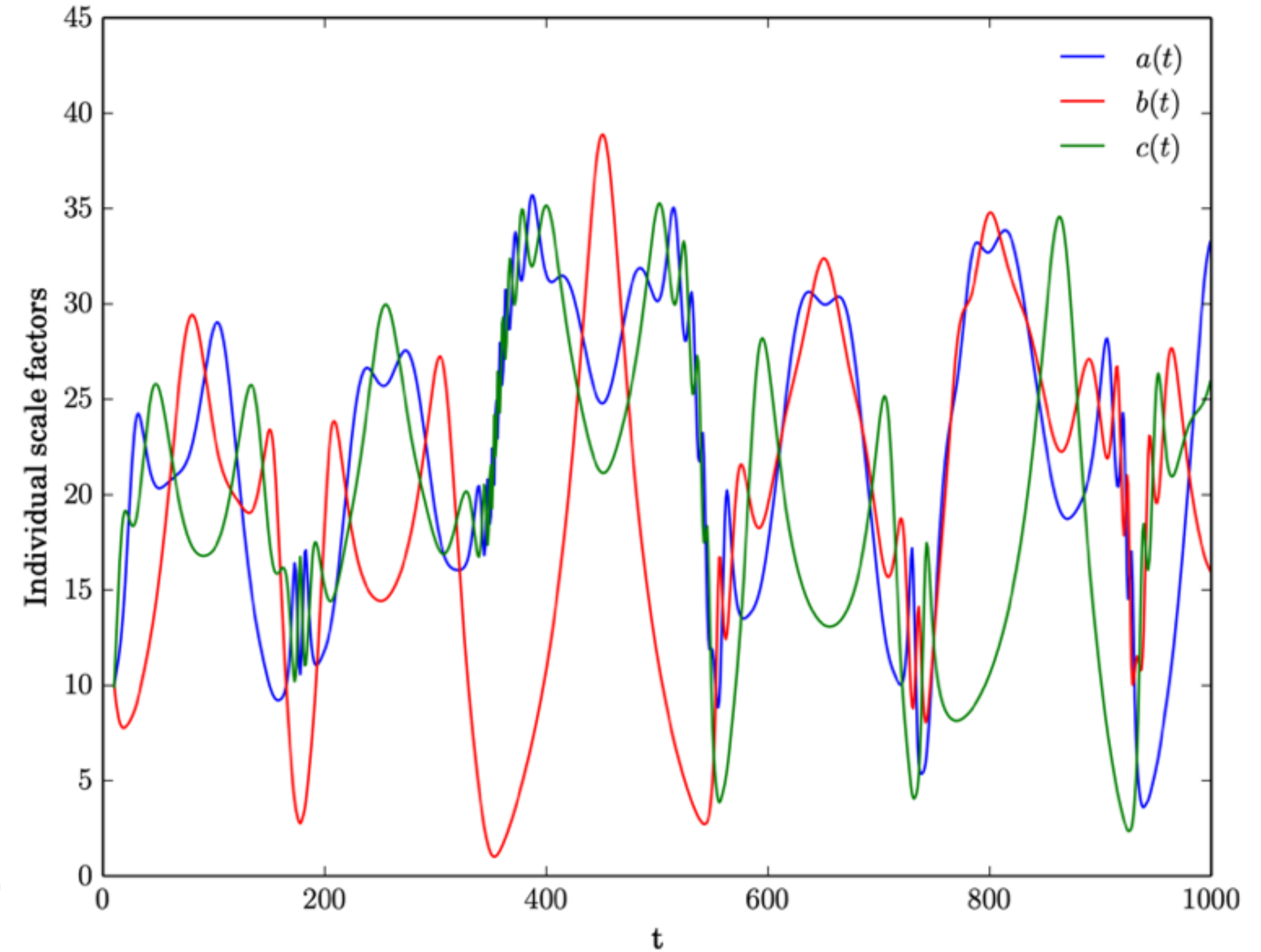
All initially contracting, spatially homogeneous, orthogonal Bianchi Type I-VIII cosmologies and all Bianchi type IX universes sourced by an ultra-stiff fluid with an equation of state such that $(\gamma - 2)$ is positive definite collapse into an isotropic singularity, where the sink is a spatially flat and isotropic FRW universe.

Anisotropic pressures and ekpyrosis

- Add stress-energy tensor anisotropies (& flat anisotropic Universe):
- $\dot{\sigma}_{ab} + 3H\sigma_{ab} = \pi_{ab}$
- Energy density anisotropies $\sigma^2 \sim \sigma_{ab}\sigma^{ab}$ can grow faster than V^{-2}
- Ekpyrotic fluid must grow faster... fine-tuned w.r.t. the anisotropic stress
- Expansion + 3-curvature anisotropy (Mixmaster) doesn't re-expand

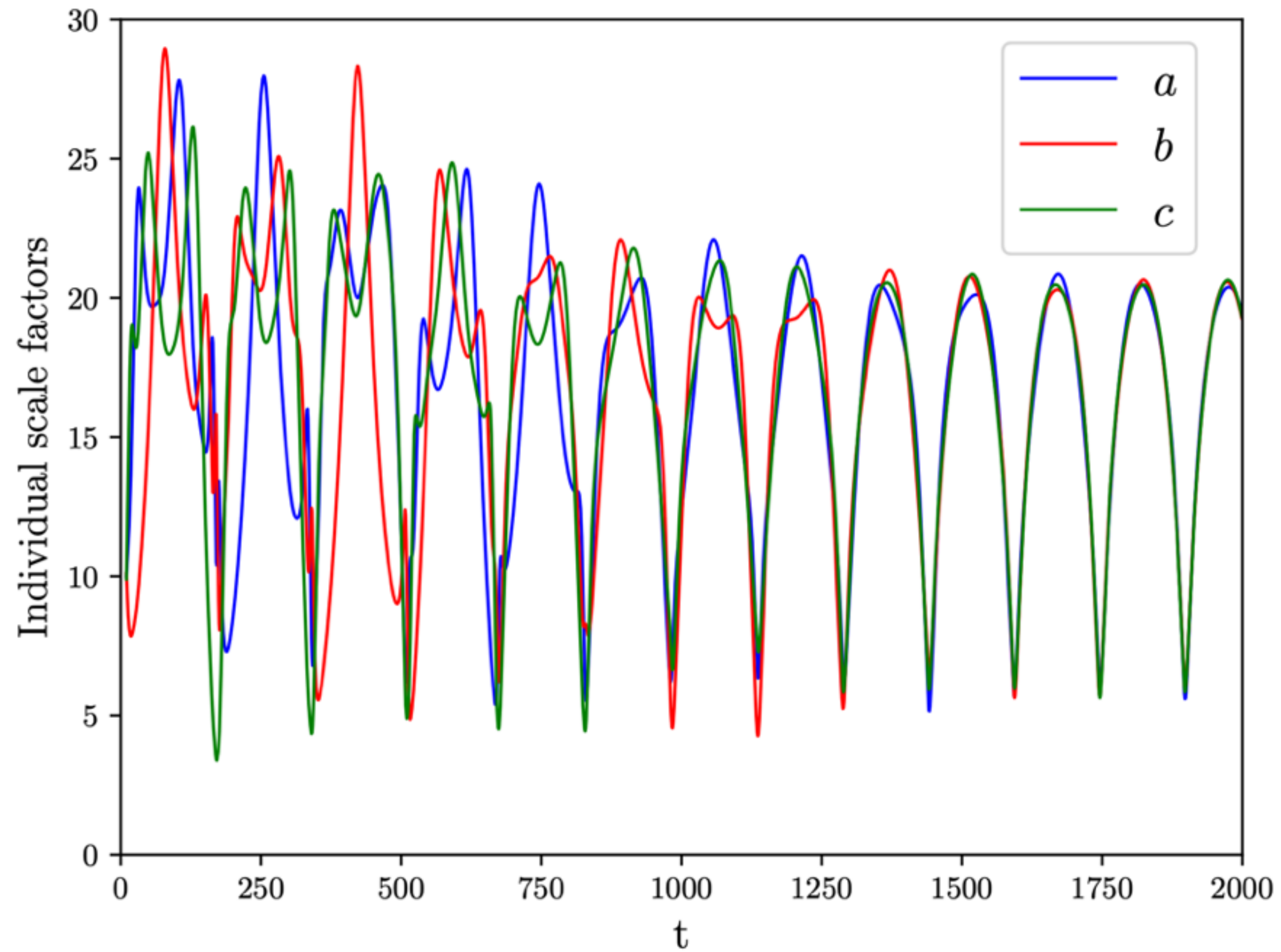
Mixmaster

- Curvature + expansion anisotropies
- Chaotic oscillations - infinite in number on approach to singularity
- Chaotic behaviour = attractor
- Anisotropy growth → collapse



Isotropisation by shear viscosity

$$\dot{\sigma}_{ab} + 3H\sigma_{ab} = \pi_{ab} \quad \& \quad \pi_{ab} = -\kappa\rho^{1/2}\sigma_{ab}$$



Cosmic no-hair reloaded

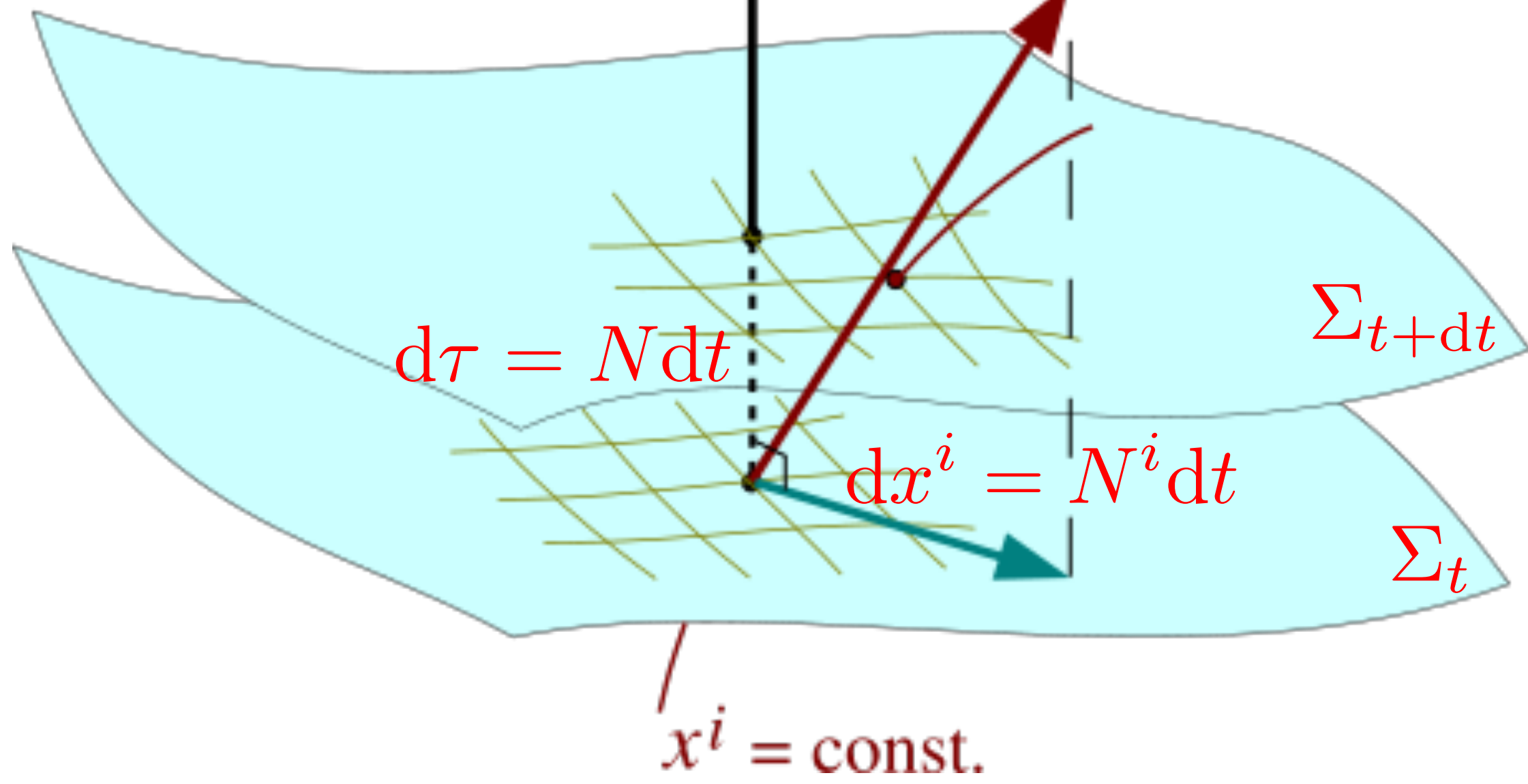
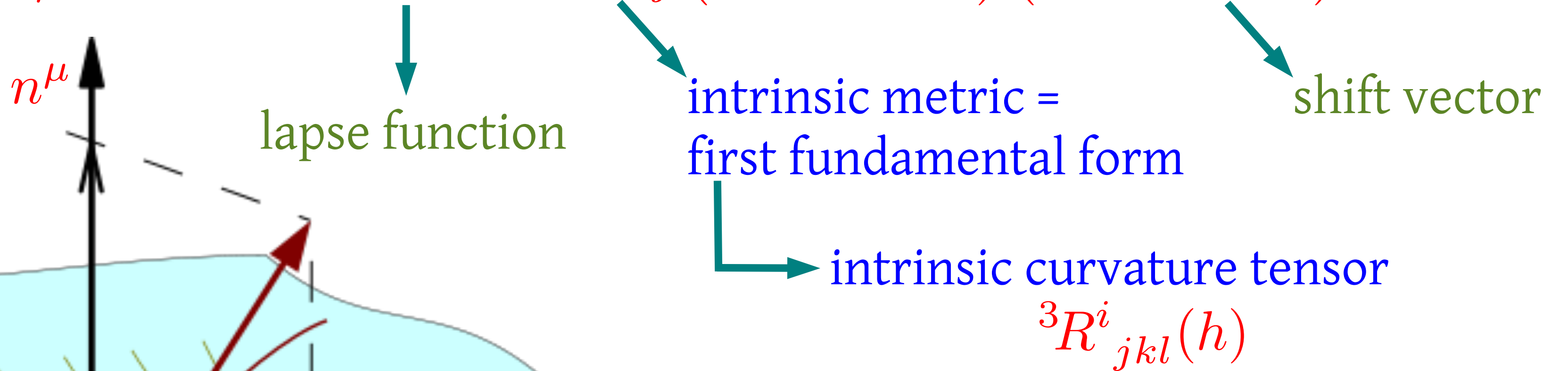
- Bianchi I-VIII curvature always negative
- impose $\rho \geq 0$
- Matter obeys Strong Energy Condition, i.e. $\rho + 3p > 0$
- Viscosity coefficient κ obeys $\kappa > 3(2 - \gamma)$

All initially contracting, spatially homogeneous, orthogonal Bianchi Type I-VIII cosmologies and all Bianchi type IX universes sourced by fluid obeying the Strong Energy Condition and with a shear viscous term such that the coefficient of viscosity κ obeys the condition that $\kappa > 3(2 - \gamma)$ collapse into an isotropic singularity, where the sink is a spatially flat and isotropic FRW universe.

Quantum cosmology

Hamiltonian GR (3+1)

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = -N^2 dt^2 + h_{ij} (dx^i + N^i dt) (dx^j + N^j dt)$$



extrinsic curvature = second fundamental form:

$$K_{ij} = -\nabla_j^{(h)} n_i = \frac{1}{2N} \left(\nabla_j^{(h)} N_i + \nabla_i^{(h)} N_j - \frac{\partial h_{ij}}{\partial t} \right)$$

Action (Einstein-Hilbert, compact space):

$$\mathcal{S} = \frac{1}{16\pi G_{\text{N}}} \left[\int_{\mathcal{M}} \sqrt{-g} (R - 2\Lambda) d^4x + 2 \int_{\partial\mathcal{M}} \sqrt{h} K^i_i d^3x \right] + \mathcal{S}_{\text{matter}} [\Phi(x)]$$

$$\longrightarrow \mathcal{S} = \int L dt = \frac{1}{16\pi G_{\text{N}}} \int dt \left[\int d^3x N \sqrt{h} (K_{ij} K^{ij} - K^2 + {}^3R - 2\Lambda) + L_{\text{matter}} \right]$$

Canonical momenta

$$\pi^{ij} \equiv \frac{\delta L}{\delta \dot{h}_{ij}} = -\frac{\sqrt{h}}{16\pi G_{\text{N}}} (K^{ij} - h^{ij} K)$$

$$\pi^{\Phi} \equiv \frac{\delta L}{\delta \dot{\Phi}} = -\frac{\sqrt{h}}{N} \left(\dot{\Phi} - N \frac{\partial \Phi}{\partial x^i} \right)$$

$$\left. \begin{aligned} \pi^0 &\equiv \frac{\delta L}{\delta \dot{N}} \approx 0 \\ \pi^i &\equiv \frac{\delta L}{\delta \dot{N}^i} \approx 0 \end{aligned} \right\} \text{primary constraints}$$

Hamiltonian

$$H \equiv \int d^3x \left(\pi^0 \dot{N} + \pi^i \dot{N}_i + \pi^{ij} \dot{h}_{ij} + \pi^\Phi \dot{\Phi} \right) - L$$

$$= \int d^3x \left(\pi^0 \dot{N} + \pi^i \dot{N}_i + N\mathcal{H} + N_i\mathcal{H}^i \right)$$

$$\mathcal{H} = \frac{1}{\sqrt{h}} \left(h_{ik}h_{jl} - \frac{1}{2}h_{ij}h_{kl} \right) \pi^{ij}\pi^{kl} - \sqrt{h} {}^3R$$


$$\mathcal{H}^i = -2\sqrt{h}\nabla_j \left(\frac{\pi^{ij}}{\sqrt{h}} \right)$$

variation wrt lapse: $\mathcal{H} = 0 \rightarrow$ Hamiltonian constraint
variation wrt shift: $\mathcal{H}^i = 0 \rightarrow$ momentum constraint

} \implies classical description complete

Superspace & canonical quantization

relevant configuration space $\text{Riem}(\Sigma) \equiv \{h_{ij}(x^\mu), \Phi(x^\mu) \mid x \in \Sigma\}$



parameters

GR \implies invariance/diffeomorphisms \implies Conf = $\frac{\text{Riem}(\Sigma)}{\text{Diff}(\Sigma)}$: superspace

Wave functional $\Psi[h_{ij}(x), \Phi(x)] = \langle h_{ij}, \Phi \mid \Psi \rangle$

+ Dirac canonical quantization procedure

$$\pi^{ij} \rightarrow -i \frac{\delta}{\delta h_{ij}}$$

$$\pi^\Phi \rightarrow -i \frac{\delta}{\delta \Phi}$$

$$\pi^0 \rightarrow -i \frac{\delta}{\delta N}$$

$$\pi^i \rightarrow -i \frac{\delta}{\delta N_i}$$

primary constraints

$$\left\{ \begin{array}{l} \hat{\pi}^0 = -i \frac{\delta \Psi}{\delta N} = 0 \\ \hat{\pi}^i = -i \frac{\delta \Psi}{\delta N_i} = 0 \end{array} \right.$$

momentum $\hat{\mathcal{H}}^i \Psi = 0 \implies i \nabla_j^{(h)} \left(\frac{\delta \Psi}{\delta h_{ij}} \right) = 8\pi G_N \hat{T}^{0i} \Psi$

same Ψ for configurations related by a coordinate transformation

Hamiltonian

$$\hat{\mathcal{H}}\Psi = \left[-16\pi G_N \mathcal{G}_{ijkl} \frac{\delta^2}{\delta h_{ij} \delta h_{kl}} + \frac{\sqrt{h}}{16\pi G_N} \left(-{}^3R + 2\Lambda + 16\pi G_N \hat{T}^{00} \right) \right] \Psi = 0$$

Wheeler - De Witt equation

$$\mathcal{G}_{ijkl} = \frac{1}{2\sqrt{h}} (h_{ik}h_{jl} + h_{il}h_{jk} - h_{ij}h_{kl})$$

De Witt metric

primary constraints

$$\left\{ \begin{array}{l} \hat{\pi}^0 = -i \frac{\delta \Psi}{\delta N} = 0 \\ \hat{\pi}^i = -i \frac{\delta \Psi}{\delta N_i} = 0 \end{array} \right.$$

momentum $\hat{\mathcal{H}}^i \Psi = 0 \implies i \nabla_j^{(h)} \left(\frac{\delta \Psi}{\delta h_{ij}} \right) = 8\pi G_N \hat{T}^{0i} \Psi$

same Ψ for configurations related by a coordinate transformation

Hamiltonian

$$\hat{\mathcal{H}} \Psi = 0$$

TIMELESS Schrödinger equation

mini-superspace

*restrict attention from an infinite dimensional configuration space to a 2 dimensional space
= mini-superspace*

$$h_{ij} dx^i dx^j \mapsto a^2(t) \left[\frac{dr^2}{1 - \mathcal{K}r^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

WDW equation becomes Schrödinger like for $\Psi [a(t), \phi(t)]$

Conceptual & technical issues

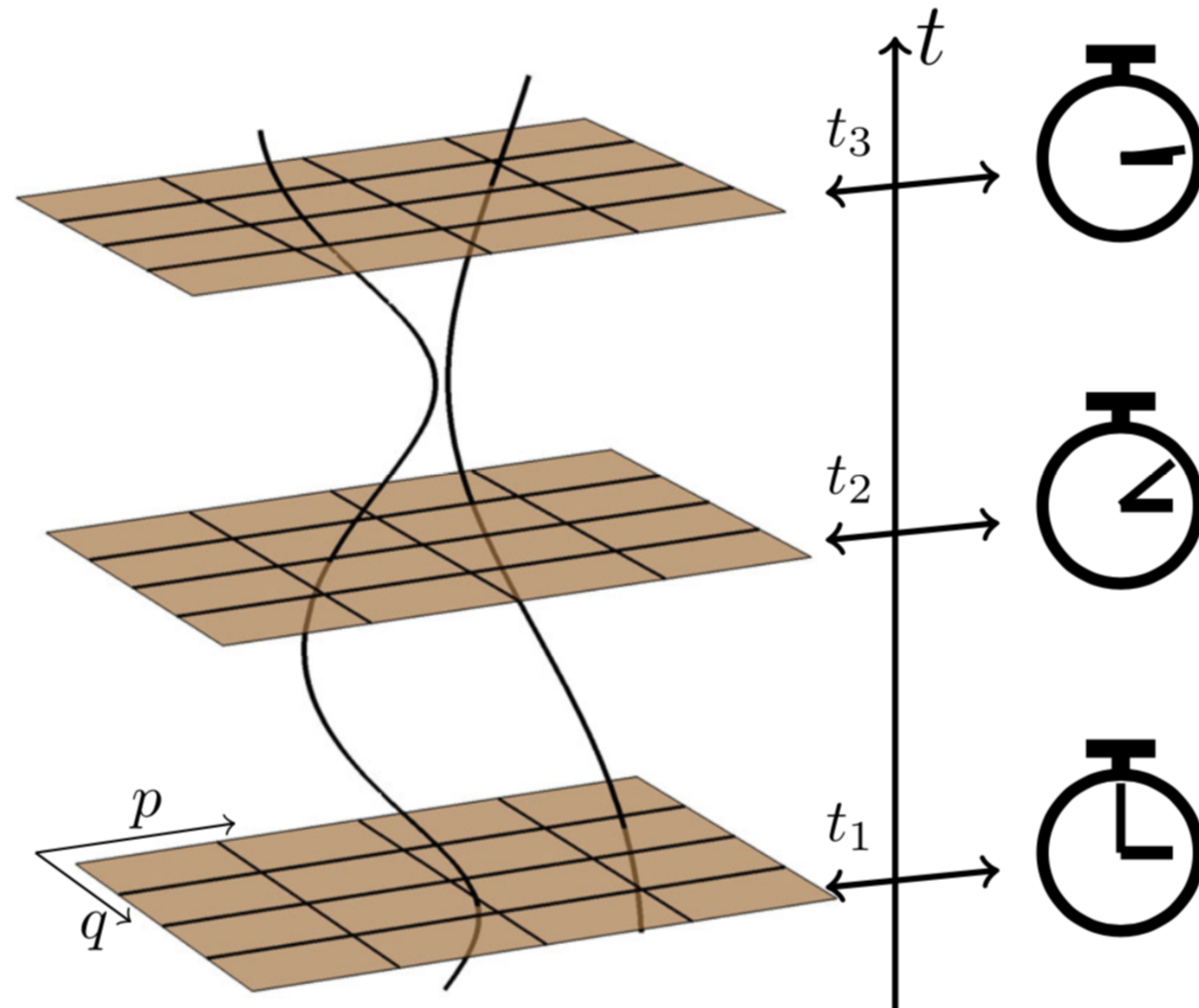
infinite # d.o.f. to a few: mathematical consistency?

freeze momenta... Heisenberg uncertainties?

[quantization, minisuperspace] $\neq 0$

The clock issue in quantum cosmology

- GR = constrained system: lack of external time
- arbitrary degree of freedom: internal clock



Classical system q_i & p_i

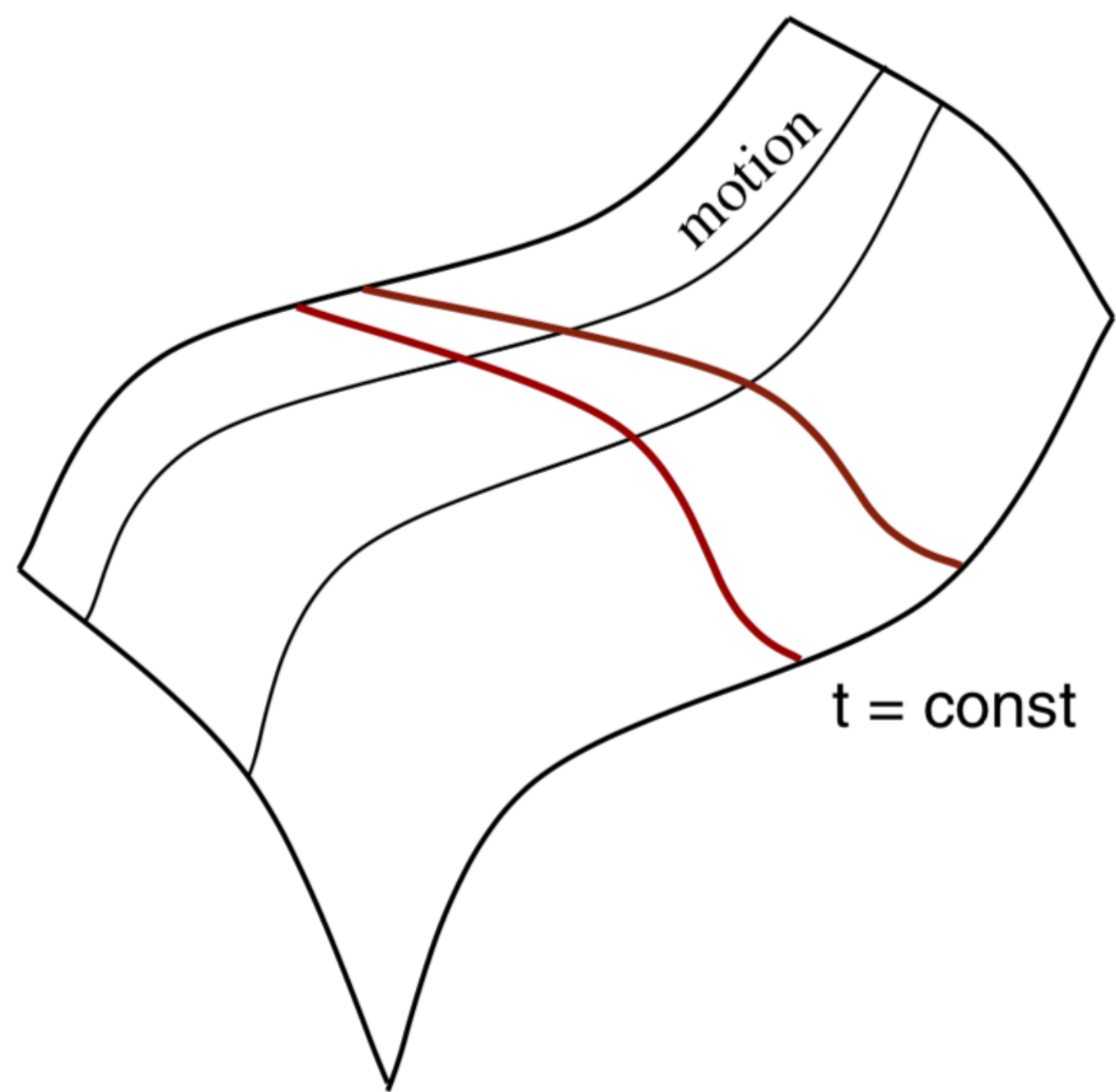
Constraint $C(q_i, p_i) = 0$ & $\frac{d}{d\tau} \mathcal{O}(q_i, p_i) = \{\mathcal{O}, C\}_{\text{P.B}}$

evolution
parameter
(time)

observable

Time parametrization invariance $\tau \rightarrow \tau' \longrightarrow N(q_i, p_i, \tau)$

arbitrary non vanishing lapse function

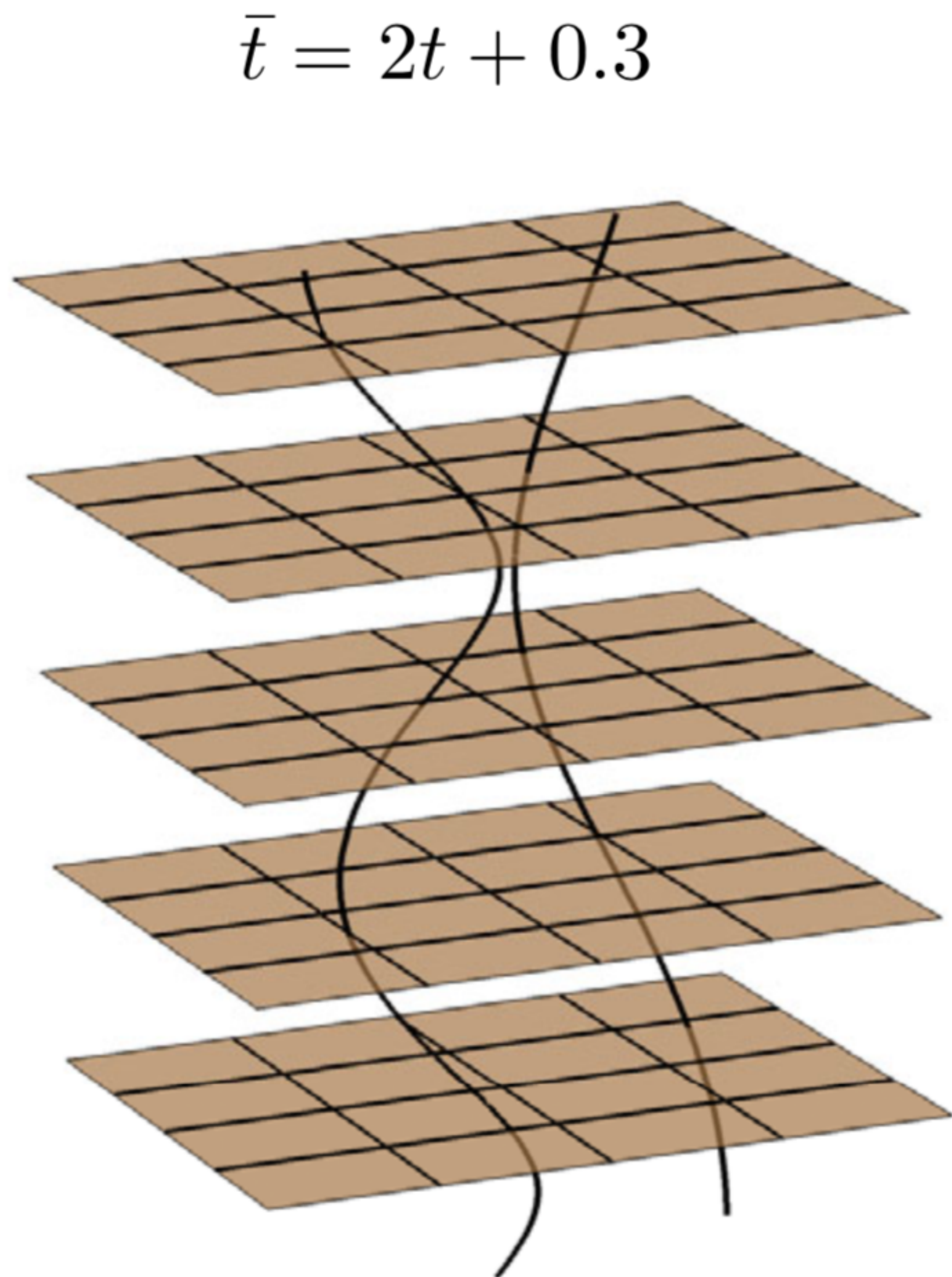
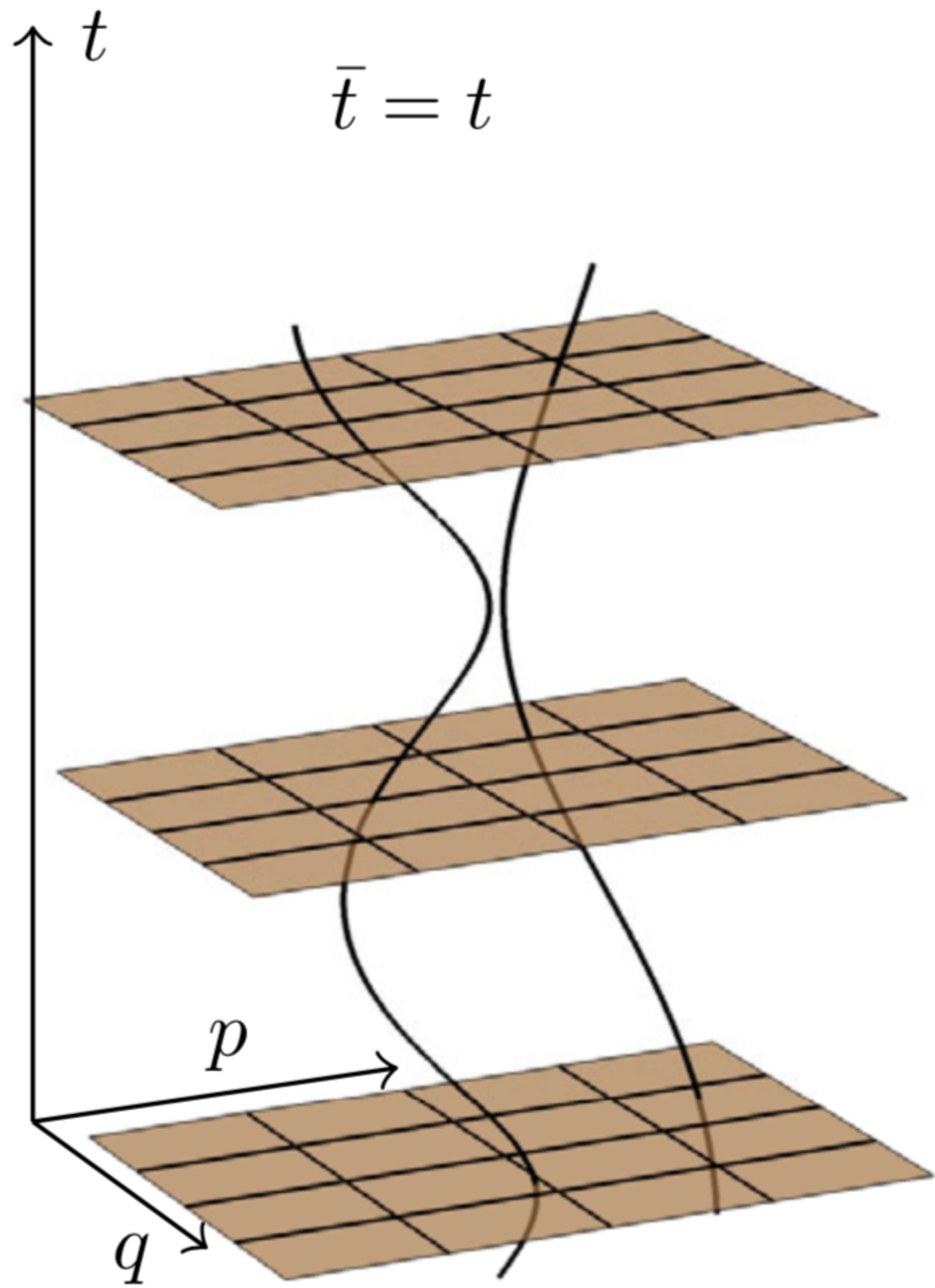


$$d\tau = N d\tau'$$

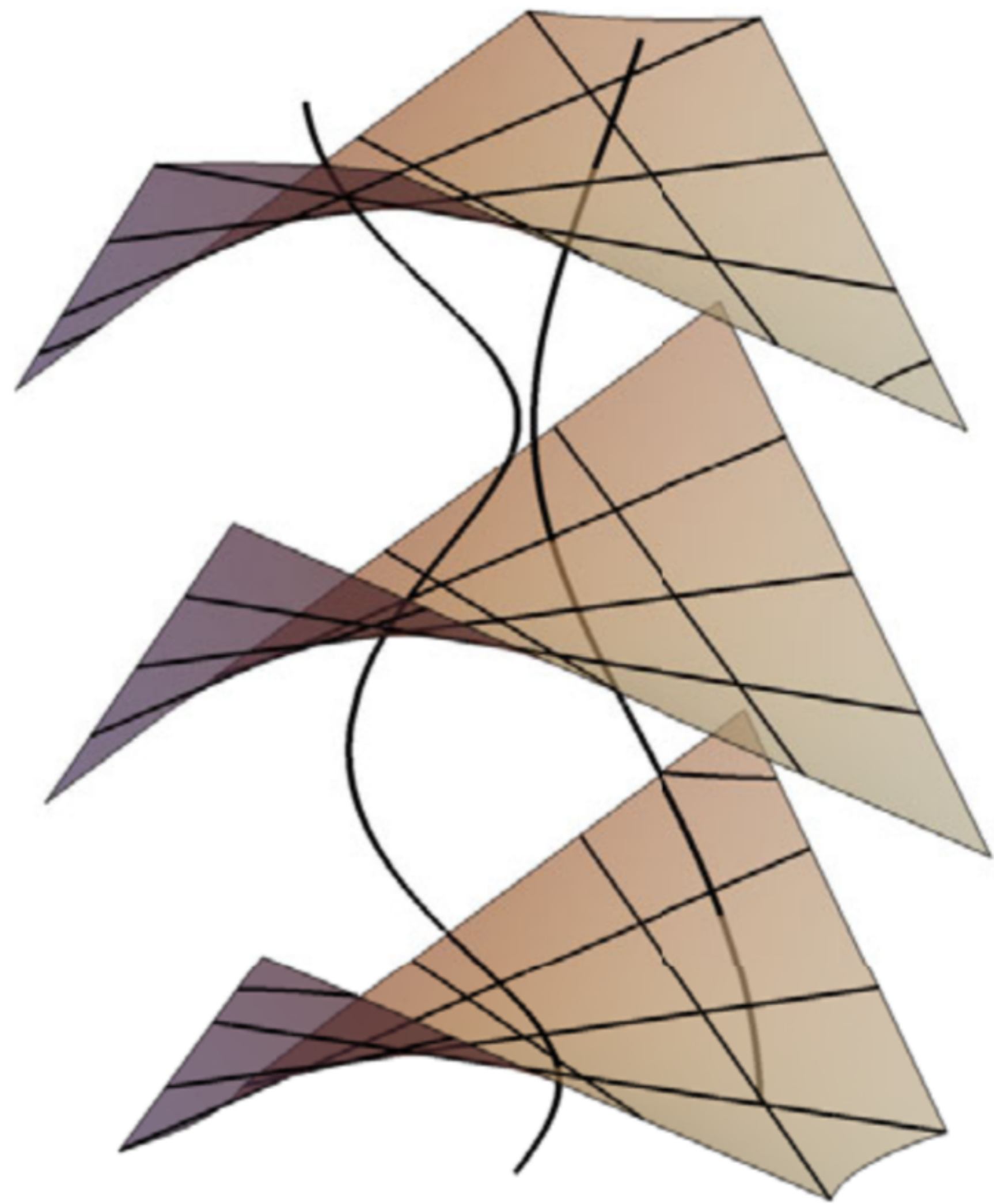
\implies

$$\frac{d}{d\tau'} \mathcal{O}(q_i, p_i) = \{\mathcal{O}, NC\}_{\text{P.B}}$$

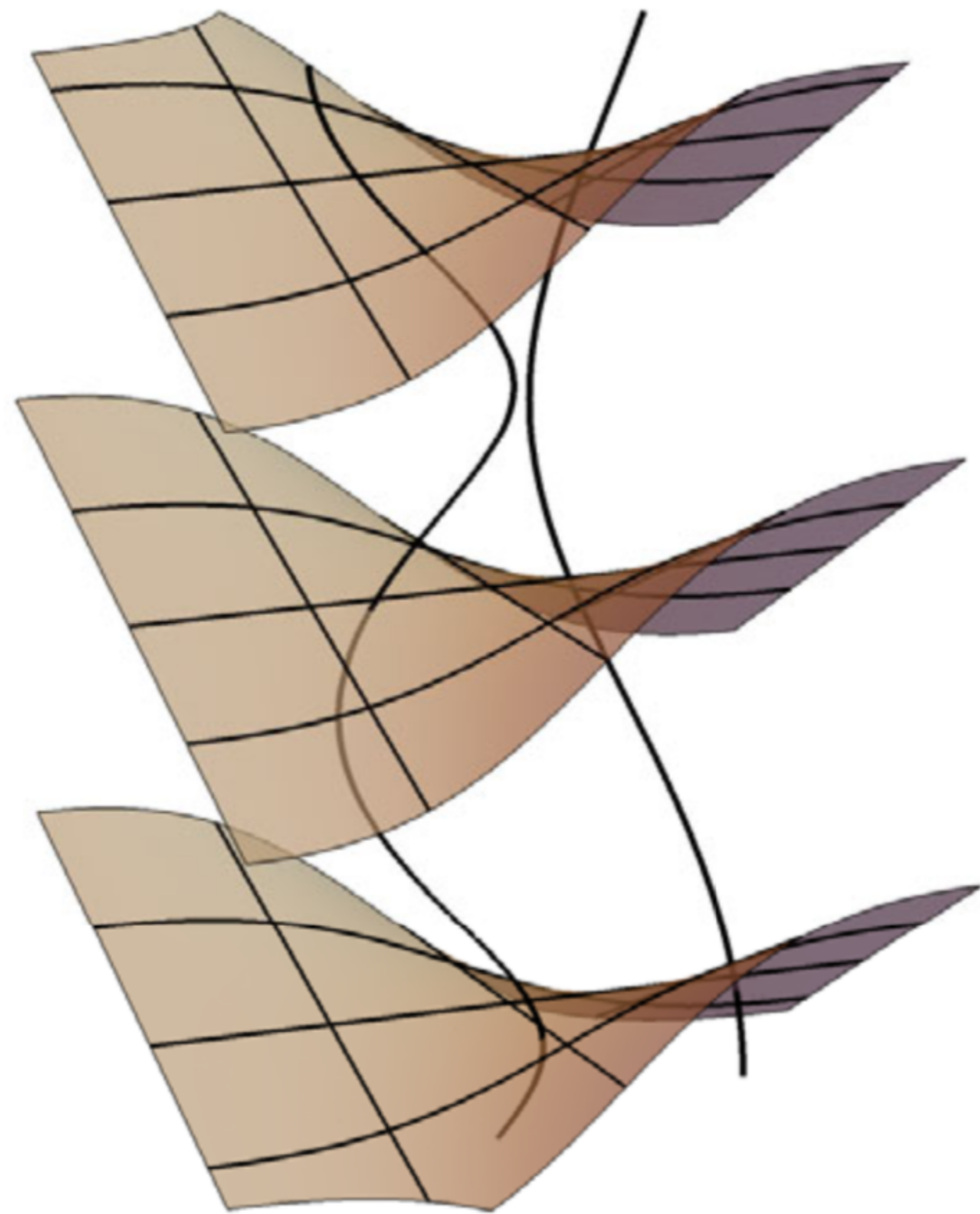
hamiltonian $H = NC$



$$\bar{t} = t + qp$$



$$\bar{t} = t - \frac{3qp}{3p^2 + 1}$$



Constrained system $C(\{q^k\}, \{p_k\}) = H_{\text{tot}}(\{q^k\}, \{p_k\}) = 0$

Canonical transformation $(\{q^k\}, \{p_k\}) \mapsto (\{Q^a\}, \{P_a\})$

$\longrightarrow \exists Q^\alpha$ such that $\{Q^\alpha, H_{\text{tot}}\}_{\text{P.B.}} = 1 = \frac{dQ^\alpha}{dt}$?

Quantum system $\hat{H}_{\text{tot}}\Psi \equiv \hat{C}\Psi(Q^a) = 0$

\hat{C} linear in $\hat{P}_\alpha = -i\frac{\partial}{\partial Q^\alpha}$

defines clock

will become time

$\longrightarrow \hat{C} = \hat{P}_\alpha + \hat{H}(P_1, \dots, P_{\alpha-1}, P_{\alpha+1}, \dots, P_n, \{Q^a\})$

$\hat{C}\Psi(Q^a) = 0 \implies$ time-dependent
Schrödinger equation

A simple example

$$\mathcal{S} = \frac{1}{2} \int \left(\frac{\dot{x}^2}{z} - zx^2 - \frac{\dot{y}^2}{z} + zy^2 \right) dt$$

First, redefine time: $d\tau = zdt \longrightarrow \mathcal{S} = \frac{1}{2} \int \left[\left(\frac{dx}{d\tau} \right)^2 - x^2 - \left(\frac{dy}{d\tau} \right)^2 + y^2 \right] d\tau$

Classical EOMs $\left\{ \begin{array}{l} \frac{d^2x}{d\tau^2} = -x \\ \frac{d^2y}{d\tau^2} = -y \end{array} \right. \longrightarrow 2 \text{ independent harmonic oscillators}$
 $H_{\text{tot}} = H_x + H_y$

Canonical transformation $T = \arctan\left(\frac{p_y}{y}\right)$ & $P_T = -\frac{1}{2}(p_y^2 + y^2) = H_y$

$$\{T, P_T\}_{\text{P.B.}} = 1$$

$$\begin{array}{c} \downarrow \\ \{T, H_{\text{tot}}\}_{\text{P.B.}} = 1 \end{array} \longrightarrow \frac{dT}{d\tau} = 1 \longrightarrow \begin{cases} \frac{dx}{dT} = p_x \\ \frac{dp_x}{dT} = -x \end{cases}$$

$$H_{\text{tot}} = P_T + H \text{ on shell } H_{\text{tot}} \approx 0$$

Quantization: x only! $\hat{H}_{\text{tot}}\psi(x, T) = 0 \implies i\frac{\partial\psi}{\partial T} = \hat{H}\psi(x, T)$

$$\& \int |\psi(x, T)|^2 dx = 1$$

y remains classical (clock)


Bianchi I case

$$ds^2 = -N^2 d\tau^2 + \sum_{i=1}^3 a_i^2 (dx^i)^2$$

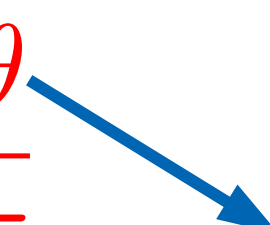
Scale factors

$$\begin{cases} a_1 &= e^{\beta_0 + \beta_+ + \sqrt{3}\beta_-} \\ a_2 &= e^{\beta_0 + \beta_+ - \sqrt{3}\beta_-} \\ a_3 &= e^{\beta_0 - 2\beta_+} \end{cases}$$

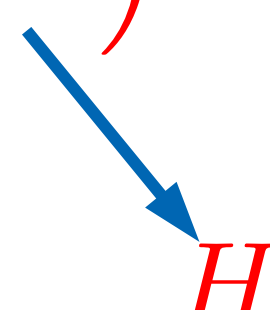
Volume $V \equiv a_1 a_2 a_3 = e^{3\beta_0}$


$$d\beta_0 = \frac{1}{3} e^{-3\beta_0} dV$$

Action $\mathcal{S} = \int d\tau \left(\underbrace{p_0 \dot{\beta}_0 + p_+ \dot{\beta}_+ + p_- \dot{\beta}_-}_{\frac{d\theta}{d\tau}} - NC \right)$



canonical one-form



H

constraint

$$C = \frac{e^{-3\beta_0}}{24} (-p_0^2 + p_+^2 + p_-^2)$$

ensure canonical one-form remains canonical $p_V \equiv \frac{e^{-3\beta_0}}{3} p_0$

$$d\theta = p_V dV + p_+ d\beta_+ + p_- d\beta_-$$

constraint $C = \frac{3V}{8} \left(-p_V^2 + \frac{p_+^2 + p_-^2}{9V^2} \right)$

cyclic variable $\dot{p}_\pm = 0$ set $p_+ = k \cos \alpha$ and $p_- = k \sin \alpha$

$$\longrightarrow d\theta = p_V dV + p_k dk + p_\alpha d\alpha + \underbrace{d(k \cos \alpha \beta_+ + k \sin \alpha \beta_-)}_{\rightarrow \text{exact... ignore!}}$$

$$p_k \equiv -(\cos \alpha \beta_+ + \sin \alpha \beta_-),$$

$$p_\alpha \equiv (k \sin \alpha \beta_+ - k \cos \alpha \beta_-)$$

neither α nor P_α in $H = NC$

the system reduces to

$$\begin{cases} d\theta &= p_V dV + p_k dk \\ C &= \frac{3V}{8} \left(-p_V^2 + \frac{k^2}{9V^2} \right) \end{cases}$$

Hamilton equations

$$\dot{k} = 0$$

$$\dot{p}_k = -N \frac{k}{12V}$$

$$\dot{V} = -N \frac{3V p_V}{4}$$

$$\dot{p}_V = -N \left[\frac{3}{8} \left(-p_V^2 + \frac{k^2}{9V^2} \right) - \frac{k^2}{12V^2} \right]$$

+ constraint

$$\frac{3V}{8} \left(-p_V^2 + \frac{k^2}{9V^2} \right) = 0$$

→ closed for V and p_V

Choosing a time

$$\frac{d}{d\tau} \left(9 \frac{p_k}{k} \right) = -\frac{3}{4} \frac{N}{V} \quad \text{monotonically increasing function}$$

valid time choice $\tau = \frac{9p_k}{k} \implies N = -\frac{4}{3}V$

Solving directly in the action $\mathcal{S} = \int d\theta = \int d\tau \left(p_V \dot{V} - \frac{V^2 p_V^2}{2} \right)$

\downarrow
 H

classical unconstrained one dimensional system

$$\frac{d}{d\tau} (V p_V) = 0 \implies V p_V = V_0 p_{V0}$$

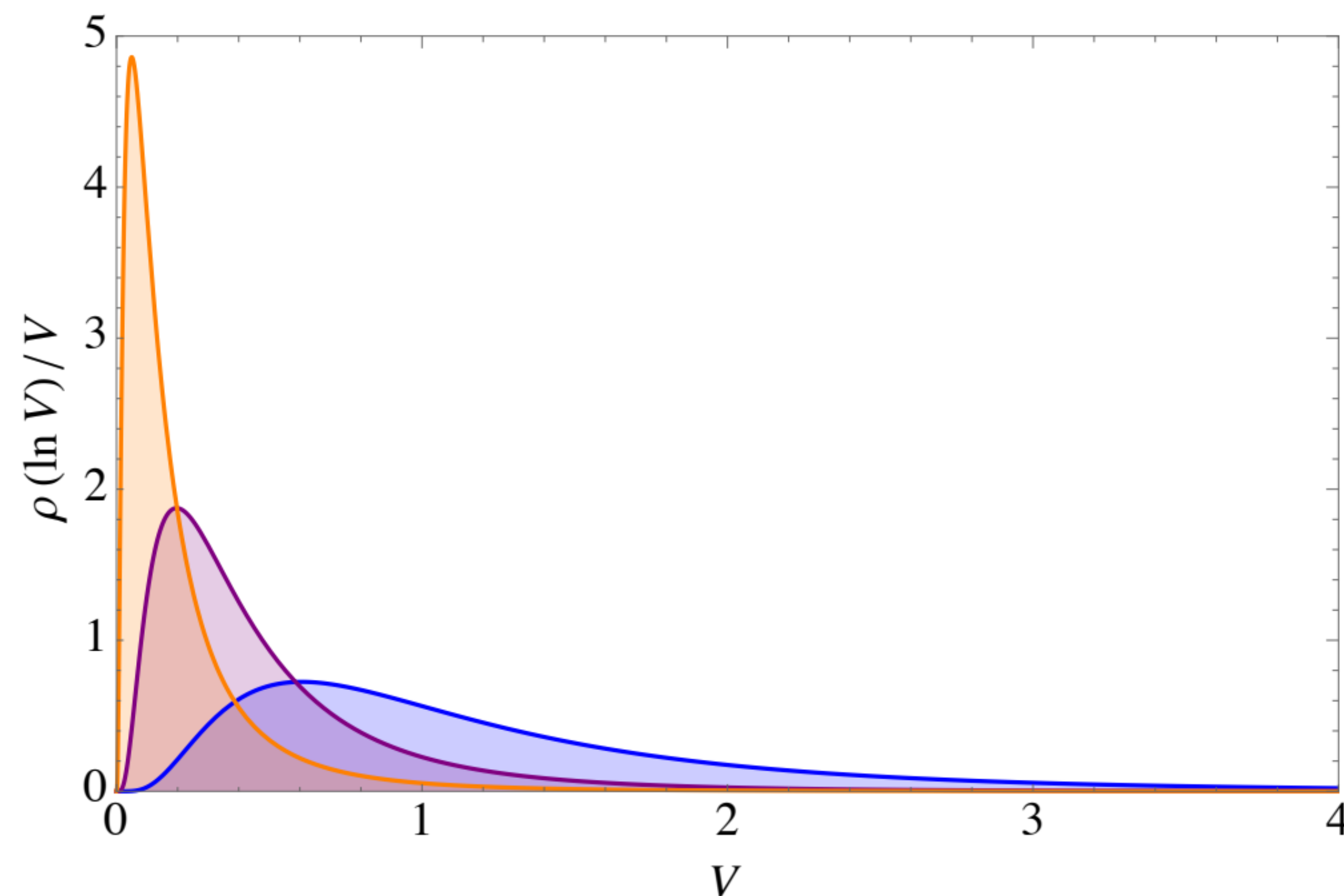
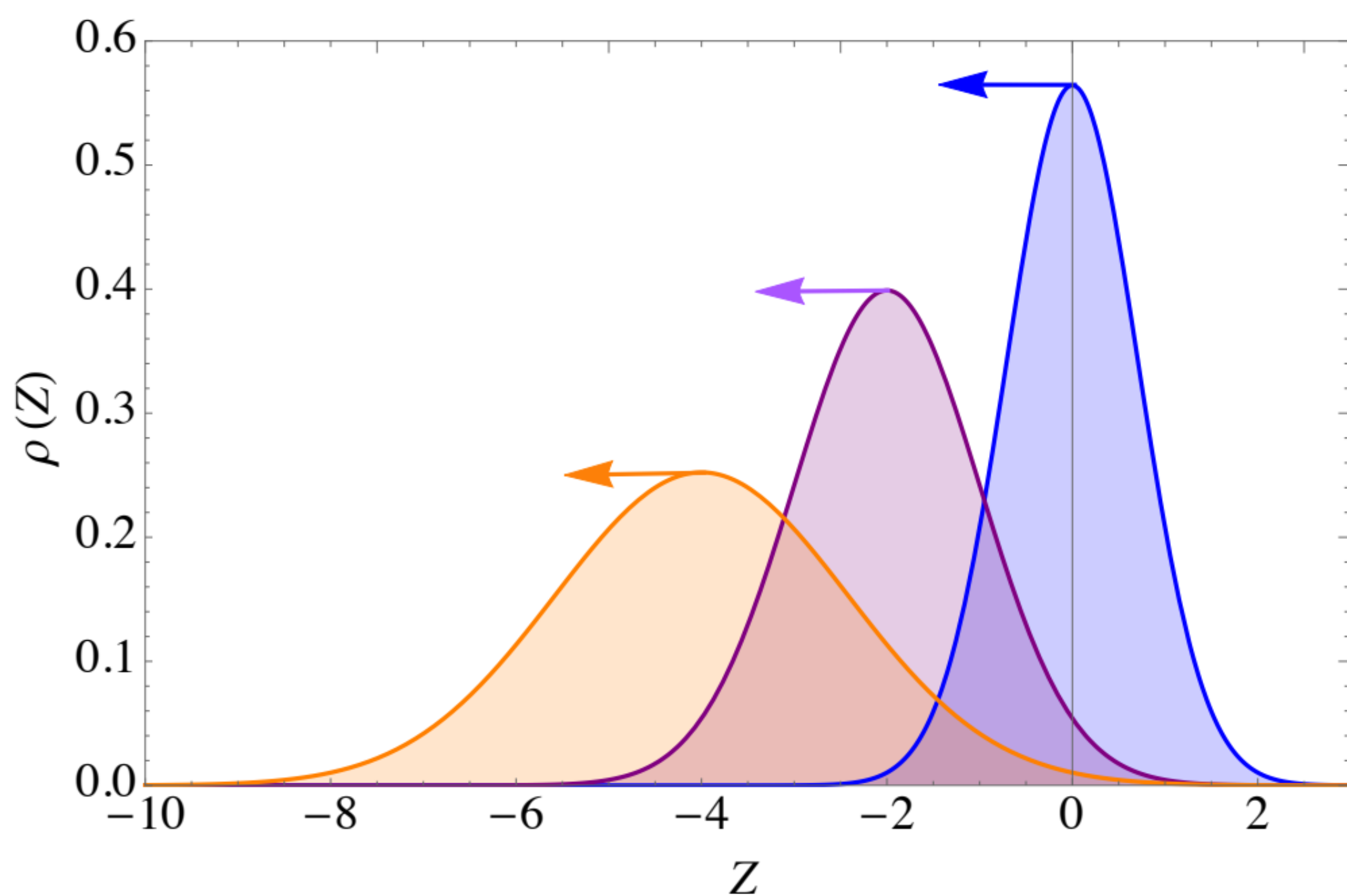
$$V = V_0 e^{(V p_V) \tau} \quad \text{and} \quad p_V = p_{V0} e^{-(V p_V) \cdot \tau}$$

symmetric ordering choice

$$H = V^2 p_V^2 \quad \mapsto \quad \hat{H} = \sqrt{V} \frac{1}{i} \partial_V \sqrt{V} \cdot \sqrt{V} \frac{1}{i} \partial_V \sqrt{V}$$

coordinate transformation $V \mapsto Z = \ln V$

→ $U \hat{H} U^{-1} = -\partial_Z^2, \quad \text{and} \quad Z \in \mathbb{R}$



slow-gauge time

$$d\theta = (V p_V) dV - \left(\frac{V^2 p_V^2}{2} \right) d \left(\frac{9p_k}{k} + \frac{V - \ln V}{V p_V} \right) \\ + d \left(\frac{9p_k}{2k} V^2 p_V^2 + \frac{1}{2} V \ln V p_V - \frac{1}{2} V^2 p_V \right)$$

→ Action

$$\mathcal{S} = \int d\theta = \int d\eta \left(V p_V \dot{V} - \frac{V^2 p_V^2}{2} \right)$$

new time variable

$$\eta \equiv \frac{9p_k}{k} + \frac{V - \ln V}{V p_V}$$

$$\dot{V} \equiv dV/d\eta$$

canonical if

$$\pi_V = p_V V$$

$$H = \frac{1}{2} V^2 p_V^2 = \frac{1}{2} \pi_V^2$$

freely moving particle...

$$(V, \pi_V) \in \mathbb{R}_+ \times \mathbb{R}$$

on the half line

Quantization: a gaussian wave packet

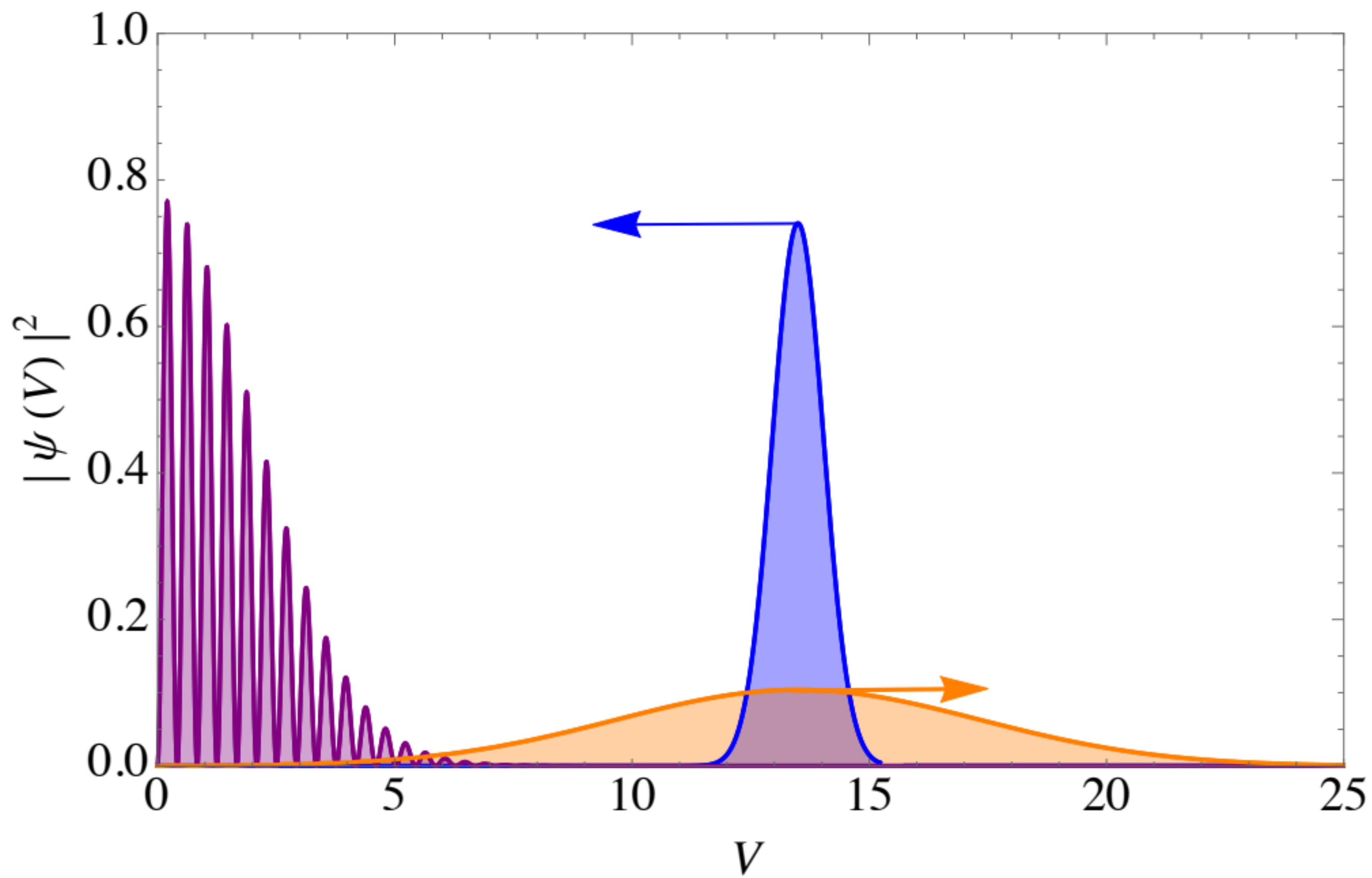
$$u(V, \eta) = \frac{e^{-k^2/4}}{\sqrt{1 + 4i\eta}} \exp \left[-\frac{(V - ik/2)^2}{1 + 4i\eta} \right]$$

implement boundary conditions to ensure self-adjointness

$$\psi(V, \eta) = \frac{u(V + V_0, \eta) - u(-V + V_0, \eta)}{\left[\sqrt{\pi/2} (1 - e^{-V_0^2 - k^2/2}) \right]^{1/2}}$$

→ *solves the Schrödinger equation*

$$i \frac{\partial}{\partial \eta} \psi = - \Delta_D \psi$$

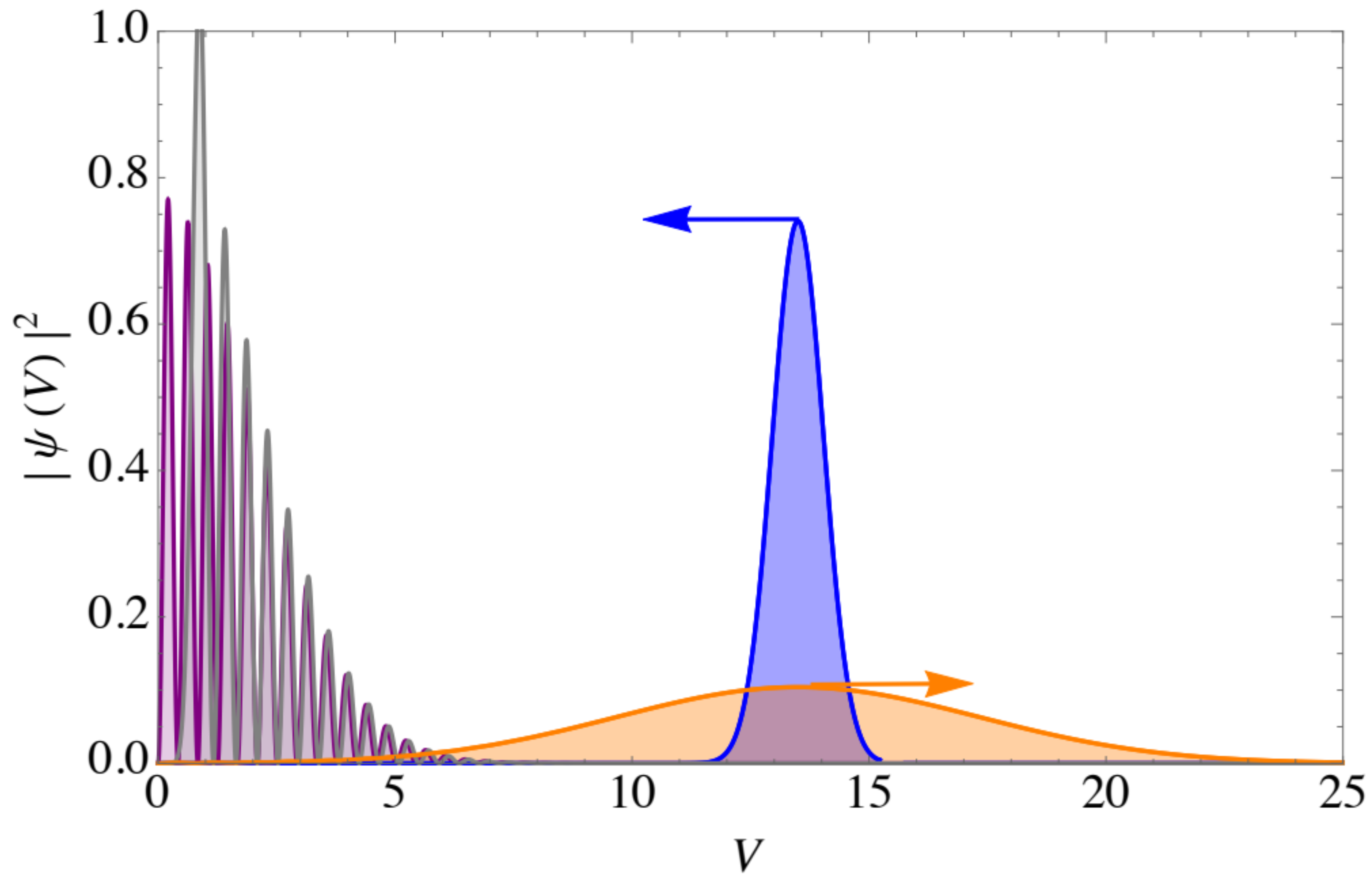


Operator ordering ambiguity

$$\pi_V^2 \mapsto \hat{V}^s \hat{\pi}_V \hat{V}^{-2s} \hat{\pi}_V \hat{V}^s$$

→ $\pi_V^2 \mapsto \hat{\pi}_V^2 + s\hat{V}^{-2}$

↓
self-adjoint hamiltonian on the half-line $s > 3/4$



Closed algebra of operators

$$\begin{cases} [\hat{V}^2, \hat{H}] = 4i\hat{D}, \\ [\hat{D}, \hat{H}] = 2i\hat{H}, \\ [\hat{V}^2, \hat{D}] = 2i\hat{V}^2, \end{cases}$$

$$\hat{D} \equiv \frac{1}{2} (\hat{V} \hat{\pi}_V + \hat{\pi}_V \hat{V})$$

→ Heisenberg equations of motion

$$\frac{d}{d\eta} \hat{V}^2 = -i[\hat{V}^2, \hat{H}] = 4\hat{D}$$

$$\frac{d}{d\eta} \hat{D} = -i[\hat{D}, \hat{H}] = 2\hat{H}$$

→ Heisenberg equations of motion

$$\frac{d}{d\eta} \hat{V}^2 = -i[\hat{V}^2, \hat{H}] = 4\hat{D}$$

$$\frac{d}{d\eta} \hat{D} = -i[\hat{D}, \hat{H}] = 2\hat{H}$$

→ solution as time-dependent operators

$$\hat{D}(\eta) = 2\hat{H}\eta + \hat{D}(0) \quad \longrightarrow \quad \hat{V}^2 = 4\hat{\eta}^2 + 4\hat{D}(0)\eta + \hat{V}^2(0)$$

expectation values follows similar equations...

→ *semi-classical variables*

$$\check{V}(t) = \sqrt{\langle \hat{V}^2(t) \rangle}$$

$$\check{\pi}_V(t) = \frac{\langle \hat{D}(t) \rangle}{\check{V}(t)}$$

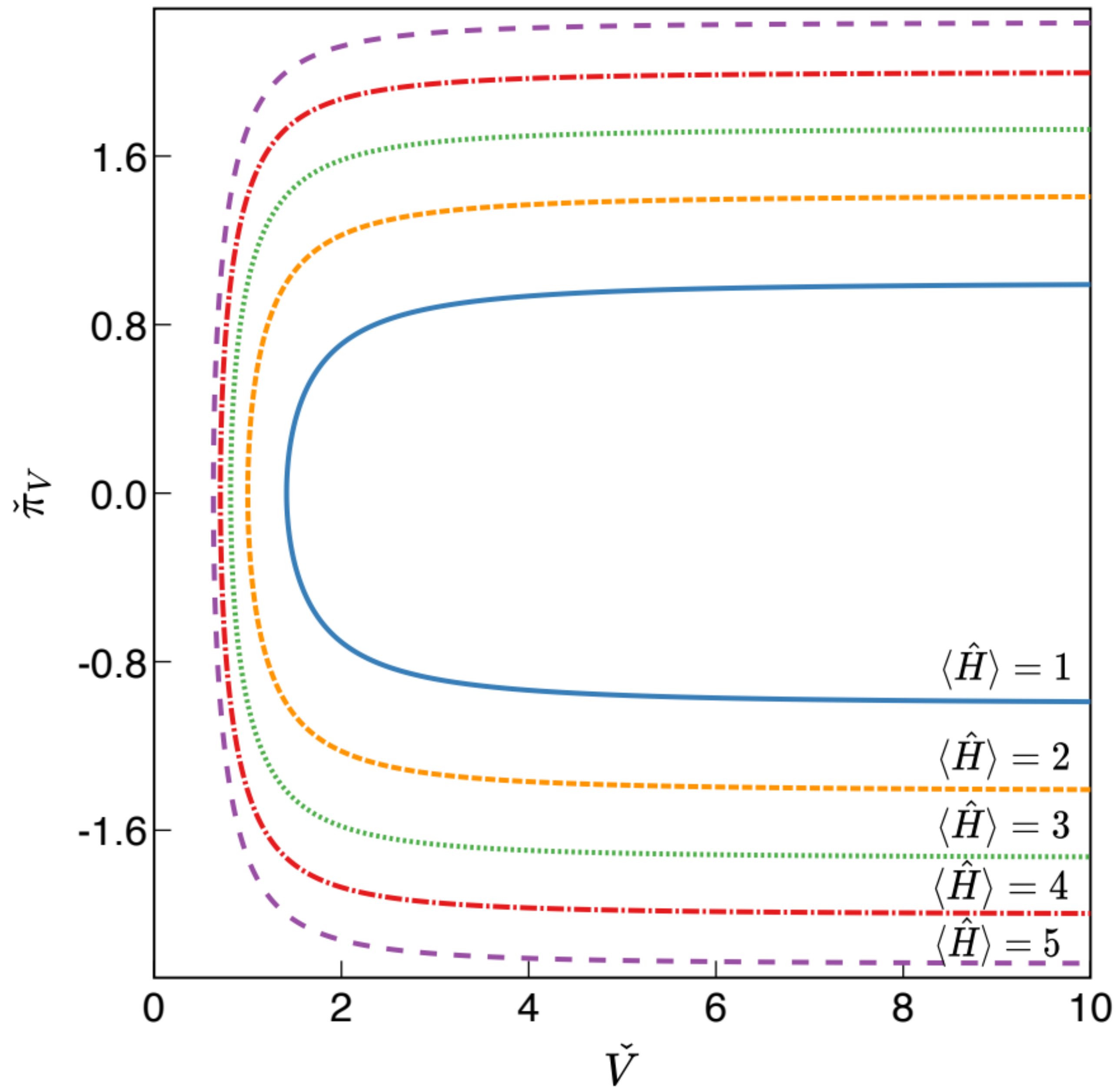
→ *phase space solution*

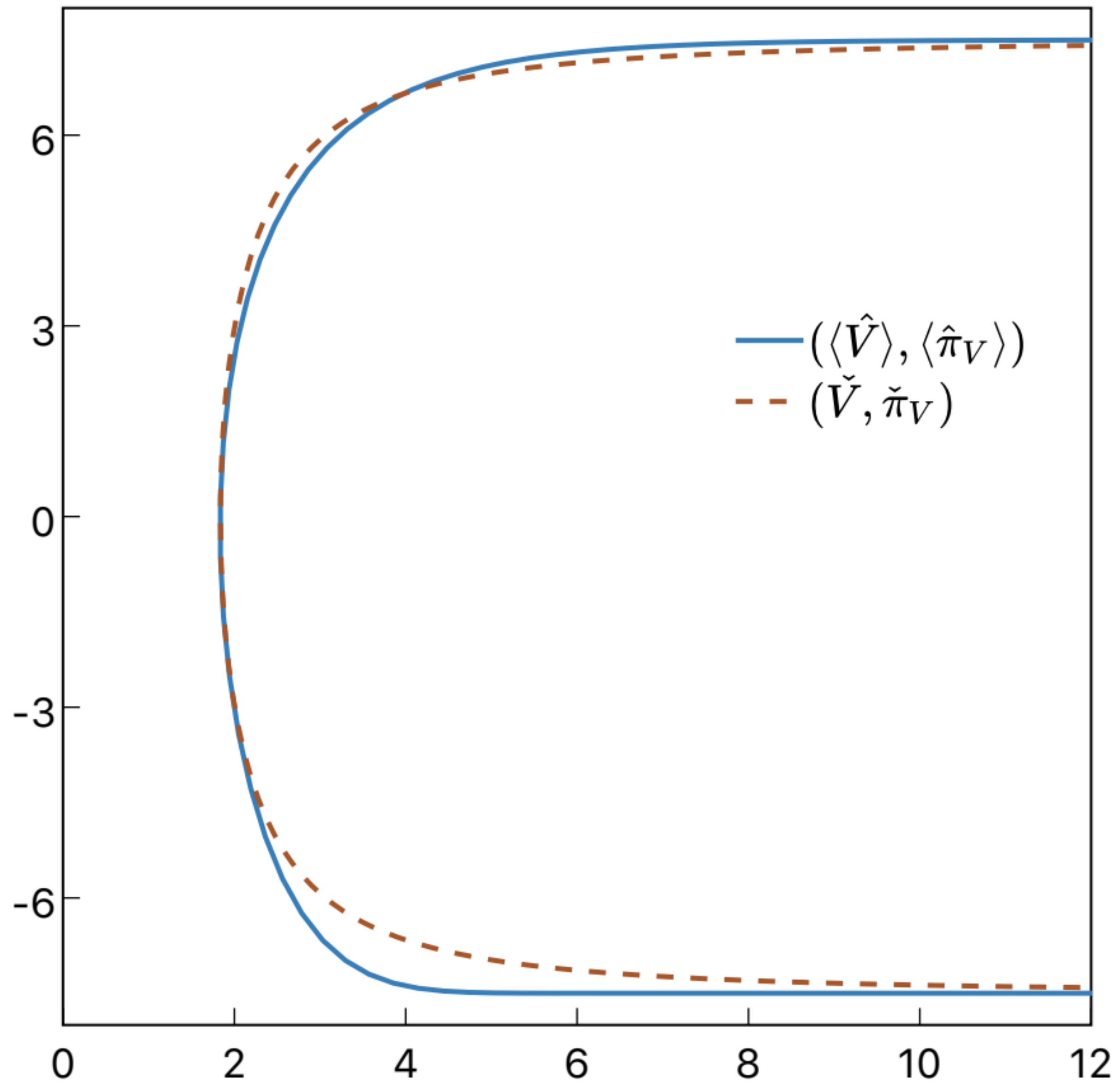
$$\check{V}(t) = \sqrt{4\langle\hat{H}\rangle t^2 + V_0^2},$$

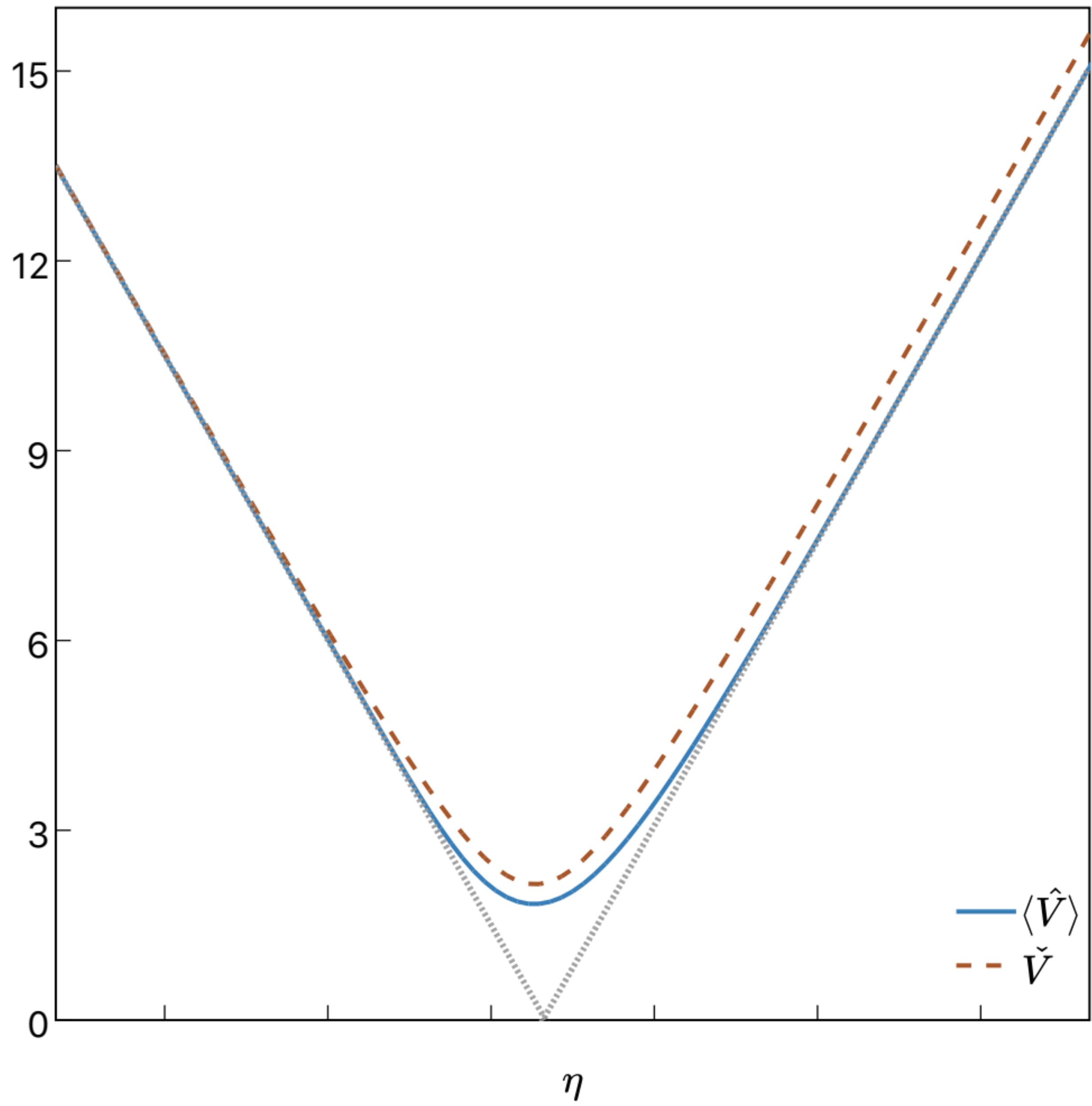
$$\check{\pi}_V(t) = \frac{2\langle\hat{H}\rangle t}{\sqrt{4\langle\hat{H}\rangle t^2 + V_0^2}}.$$



NO SINGULARITY







Changing the time variable $\eta' = \eta'(\eta, V, \pi_V)$

redefining the dynamical variables in the process

$$\pi'_V = \pi_V \quad \& \quad V' = V + \pi_V(\eta' - \eta) \quad \text{no change of range...}$$

change the canonical one-form

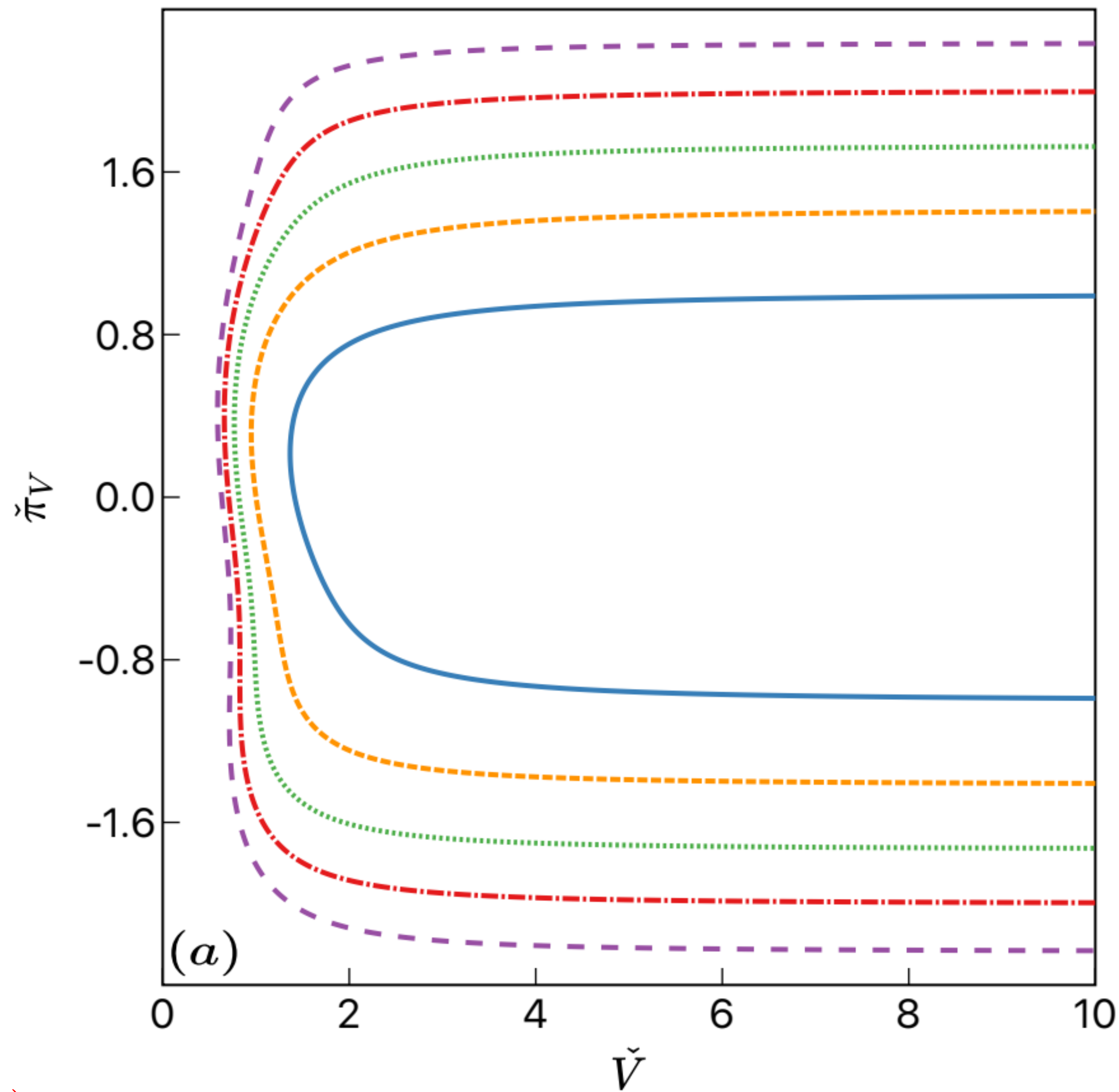
$$d\theta = \pi_V dV - \frac{\pi_V^2}{2} d\eta = \pi'_V dV' - \frac{\pi_V'^2}{2} d\eta' + d \left[(\eta - \eta') \frac{\pi_V'^2}{2} \right]$$

→ same system!

delay function $\Delta(V, \pi_V) = \eta' - \eta$ no dependency on time

Delay function

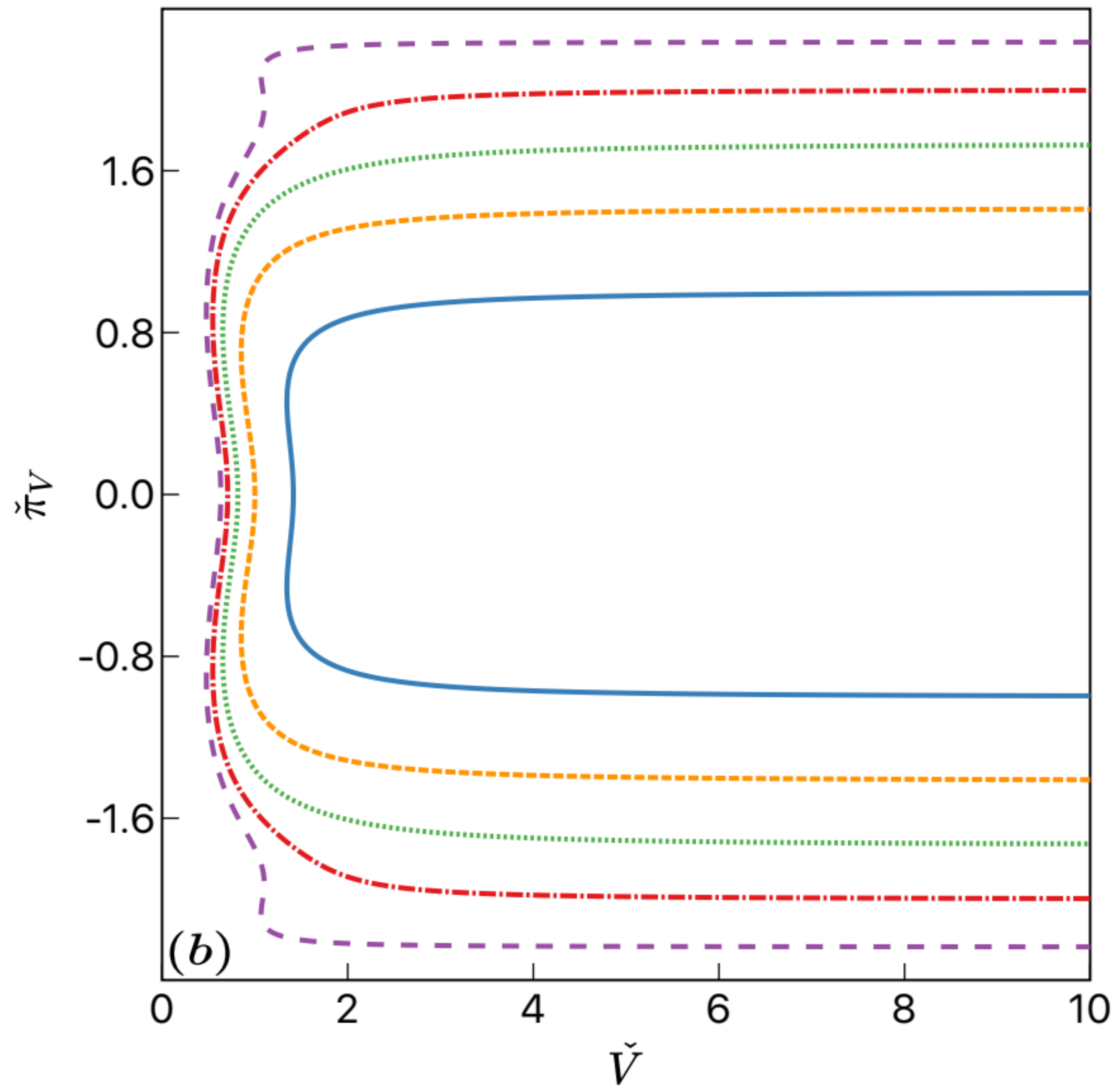
$$\Delta = V e^{-2|\pi_V|/3} \sin(3V\pi_V)/(10\pi_V)$$



(a)

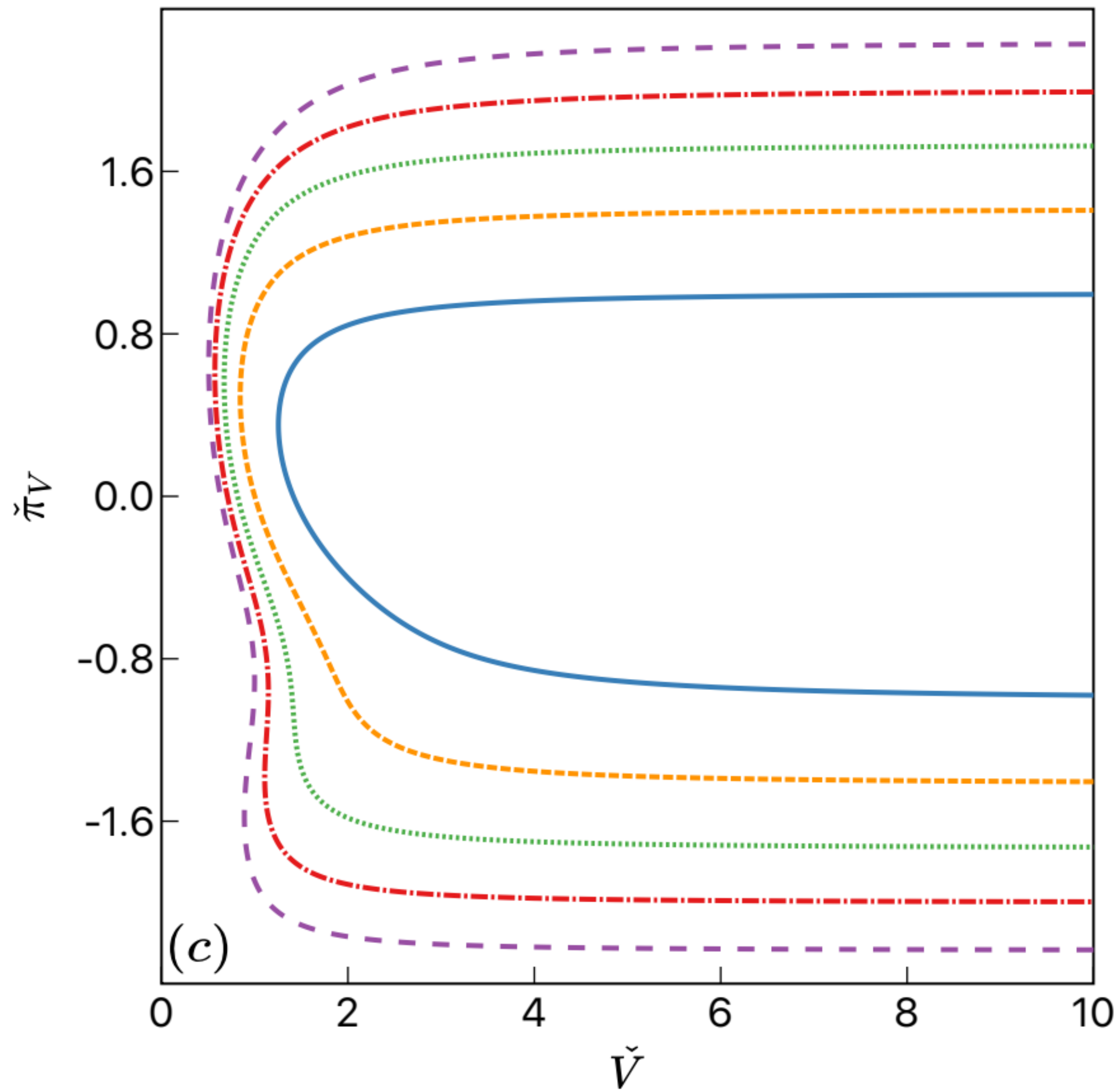
Delay function

$$\Delta = V(\pi_V - 10^{-0.2}\pi_V^3 + \pi_V^5/10)$$



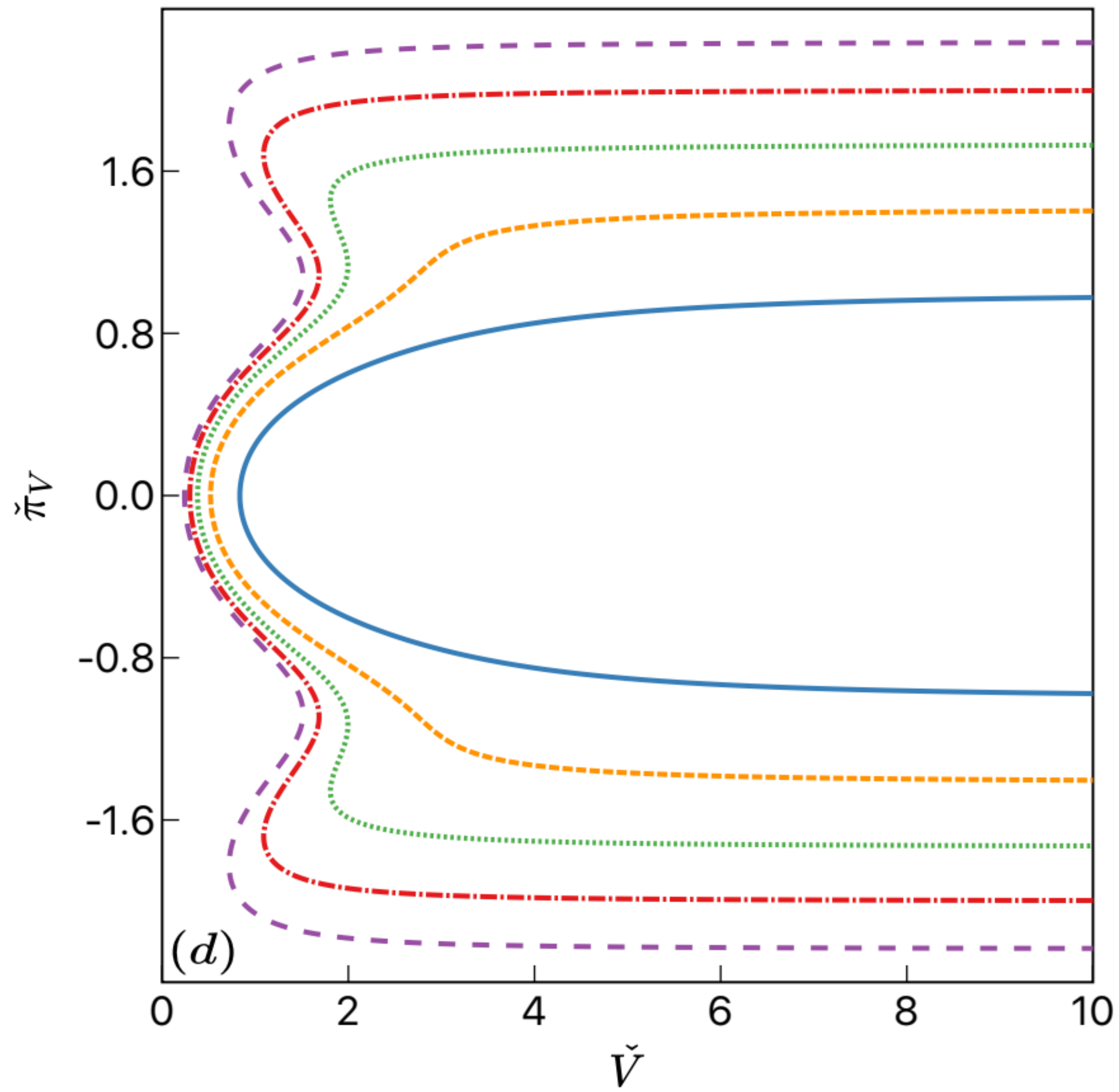
Delay function

$$\Delta = 10^{-0.5} V \sin(2\pi V) / \pi V$$



Delay function

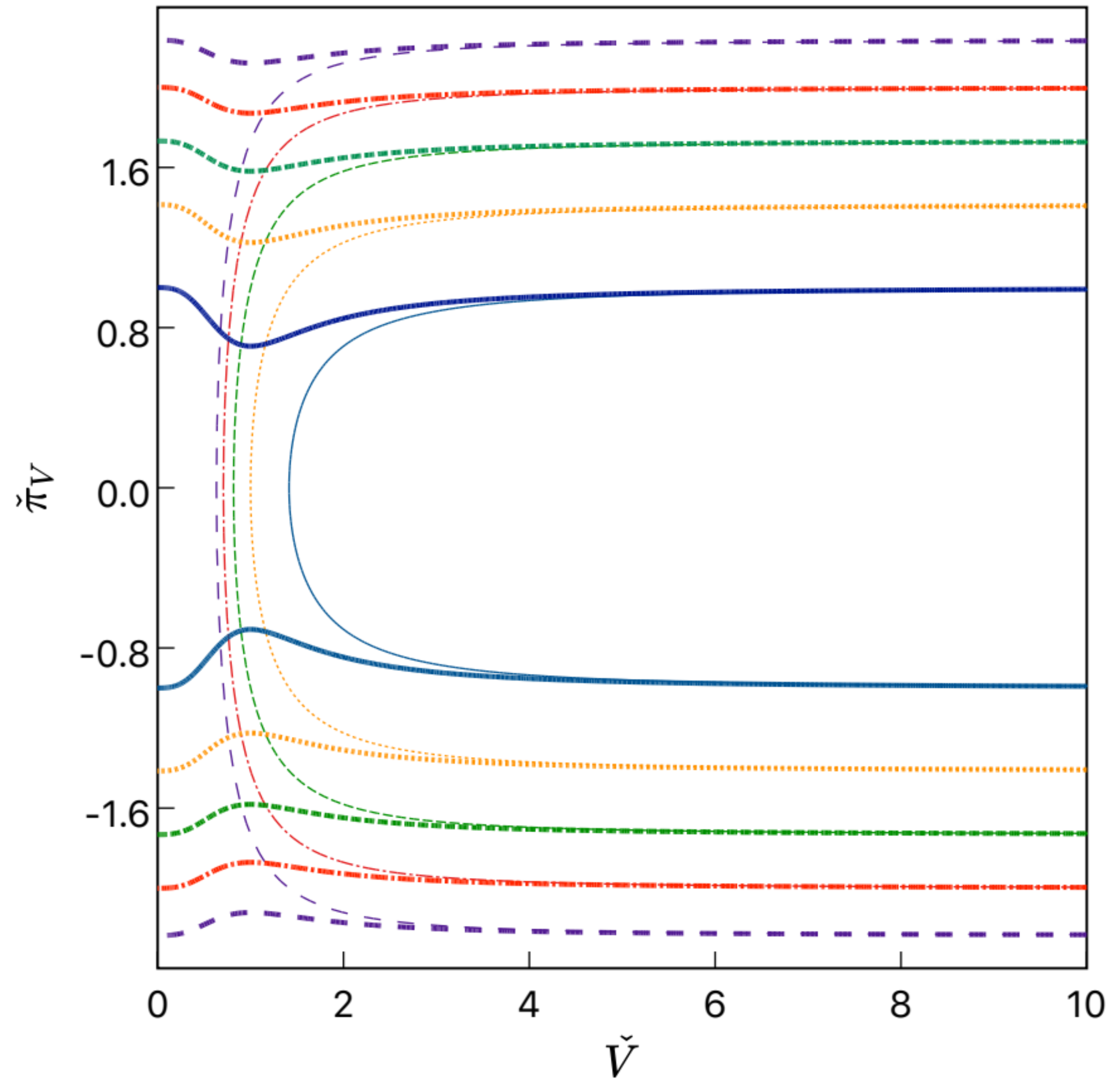
$$\Delta = 10^{-0.5} (V + 1) \cos(3\pi_V) / \pi_V$$



Delay function (slow to fast)

[regular to singular]

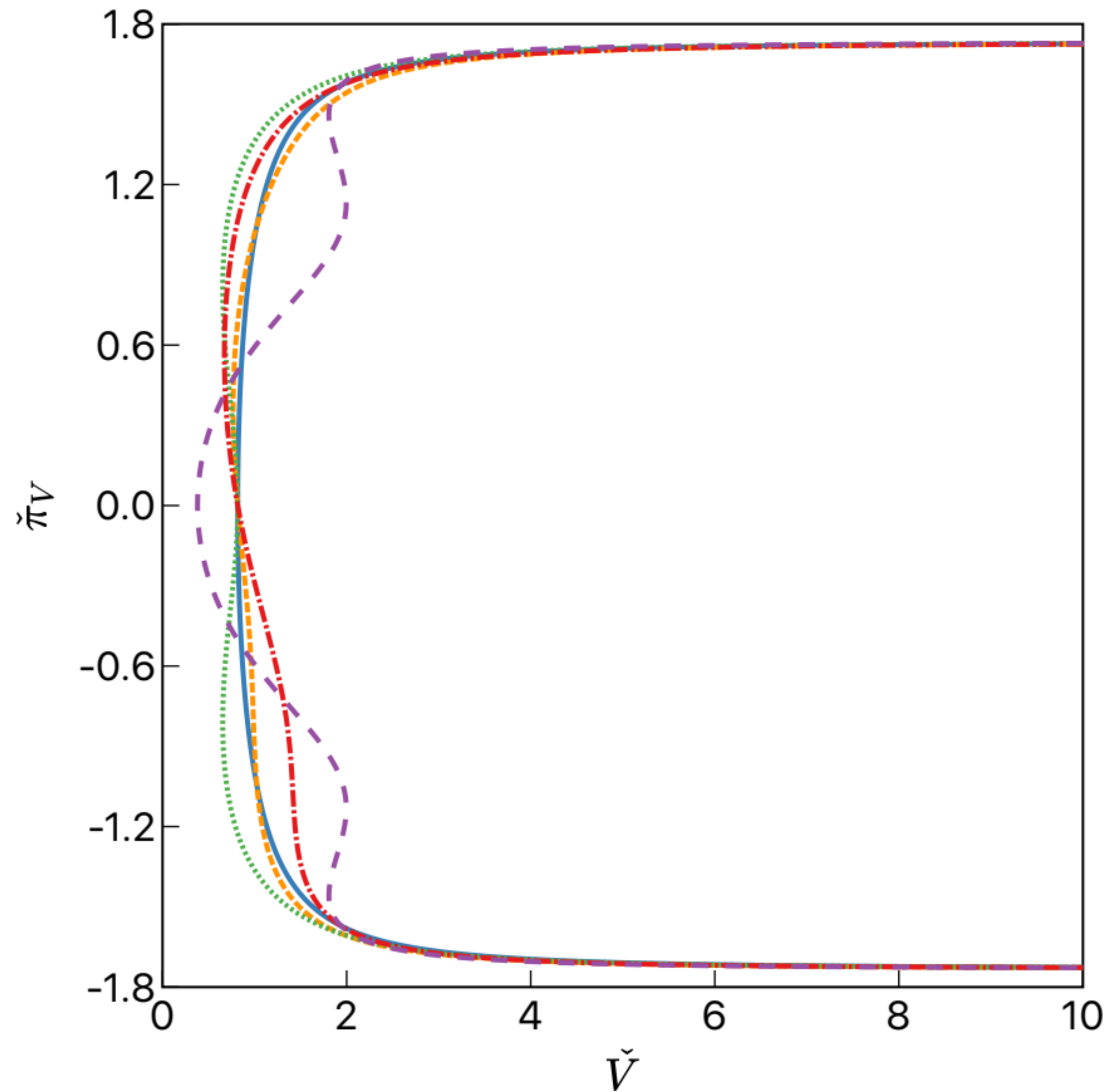
$$\Delta_{\text{slow} \rightarrow \text{fast}} = \frac{V - \ln V}{V p_V}$$



Comparison between
different delay functions



Same asymptotics



Conclusions

- *Bouncing alternative to inflation still alive*
- *Shear issue: ekpyrosis or shear viscosity*
- *Classical / quantum bounce*
- *Observational consequences & perturbations*
(talks by Shiv Sethi, Yi Fu Cai...)

Thank you for your attention!