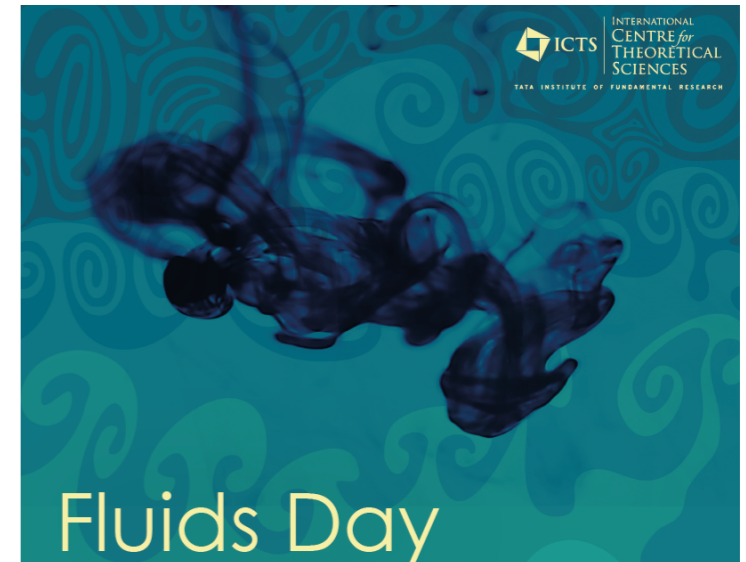




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# Lagrangian statistics in rotating turbulent flows

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FLUIDS DAY, 20th January 2020

In collaboration with  
Rama Govindarajan & Samriddhi Sankar Ray

# Turbulent particle transport



# Time irreversibility

$$\partial_t \mathbf{u} + ((\mathbf{u} \cdot \nabla) \mathbf{u}) = -\nabla P + \nu \nabla^2 \mathbf{u} + \mathbf{f} \quad \dots (1),$$

$$\nabla \cdot \mathbf{u} = 0 \quad \dots (2),$$

$$t \rightarrow t', \quad \mathbf{u} \rightarrow -\mathbf{u}$$

**Time reversal symmetry is broken**

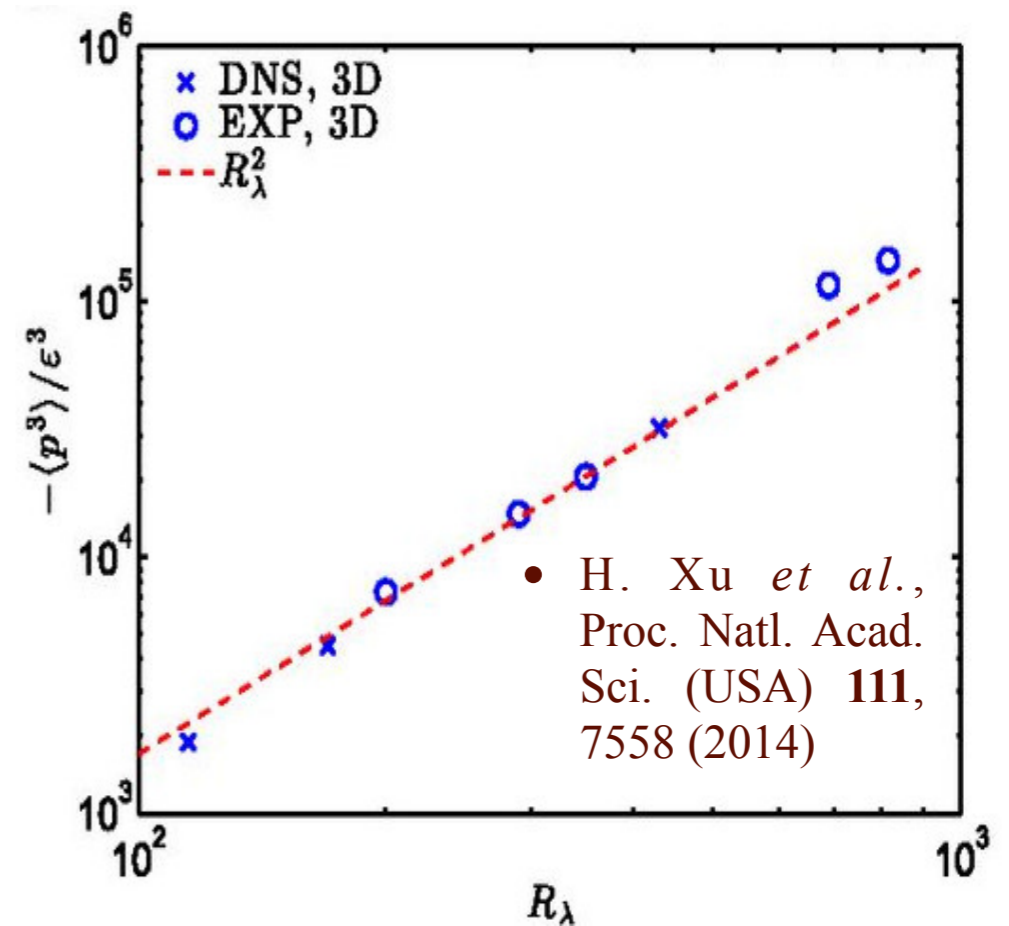
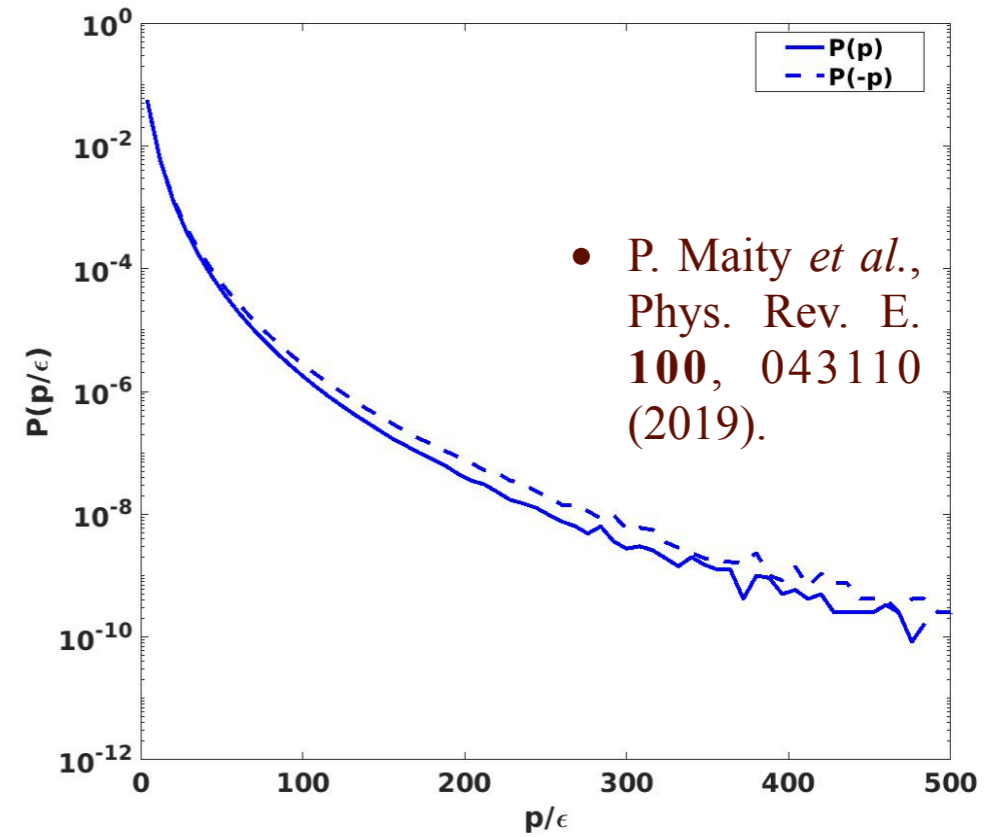
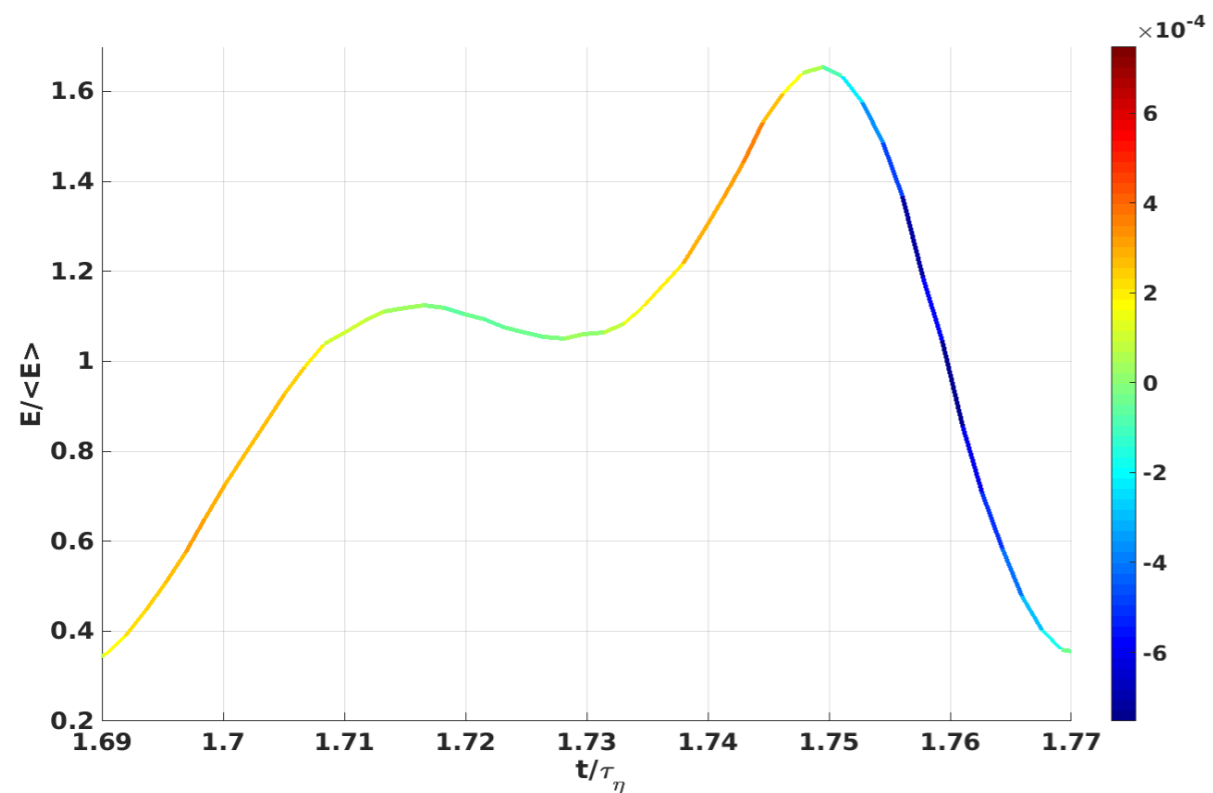
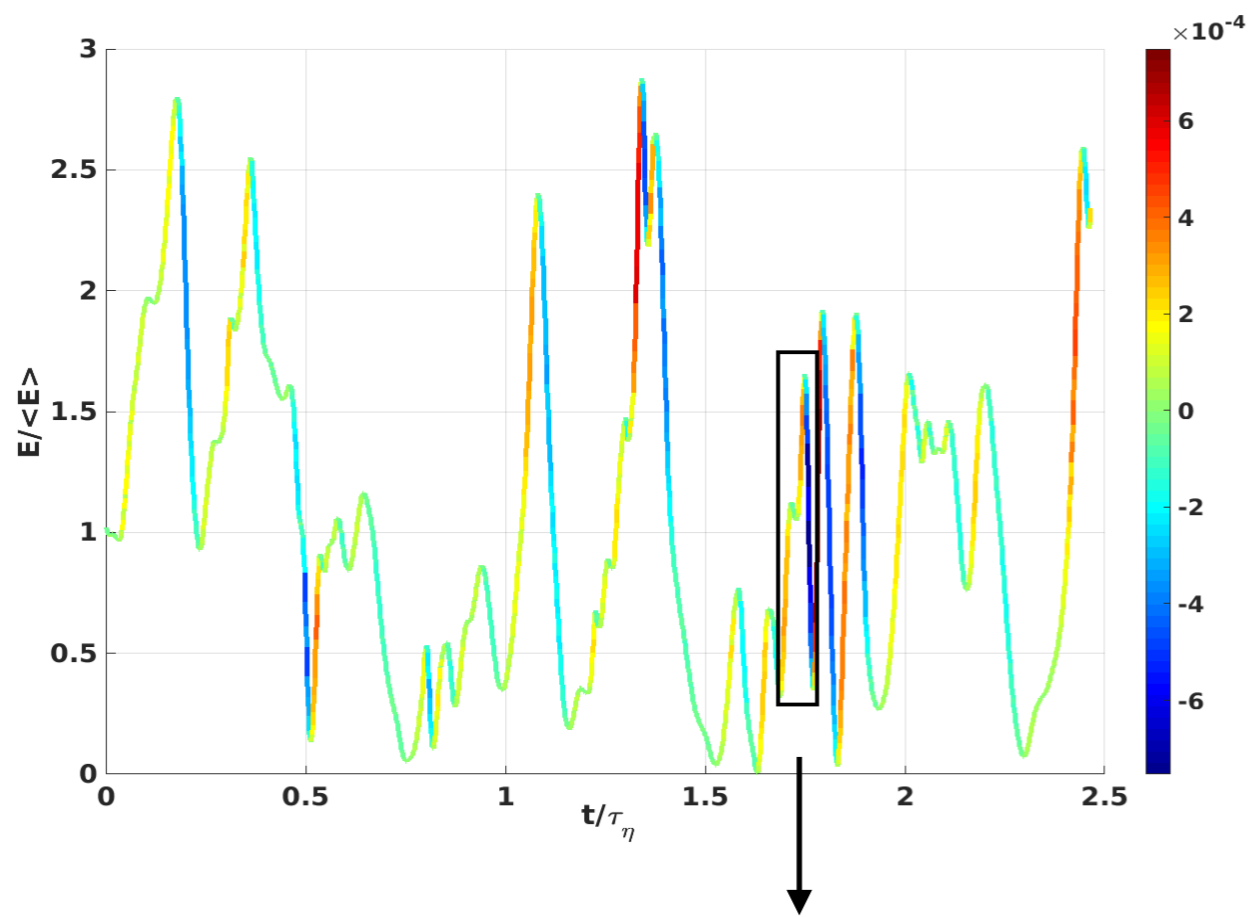
## Eulerian Framework

1. M. K. Verma, “Asymmetric energy transfers in driven nonequilibrium systems and arrow of time”, Eur. Phys. J. B **92**, 190 (2019).

## Lagrangian Framework

1. H. Xu *et al.*, Proc. Natl. Acad. Sci. (USA) **111**, 7558 (2014).
2. J. Jucha, H. Xu, A. Pumar, Phys. Rev. Lett. **113**, 054501 (2014).
3. M. Cencini, L. Biferale, G. Boffetta, M. De Pietro, Phys. Rev. Fluids **2**, 104604 (2017).
4. A. Bhatnagar, A. Gupta, D. Mitra, and R. Pandit, Phys. Rev. E. **97**, 033102 (2018).

# Particle trajectories of tracers in turbulent flow





# Rotating turbulence

## Eulerian fluid fields

$$\partial_t \mathbf{u} + ((\mathbf{u} \cdot \nabla) \mathbf{u}) + 2(\boldsymbol{\Omega} \times \mathbf{u}) = -\nabla P + \nu \nabla^2 \mathbf{u} + \mathbf{f} + \alpha \Delta^{-1} \mathbf{u} \quad \dots (3),$$

$$\nabla \cdot \mathbf{u} = 0 \quad \dots (4),$$

Coriolis force
Large scale friction

$$P = p - \frac{1}{2} |\boldsymbol{\Omega} \times \mathbf{r}|^2 \quad \dots (5),$$

Centrifugal force

Where  $\mathbf{u}(x,y,z,t) = (u_1, u_2, u_3)$  is the flow velocity,  $\nu$  is the kinematic viscosity, and  $P$  is the pressure field. The external forcing term  $\mathbf{f}$  is active for  $k \leq 3$ .

## Lagrangian Particle

$$Re = \frac{uL}{\nu} \quad \frac{d\mathbf{r}_p}{dt} = \mathbf{v}_p \quad \dots (6)$$

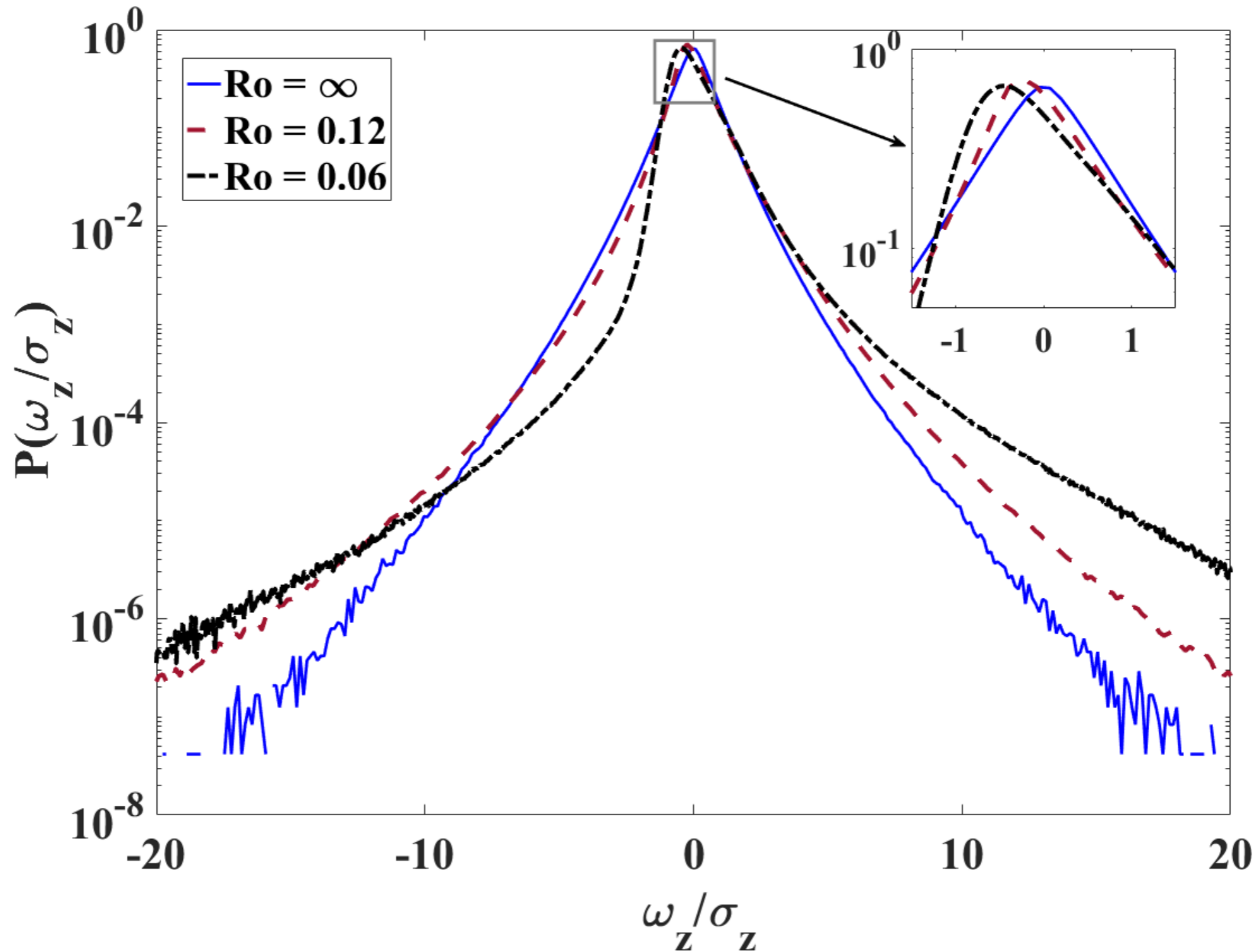
$$St = \frac{\tau_p}{\tau_\eta} \quad \frac{d\mathbf{v}_p}{dt} = \frac{\mathbf{v}_p - \mathbf{u}_p}{\tau_p} - 2\boldsymbol{\Omega} \times \mathbf{v}_p - (\boldsymbol{\Omega} \times \boldsymbol{\Omega} \times \mathbf{r}_p) \quad \dots (7)$$

$$Ro = \frac{u}{L\Omega} \quad \tau_p \rightarrow 0, \quad \mathbf{v}_p = \mathbf{u}_p \quad \dots (8)$$

$\mathbf{u}_p$  is the flow velocity at the particle position,  $\mathbf{v}_p$  and  $\mathbf{r}_p$  are the particle velocity and position vector respectively.

1. Luca Biferale *et. al.*, "Coherent Structures and Extreme Events in Rotating Multiphase Turbulent Flows", Phys. Rev. E **6**, 041036 (2016).
2. L. Del Castello and H. J. H. Clercx, Phys. Rev. Lett. **107**, 214502 (2011).
3. L. Del Castello and H. J. H. Clercx, Phys. Rev. E **83**, 056316 (2011)

# Rotational effects



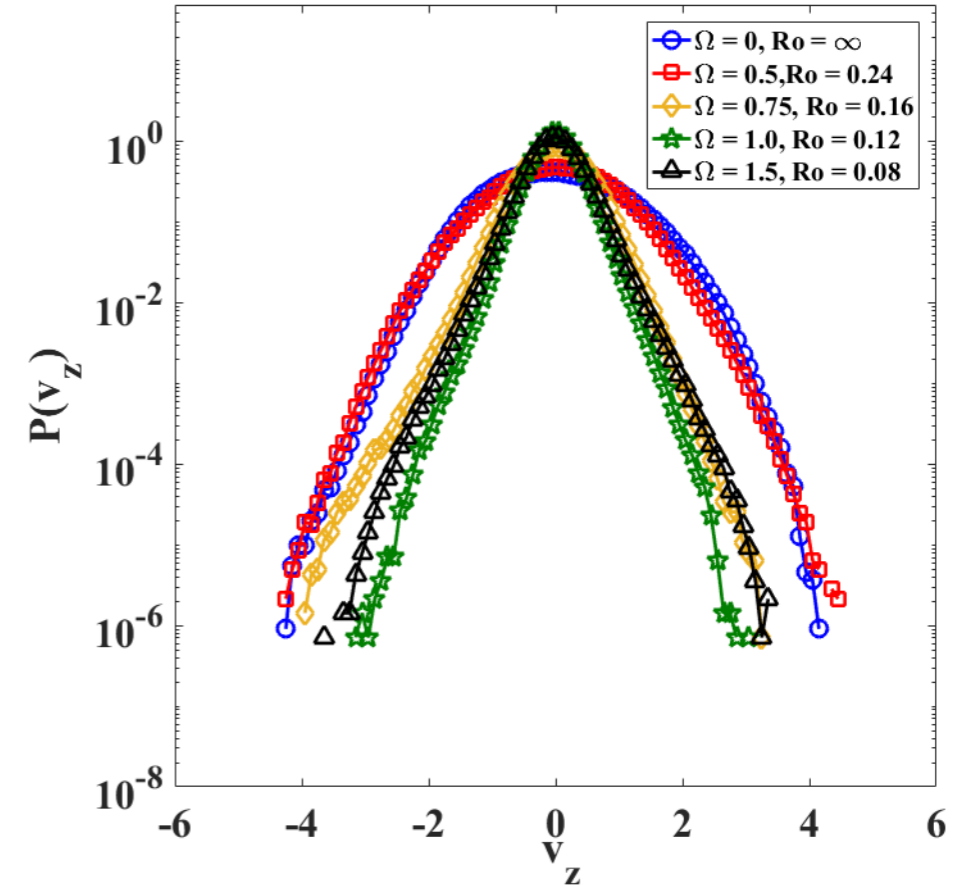
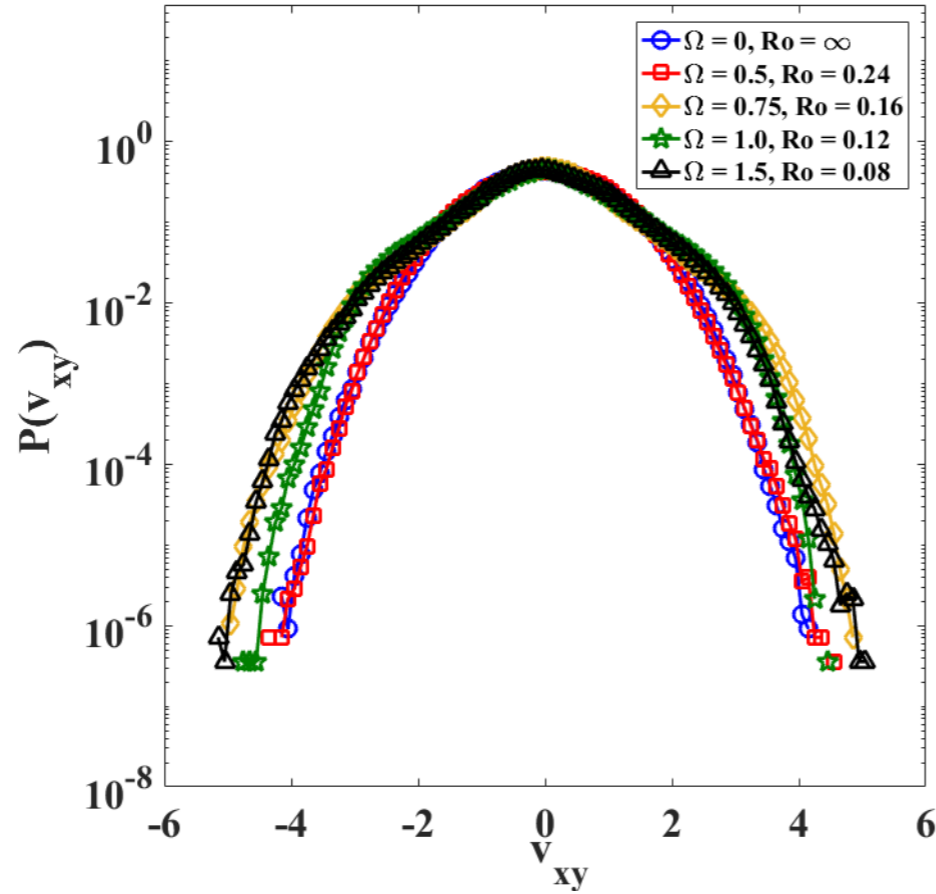
- **P. Maity**, R. Govindarajan, and S. S. Ray, “*Statistics of Lagrangian trajectories in a rotating turbulent flow*”, Phys. Rev. E. **100**, 043110 (2019).
- C. Morize, F. Moisy, and M. Rabaud, Phys. Fluids **17**, 095105 (2005).

# Anisotropy due to rotation

## Simulations

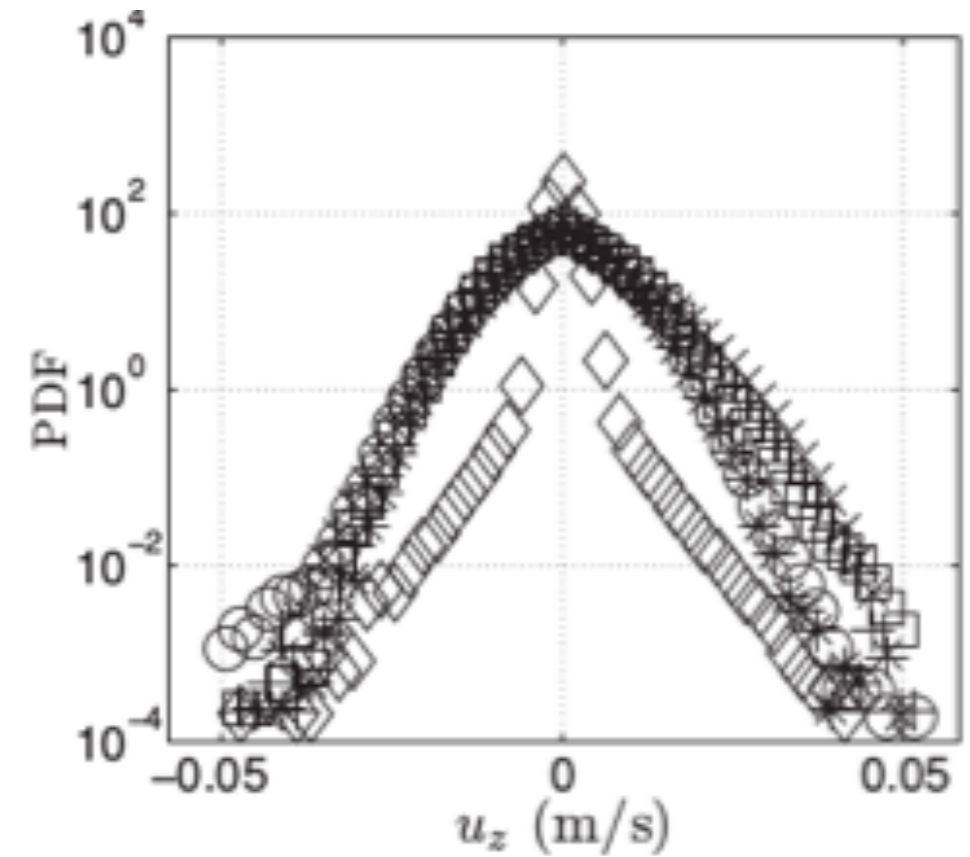
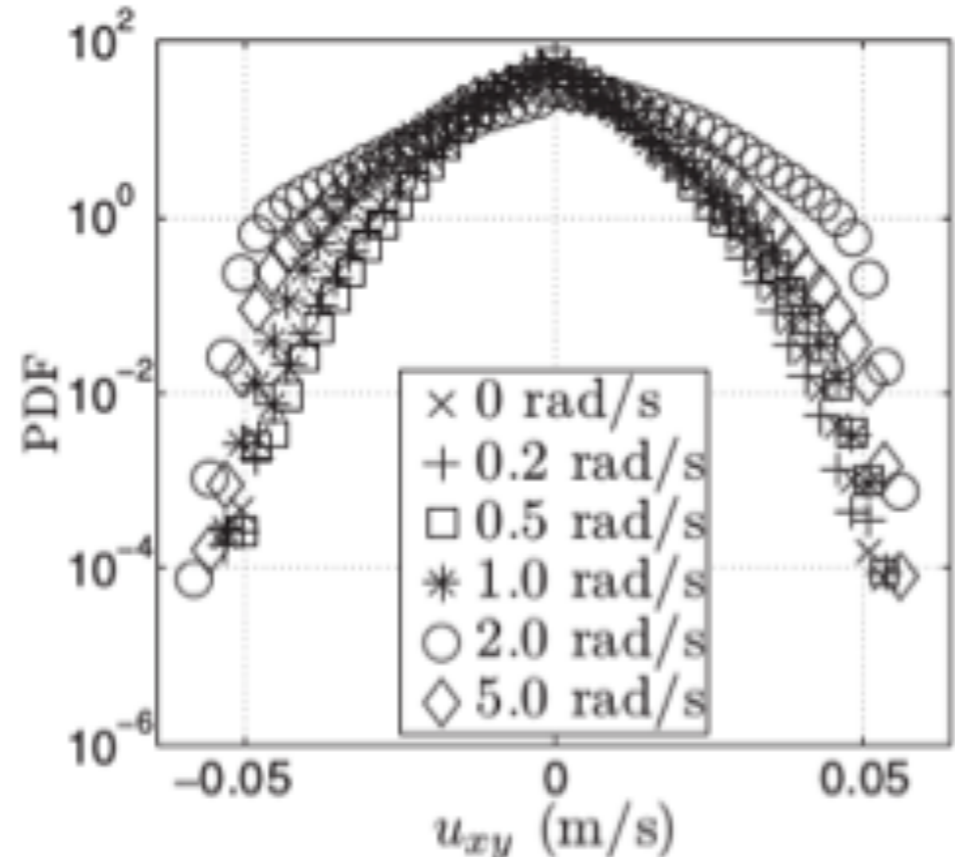
## Results

Maity et. al 2020.



## Experimental results

- L. D. Castello & H. J. H. Clercx, PRE **83**, 056316 (2011)



# Simulation details

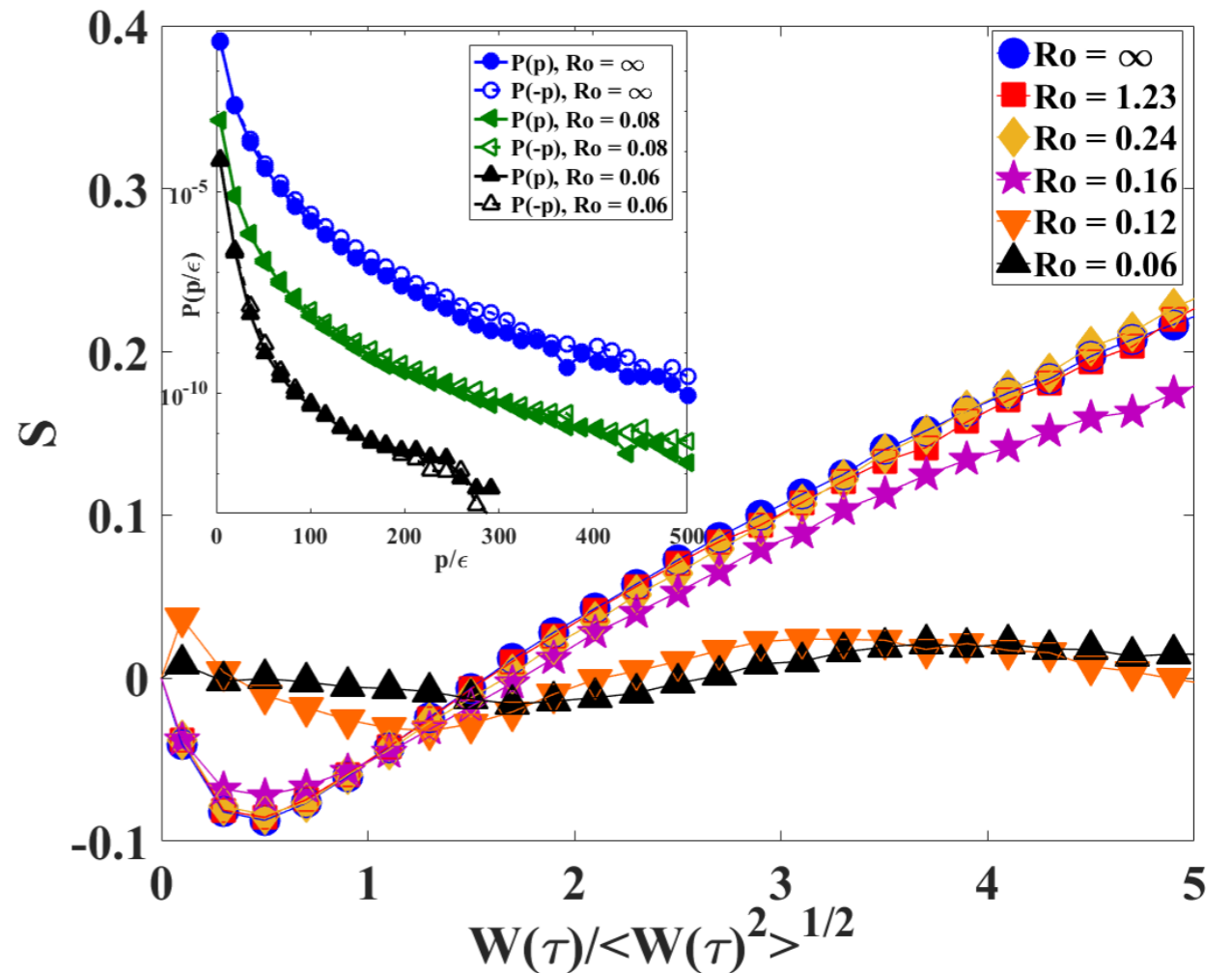
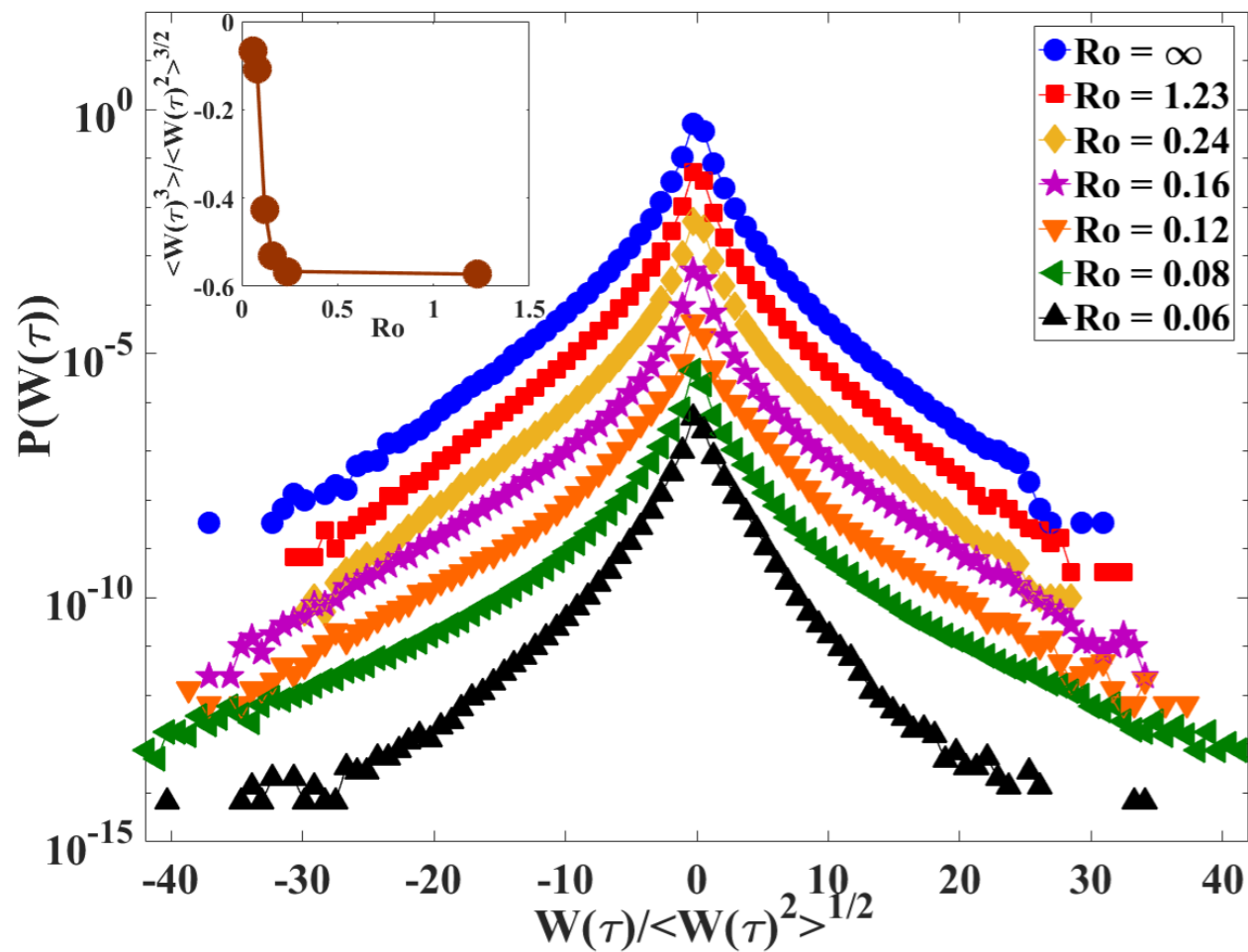
<b>N Coallocation points</b>	$\nu$ <b>Kinema tic viscosity</b>	$\delta t$ <b>Time step</b>	$N_p$ <b>Number of tracers</b>	$Re_\lambda$ <b>Taylor Reynold's number</b>	$k_{max}\eta$	$\tau_\eta$ <b>Kolmogorov time</b>	$\alpha$ <b>Coefficient of large scale friction</b>	$\varepsilon$ <b>Mean energy dissipation</b>
512	$10^{-3}$	$4 \times 10^{-4}$	$10^6$	90	2.56	$3.33 \times 10^{-3}$	$5 \times 10^{-3}$	0.89

$$\mathbf{\Omega}=(0,0,\Omega)$$

$\Omega$	<b>Rossby number <math>Ro=U/L\Omega</math></b>
0	$\infty$
0.1	1.43
0.5	0.24
0.75	0.16
1.0	0.12
1.5	0.08
2.0	0.06



# Effect of rotation on energy increment statistics



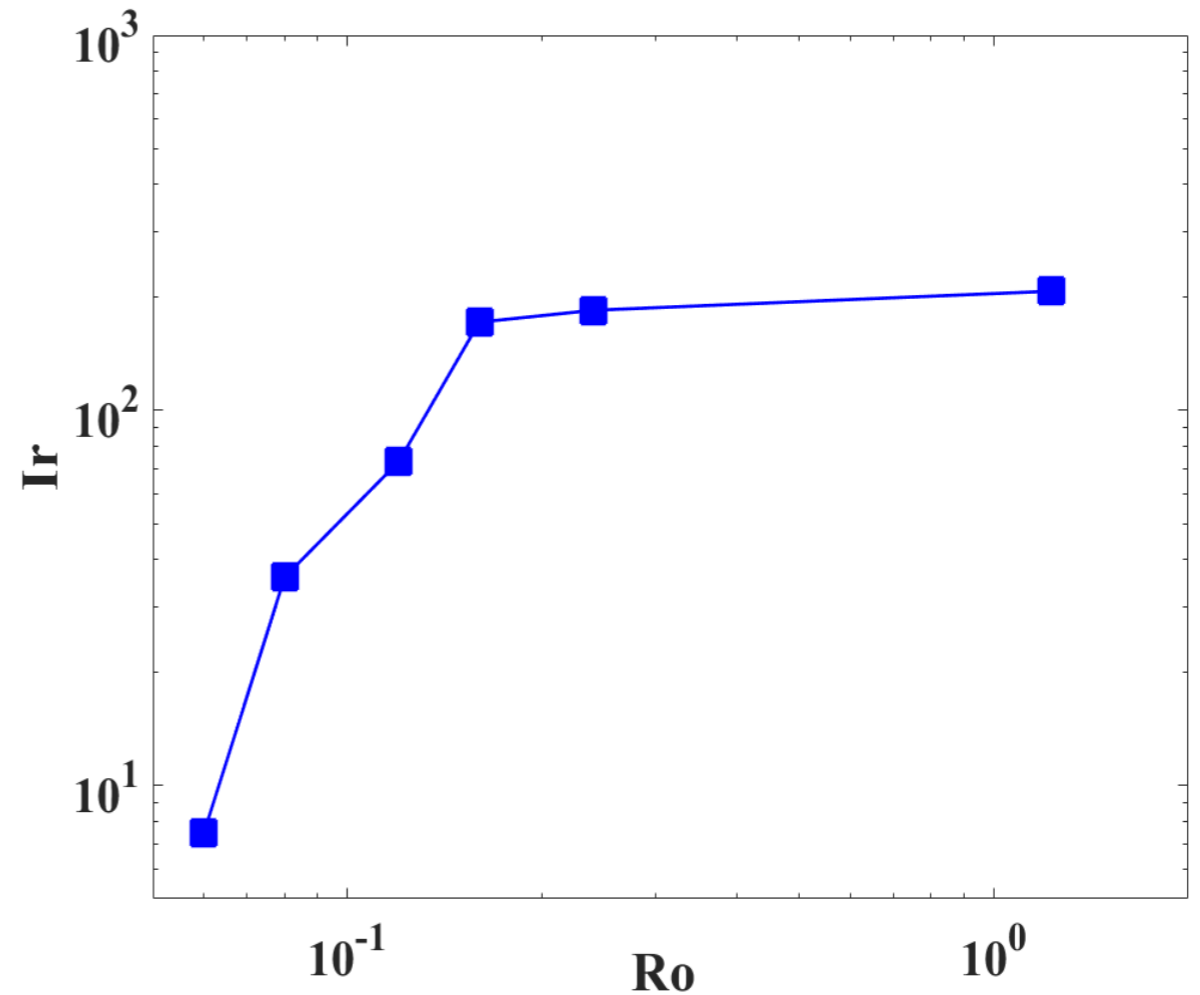
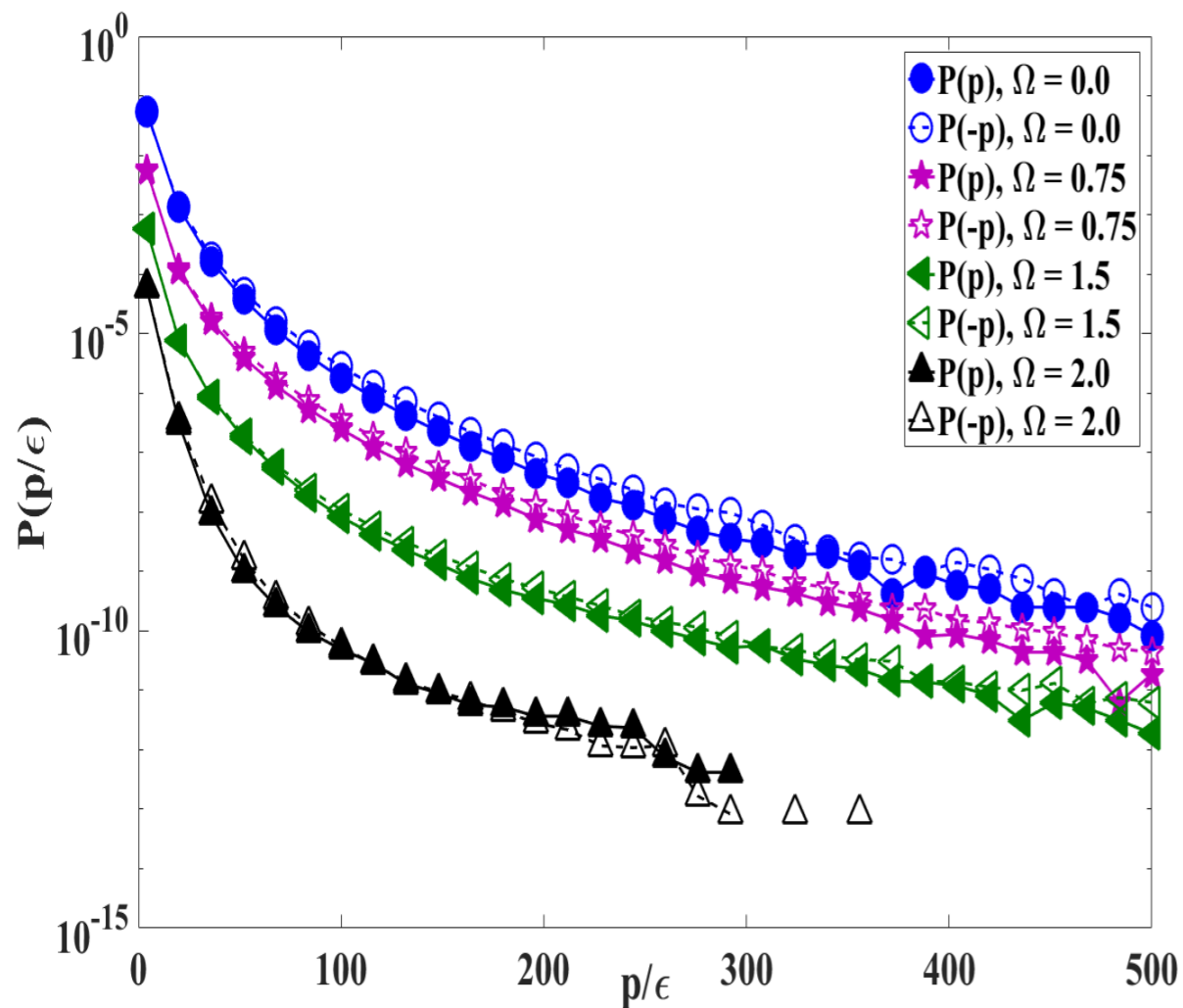
PDFs of energy increment  $W(\tau) = E(t+\tau) - E(t)$ , for time increments  $\tau/\tau_\eta = 2$ ,  $\tau_\eta$  being the Kolmogorov time scale.

The PDFs have been shifted by a factor of 10 for clarity.

Variation of symmetry function  $S = \ln(P(-W)/P(W))$  as a function of  $W(\tau)$ ,  $\tau/\tau_\eta = 2$  for various values of rotation rates given by  $Ro$ .

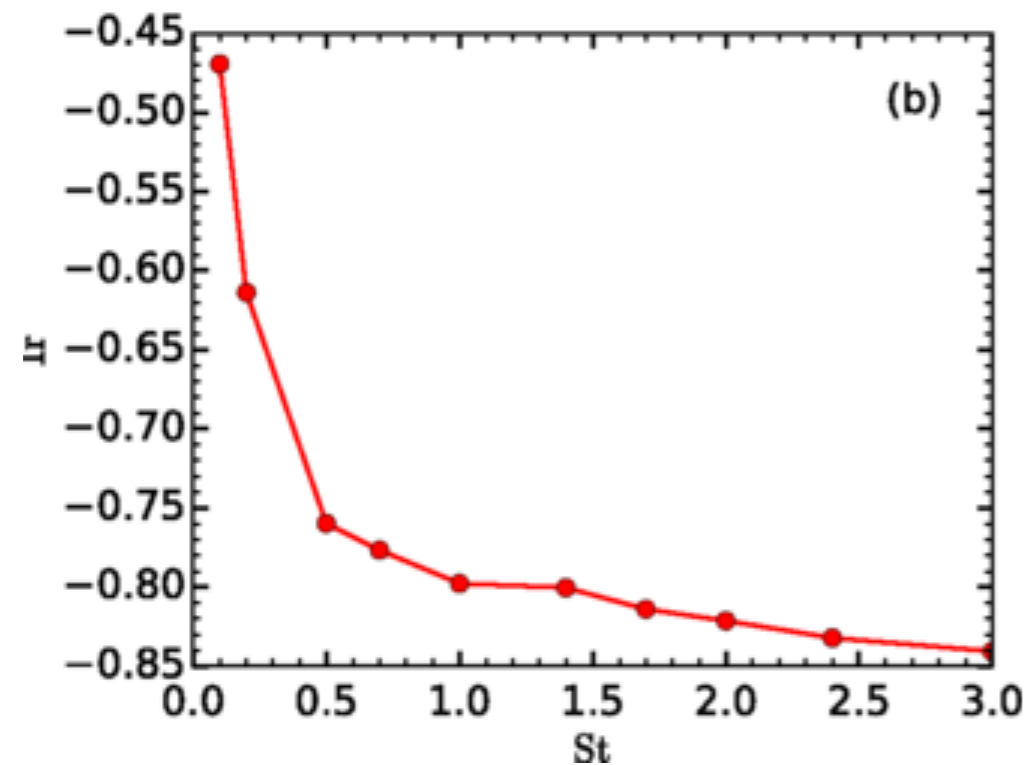
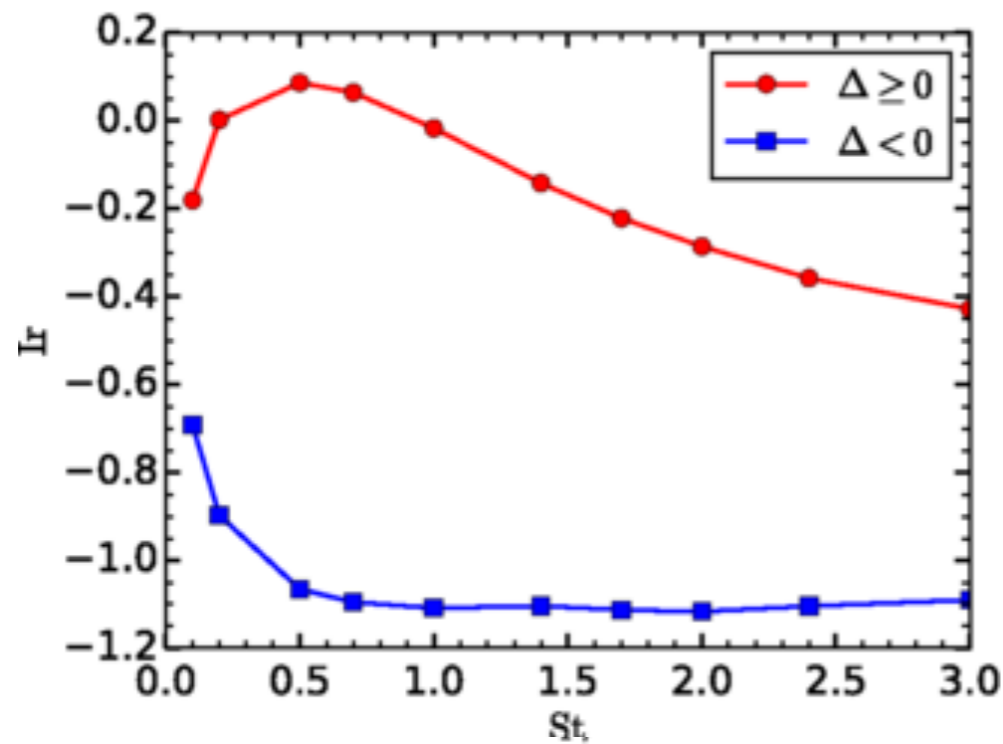
- P. Maity, R. Govindarajan, and S. S. Ray, “Statistics of Lagrangian trajectories in a rotating turbulent flow”, Phys. Rev. E. **100**, 043110 (2019).
- S. Ciliberto, S. Joubaud, and A. Petrosyan, J. Stat. Mech. P12003 (2010).

# Irreversibility and rotation



**P. Maity**, R. Govindarajan, and S. S. Ray, “*Statistics of Lagrangian trajectories in a rotating turbulent flow*”, Phys. Rev. E. **100**, 043110 (2019).

# Lagrangian Irreversibility (finite St effects)



Irreversibility as a function of Stokes number from Bhatnagar *et al.*

- A. Bhatnagar, A. Gupta, D. Mitra, and R. Pandit, Phys. Rev. E **97**, 033102 (2018).

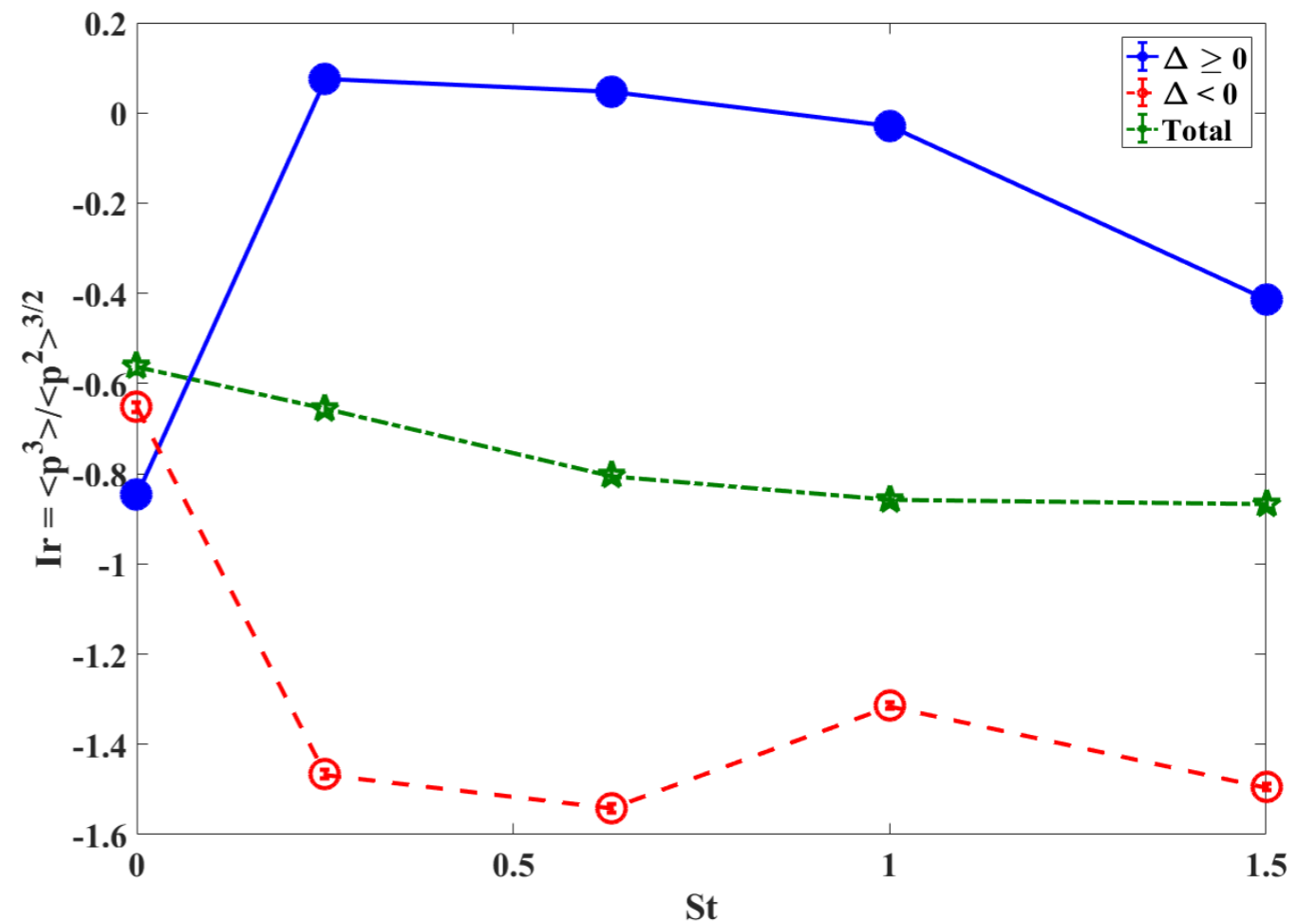
$$\Delta = \frac{27}{4}R^2 + Q^3$$

$\Delta \geq 0$  vortical,  $\Delta < 0$  straining

$$Q = \frac{1}{2}(\|\omega\|^2 - \|S\|^2), R = -\frac{1}{4}\omega_i S_{ij} \omega_j - \|S\|^3$$

$$\|\omega\| = \text{tr}[\omega\omega^T], \|S\| = \text{tr}[SS^T]$$

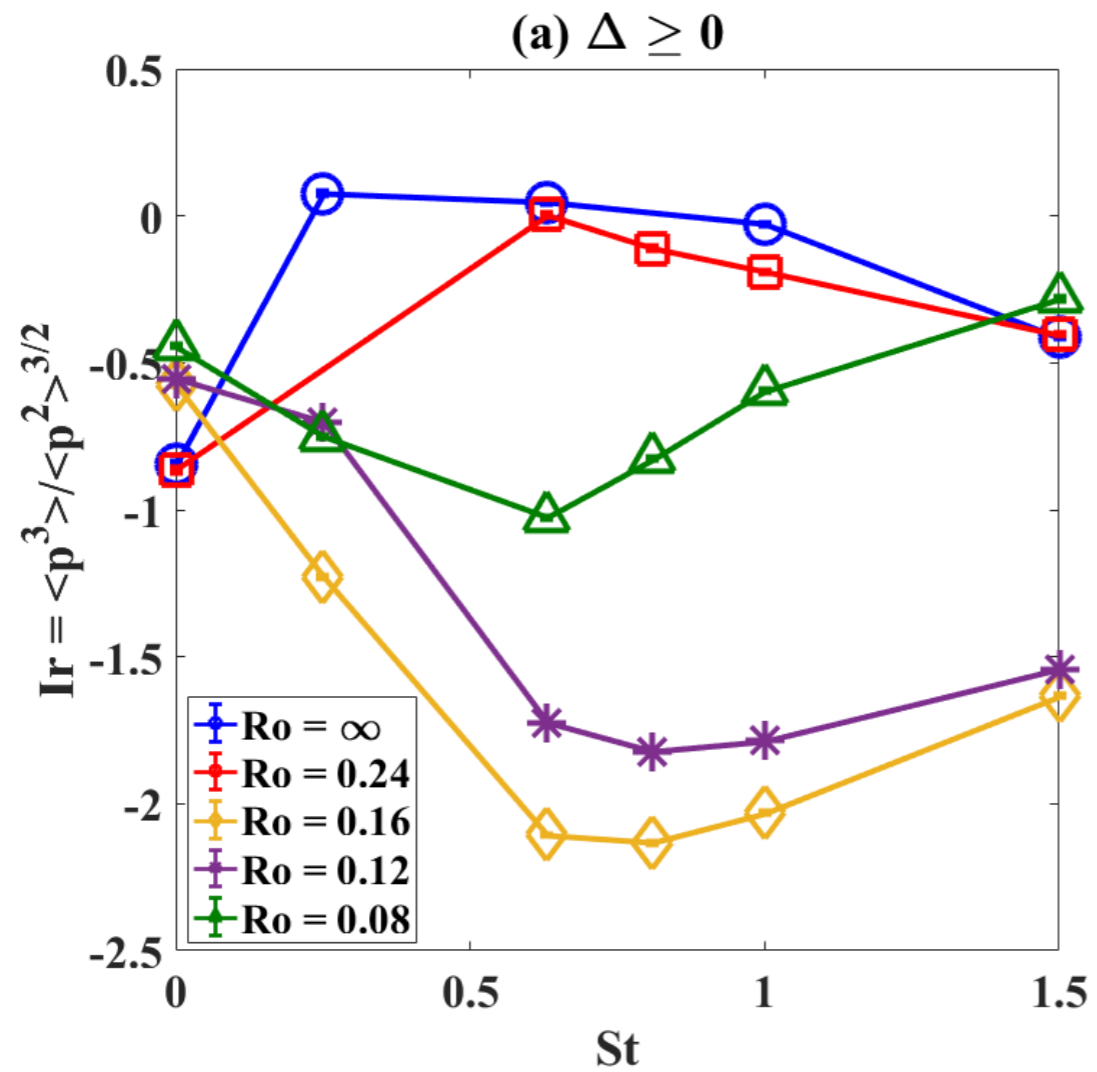
$$S = \frac{1}{2}[\nabla\mathbf{u} + \nabla\mathbf{u}^T], \omega = \frac{1}{2}[\nabla\mathbf{u} - \nabla\mathbf{u}^T]$$



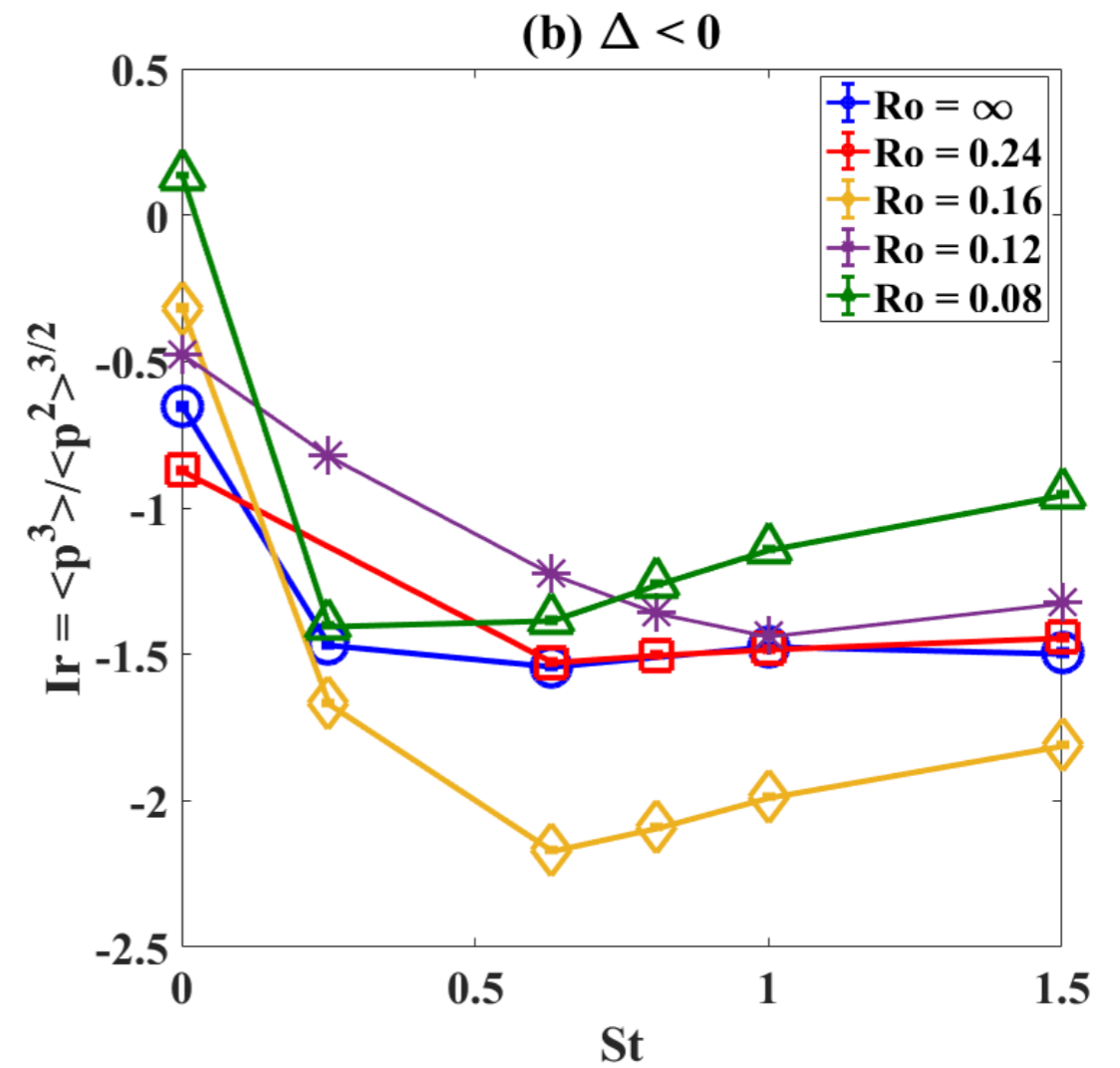
Skewness of Lagrangian power as a function of Stokes number (our simulations).

# Lagrangian irreversibility with finite St with rotation

## Vortical Region



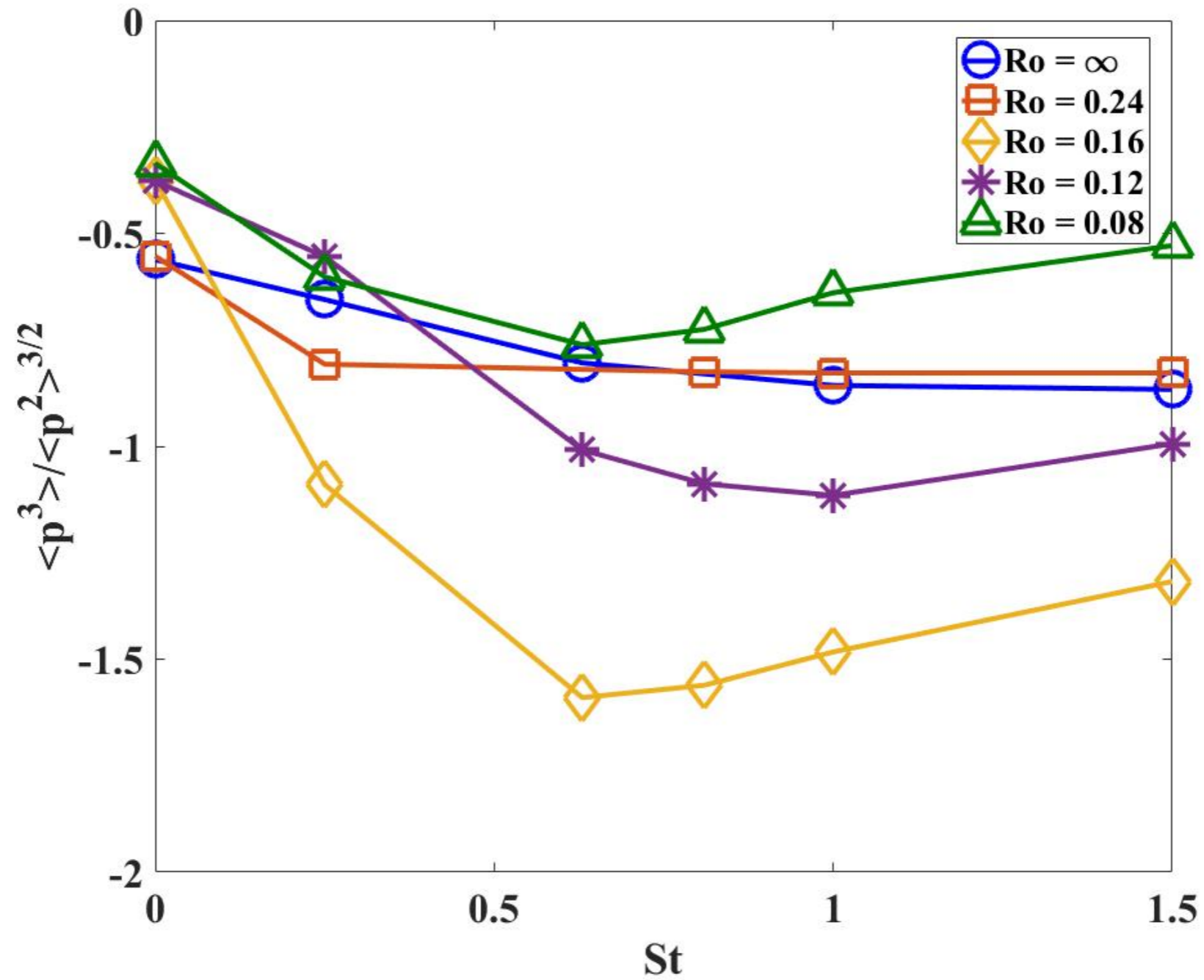
## Straining Region



Lagrangian irreversibility as a function of Stokes number for various rotation rates (Maity *et al.* 2020)

# Lagrangian irreversibility with finite St with rotation

Total



Total Lagrangian irreversibility as a function of Stokes number for various values of rotation rates (Maity *et. al.* 2020)



## References:

1. U. Frisch, *Turbulence : The legacy of A. N. Kolmogorov* (Cambridge University Press, 1995).
2. H. Xu *et al.*, *Proc. Natl. Acad. Sci. (USA)* **111**, 7558 (2018).
3. A. Bhatnagar, A. Gupta, D. Mitra, and R. Pandit, *Phys. Rev. E.* **97**, 033102 (2018).
4. **P. Maity, R. Govindarajan, and S. S. Ray**, “Statistics of Lagrangian trajectories in a rotating turbulent flow”, *Phys. Rev. E.* **100**, 043110 (2019).
5. Luca Biferale *et. al.*, “Coherent Structures and Extreme Events in Rotating Multiphase Turbulent Flows”, *Phys. Rev. E* **6**, 041036 (2016).
6. L. Del Castello and H. J. H. Clercx, *Phys. Rev. Lett.* **107**, 214502 (2011).
7. L. Del Castello and H. J. H. Clercx, *Phys. Rev. E* **83**, 056316 (2011)

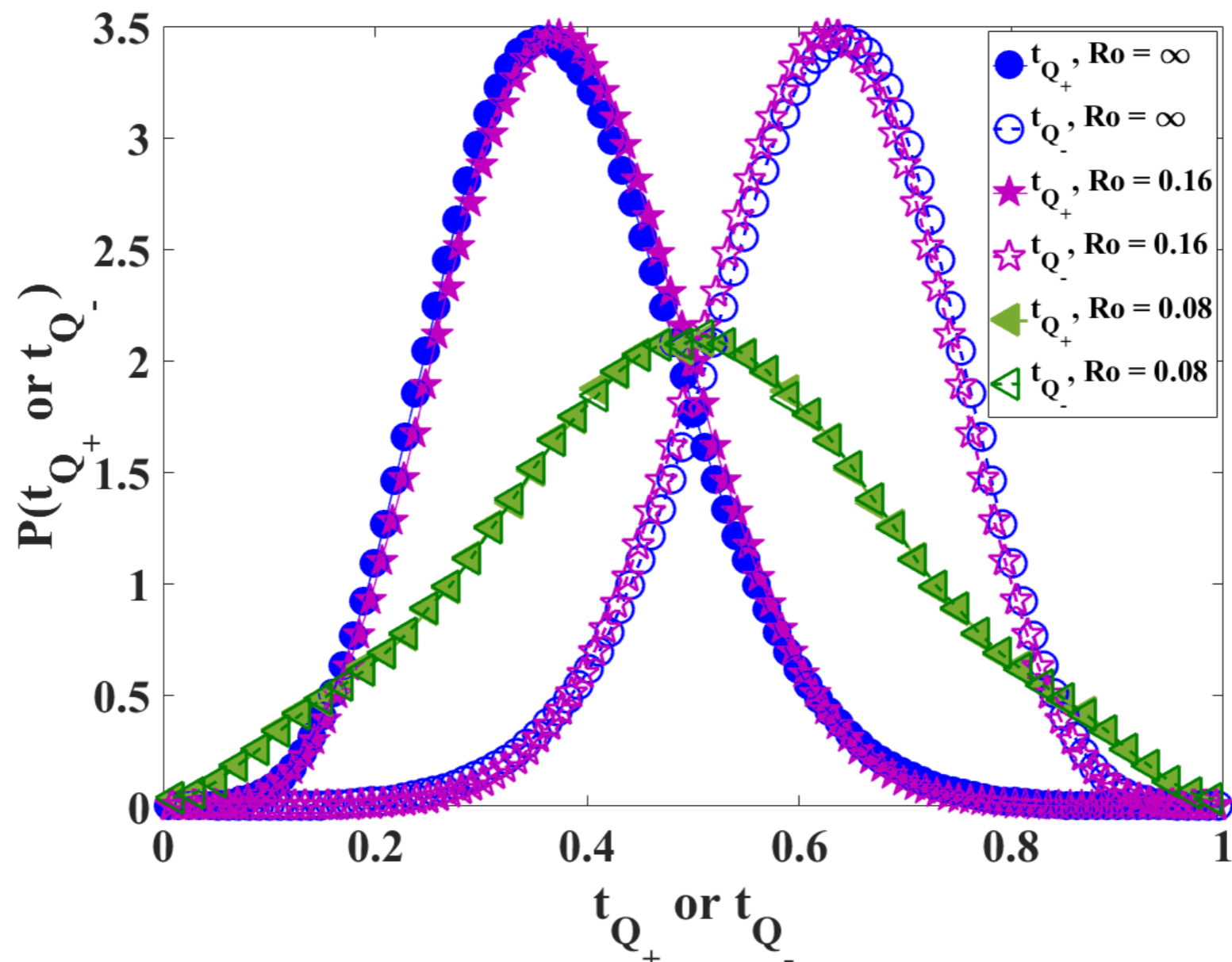
**Thanks for you  
attention**

# Statistics of Lagrangian Trajectories

$$Q = \frac{1}{2}(\|\omega\|^2 - \|S\|^2), R = -\frac{1}{4}\omega_i S_{ij} \omega_j - \|S\|^3$$

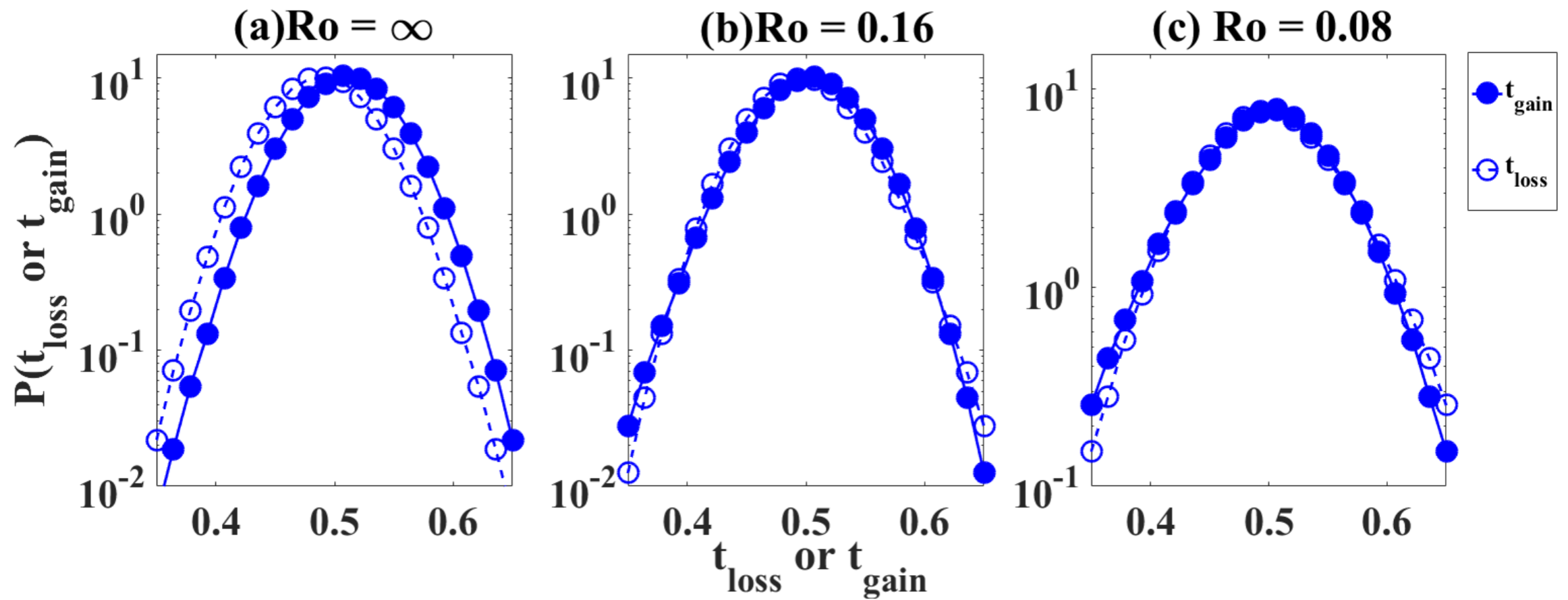
$$\|\omega\| = \text{tr}[\omega\omega^T], \|S\| = \text{tr}[SS^T]$$

$$S = \frac{1}{2}[\nabla\mathbf{u} + \nabla\mathbf{u}^T], \omega = \frac{1}{2}[\nabla\mathbf{u} - \nabla\mathbf{u}^T]$$



- Akshay Bhatnagar, Anupam Gupta, Dhruvadya Mitra, Rahul Pandit, and Prasad Perlekar, “How long do particles spend in vortical regions in turbulent flows?”, Phys. Rev. E **94**, 053119 (2016).

# Statistics of Lagrangian power



**P. Maity**, R. Govindarajan, and S. S. Ray, “*Statistics of Lagrangian trajectories in a rotating turbulent flow*”, Phys. Rev. E. **100**, 043110 (2019).

# Flow dynamics

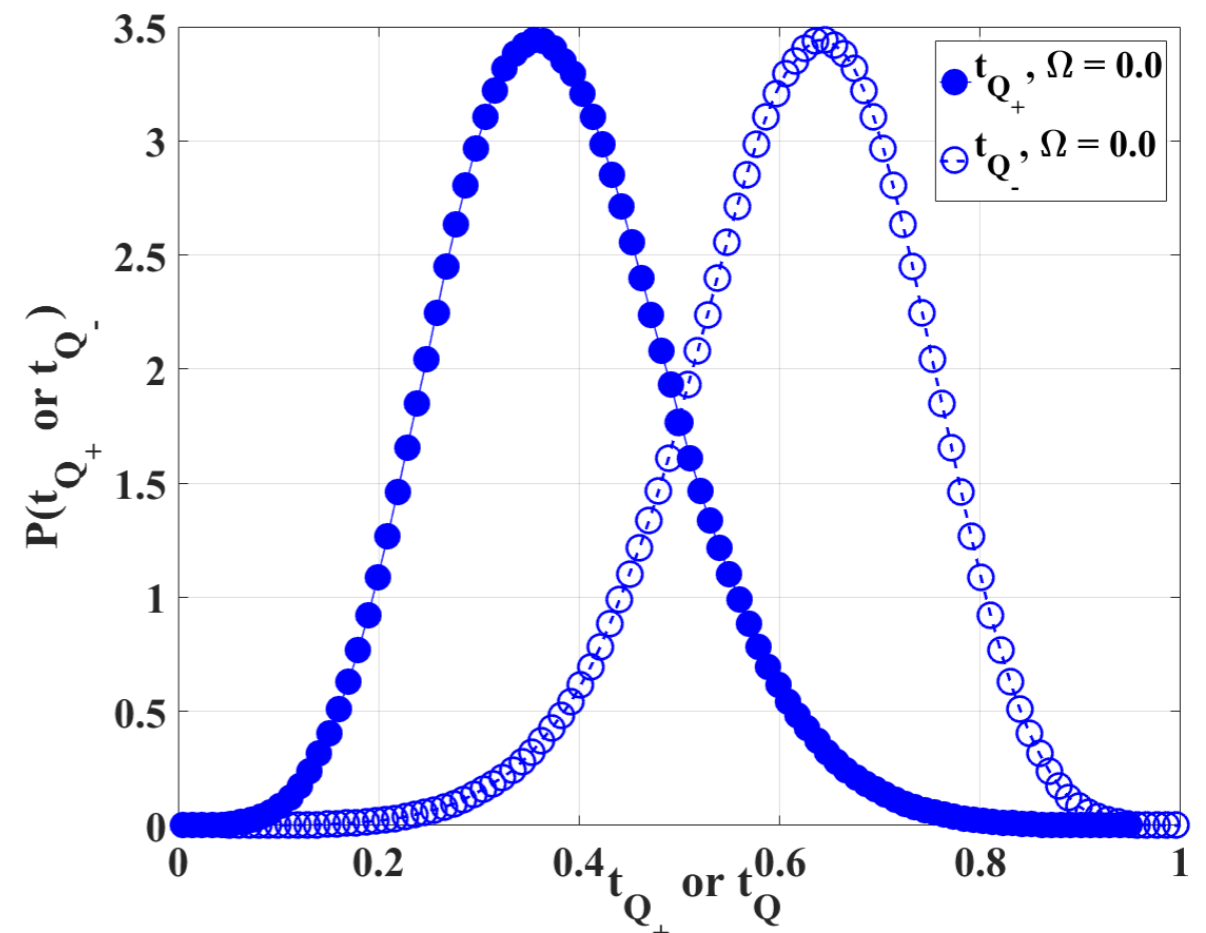
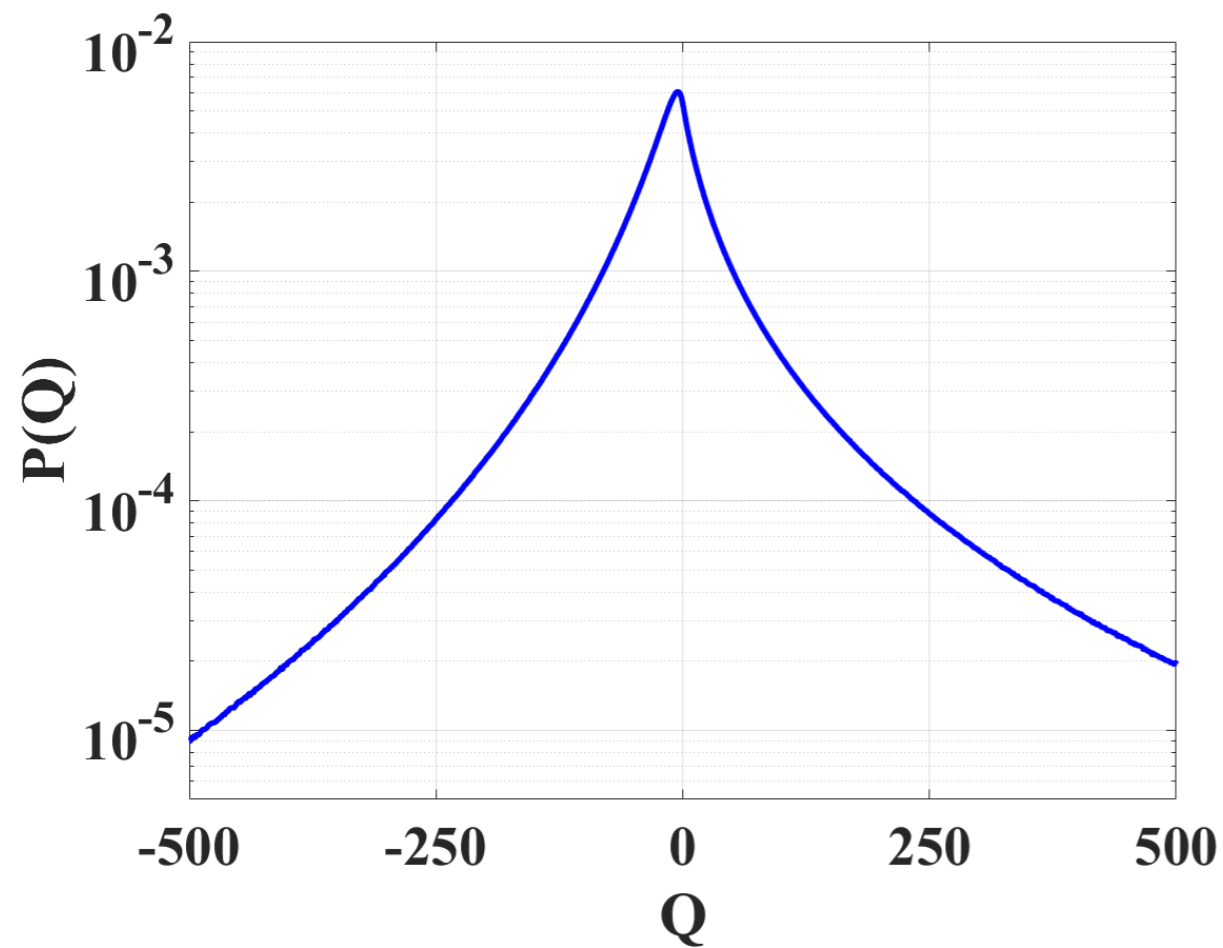
$$Q = \frac{1}{2}(\|\omega\|^2 - \|S\|^2)$$

$$\|\omega\| = \text{tr}[\omega\omega^T], \|S\| = \text{tr}[SS^T]$$

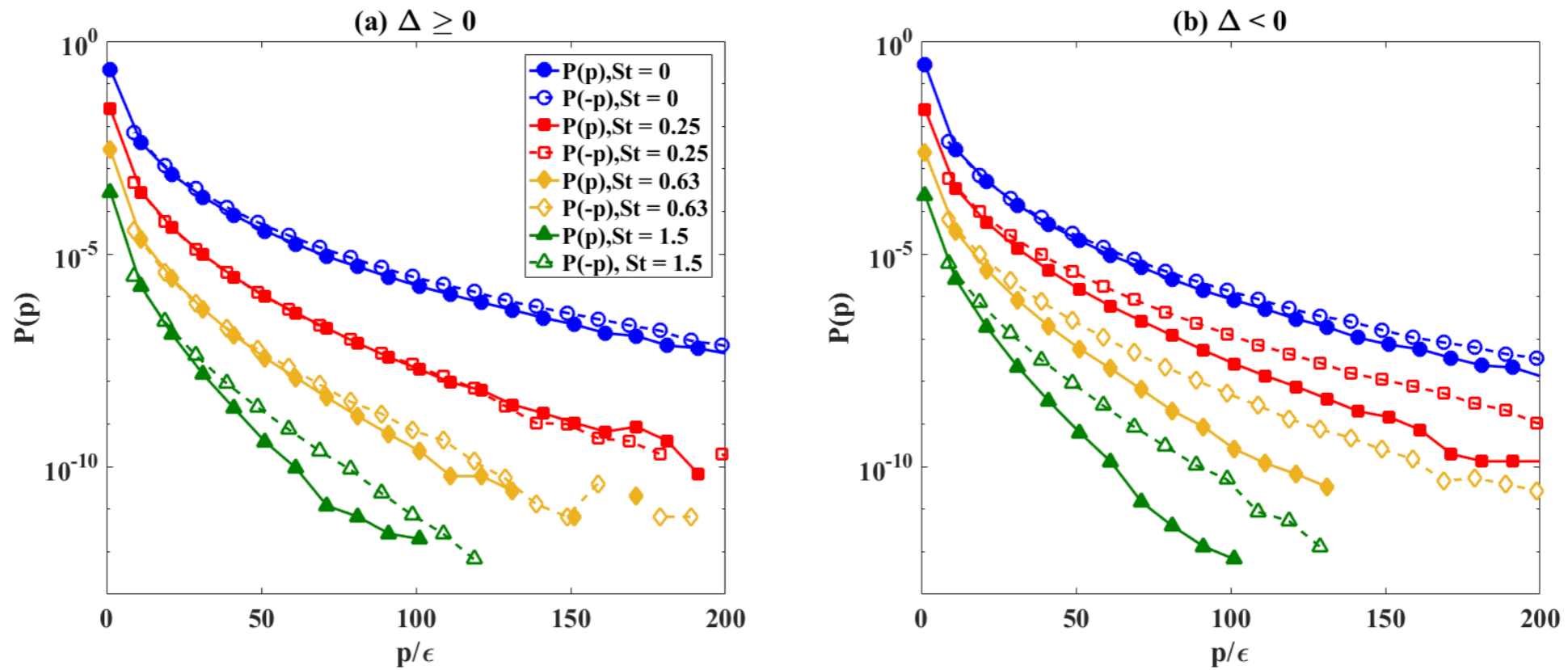
$$S = \frac{1}{2}[\nabla\mathbf{u} + \nabla\mathbf{u}^T], \omega = \frac{1}{2}[\nabla\mathbf{u} - \nabla\mathbf{u}^T]$$

$Q \geq 0 \rightarrow$  vortical

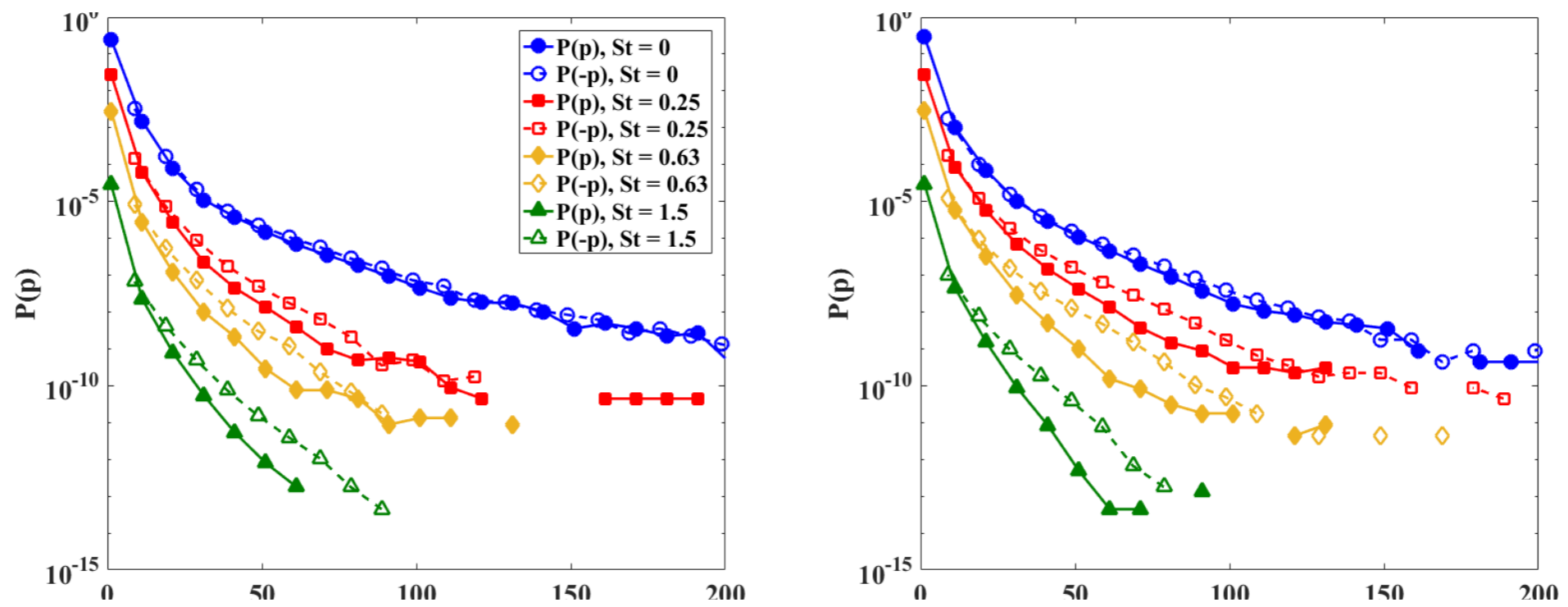
$Q < 0 \rightarrow$  straining







PDFs of instantaneous power  $p$  in (a) vortical and (b) straining region for  $Ro = \infty, \Omega = 0$ .



PDFs of instantaneous power  $p$  in (a) vortical and (b) straining region for  $Ro = 0.12, \Omega = 1.0$ .