

Kelvin circulation theorem, dynamic metric, and FQHE

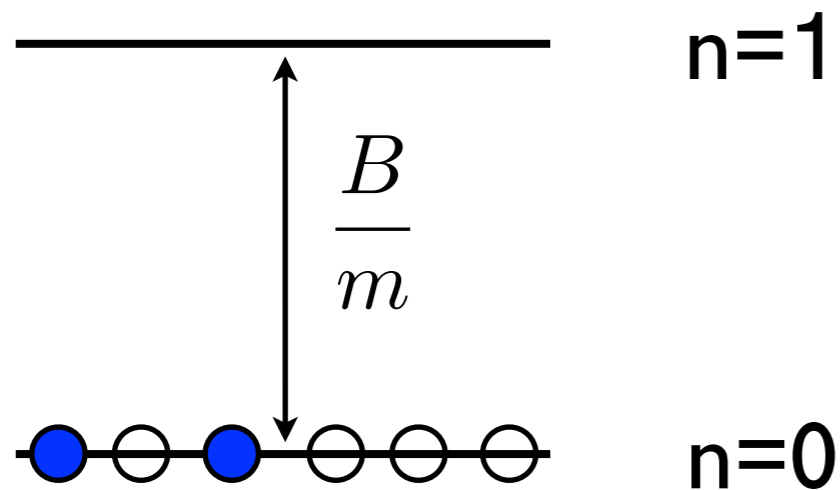
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Novel Phases of Quantum Matter
ICTS, 27 December 2019

Broad view

- Fractional quantum Hall effect, rich physics
 - wave function, CFT
 - TQFT, anyons
 - flux attachment
 - LLL approaches (Shankar-Murthy...)
 - duality

Fractional quantum Hall effect

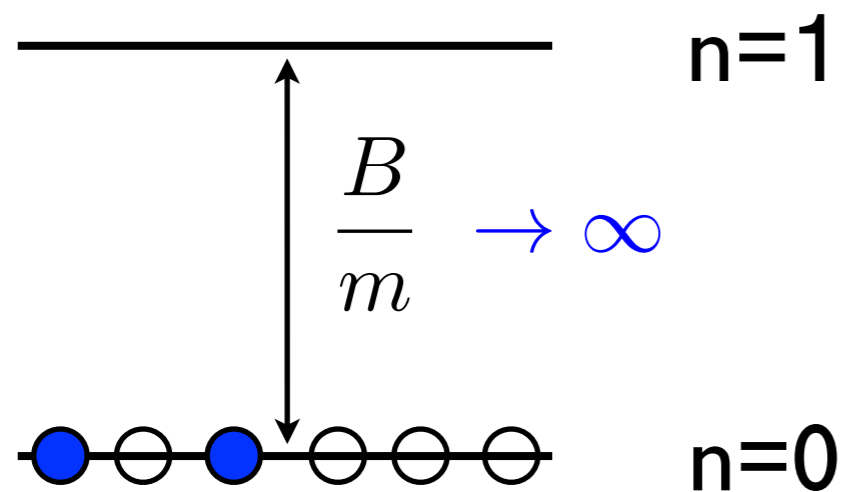
$$H = \sum_a \frac{(\mathbf{p}_a + e\mathbf{A}_a)^2}{2m} + \sum_{\langle a,b \rangle} \frac{e^2}{|\mathbf{x}_a - \mathbf{x}_b|}$$



Fractional quantum Hall effect

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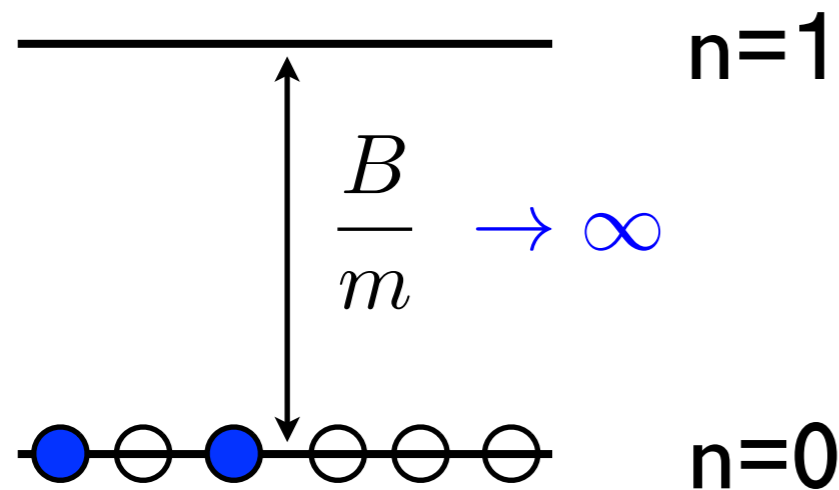
$m \rightarrow 0$



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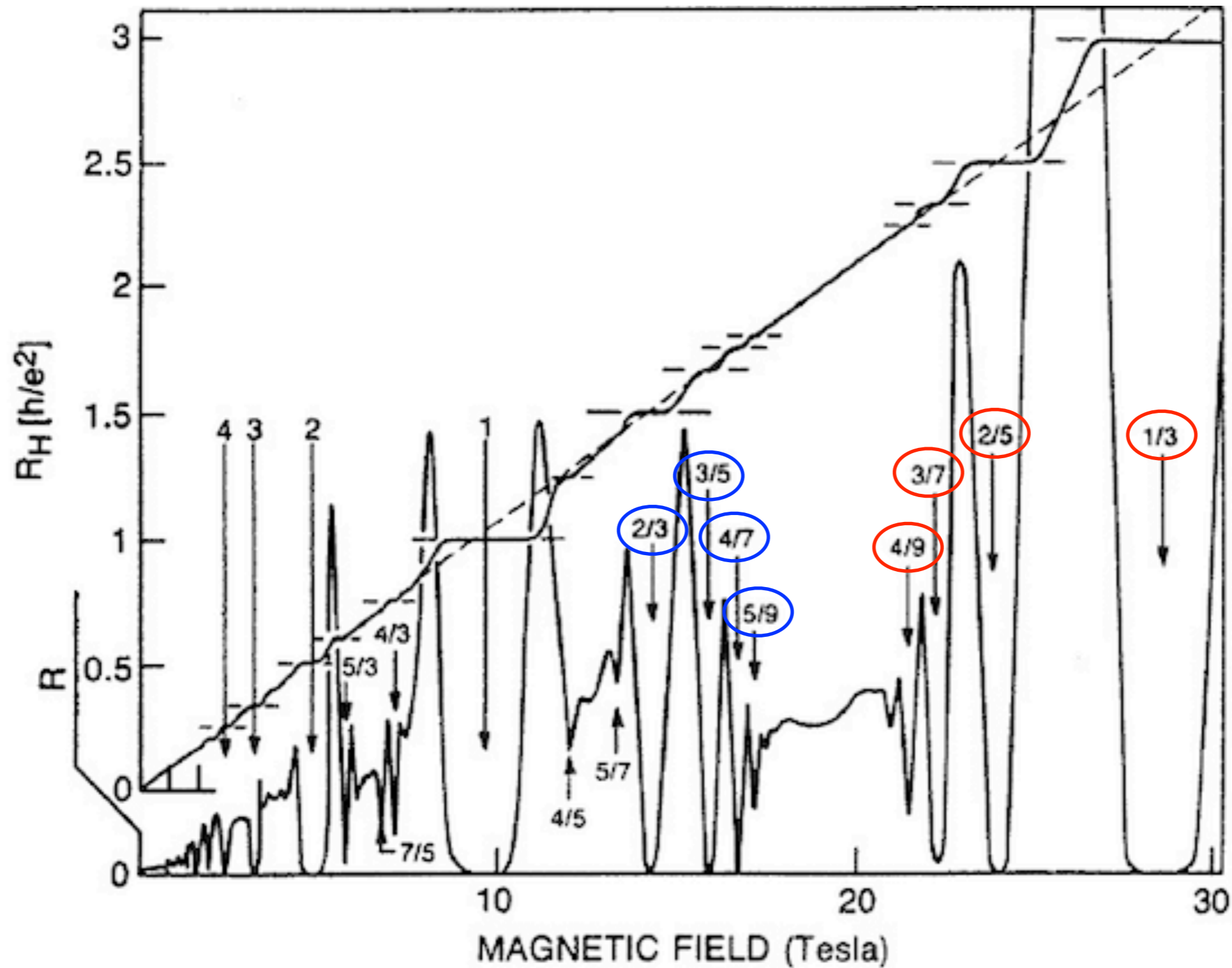
$m \rightarrow 0$



$$H = P_{\text{LLL}} \sum_{a,b} \frac{e^2}{|\mathbf{x}_a - \mathbf{x}_b|}$$

Projection to lowest Landau level

Rich phenomenology



gapped QH states at $\nu = \frac{N+1}{2N+1}$ $\nu = \frac{N}{2N+1}$

Composite fermion

- Near half filling: a new quasiparticle
- Traditional view “composite fermion” = electron + 2 flux quanta
- That picture however is not particle-hole symmetric
- New view: composite fermion as a “dual fermion”

Dirac composite fermion

- Composite fermions are “massless Dirac fermions”:
have Berry phase π around any loop around $p=0$
- Number of CFs = $1/2$ number of magnetic flux quanta
- CFs live in a magnetic field $b = B - 4\pi n_e$
- Number of electron is NOT number of CFs

LLL projection

- Some fundamental questions remain :
 - is there projection to LLL in the composite fermion theory?
 - in particular, does the theory realize the algebra of the projected electron density operator (GMP algebra)
 - does density-density correlator $\sim q^4$?

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

$$\rho \sim \frac{qj}{\omega}$$

$$\langle \rho \rho \rangle \sim q^2$$

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Spin of magnetoroton

- Lowest neutral excitation: magnetoroton (Girvin, MacDonald, Platzman 1986)
- near $q=0$: excitations can be classified by orbital angular momentum (spin)

$$\langle 0|\rho|0\rangle \sim \underbrace{\langle 0|\rho|\text{MR}, q\rangle}_{q^2} \underbrace{\langle \text{MR}, q|\rho|0\rangle}_{q^2} \sim q^4$$

$$\langle \text{spin} \pm n, q|\rho|0\rangle \sim (q_x \pm iq_y)^n$$

Magnetoroton at $q=0$ has spin 2
a dynamical graviton (Haldane)?

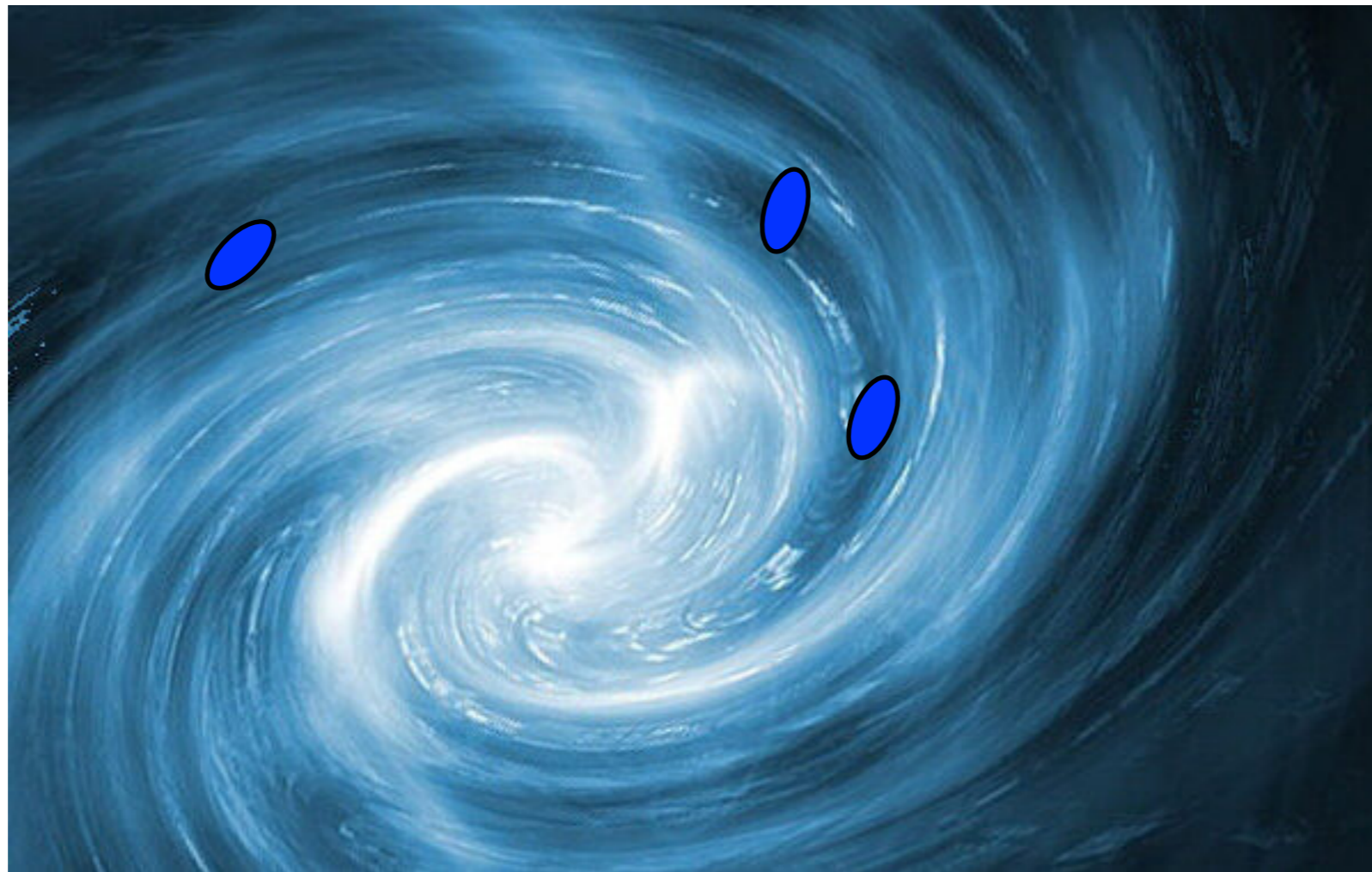
A fluid view on FQHE

- Quantum Hall fluid is a fluid, described by a hydrodynamic theory
- The fluid has an internal tensor degree of freedom: a dynamic “metric”
- Particles of the fluid are dipoles
- Provide a consistent and predictive picture of the long-distance behavior of QH fluids

Ref.: DTS 1907.07187

works with A. Gromov and D.X. Nguyen

A crash course in hydrodynamics



Hydrodynamics

Landau 1941

- can be formulated as a dynamical system with the Poisson brackets

$$\{\pi_i(\mathbf{x}), n(\mathbf{y})\} = n(\mathbf{x})\partial_i\delta(\mathbf{x} - \mathbf{y})$$

$$\{\pi_i(\mathbf{x}), \pi_j(\mathbf{y})\} = [\pi_j(\mathbf{x})\partial_i + \pi_i(\mathbf{y})\partial_j]\delta(\mathbf{x} - \mathbf{y})$$

- and Hamiltonian

$$H = \int d\mathbf{x} \left[\frac{1}{2m} \frac{\vec{\pi}^2(\mathbf{x})}{n(\mathbf{x})} + \epsilon(n(\mathbf{x})) \right]$$

$$\dot{n} = \{H, n\}$$

$$\dot{\pi}_i = \{H, \pi_i\}$$

Extending Poisson algebra

- Let's introduce a “dynamical metric” $G_{ij}(\mathbf{x})$
- The fact that G_{ij} transforms like a tensor fixes the Poisson bracket

$$\{G_{ij}(\mathbf{x}), \pi_k(\mathbf{y})\} = (G_{ik}(\mathbf{x})\partial_j + G_{jk}(\mathbf{x})\partial_i + \partial_k G_{ij})\delta(\mathbf{x} - \mathbf{y})$$

- $\{G, G\} = ?$

Chiral metric hydro

In 2 spatial dimensions

$$\{G_{ij}(\mathbf{x}), G_{kl}(\mathbf{y})\} = -\frac{1}{s}(\varepsilon_{ik}G_{jk} + \varepsilon_{il}G_{jk} + \varepsilon_{jk}G_{il} + \varepsilon_{jl}G_{ik})\delta(\mathbf{x} - \mathbf{y})$$

- We can consistently impose $(\det \mathbf{G})^{1/2} = n$
- Hydrodynamics equations

$$\dot{A} = \{H, A\}$$

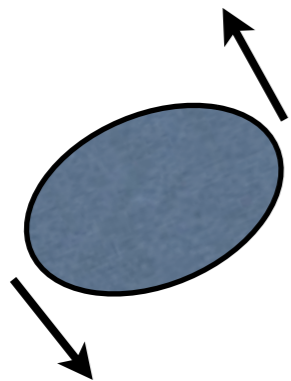
Qualitative discussion

- “Metric” perturbation: a gapped spin-2 mode

$$G_{ij} = n(\delta_{ij} + Q_{ij})$$

$$Q_{xx} = -Q_{yy} \sim \cos \omega t$$

$$Q_{xy} \sim \sin \omega t$$



$$\omega = \frac{2\mu}{ns} \leftarrow \text{“Lamé constant”}$$

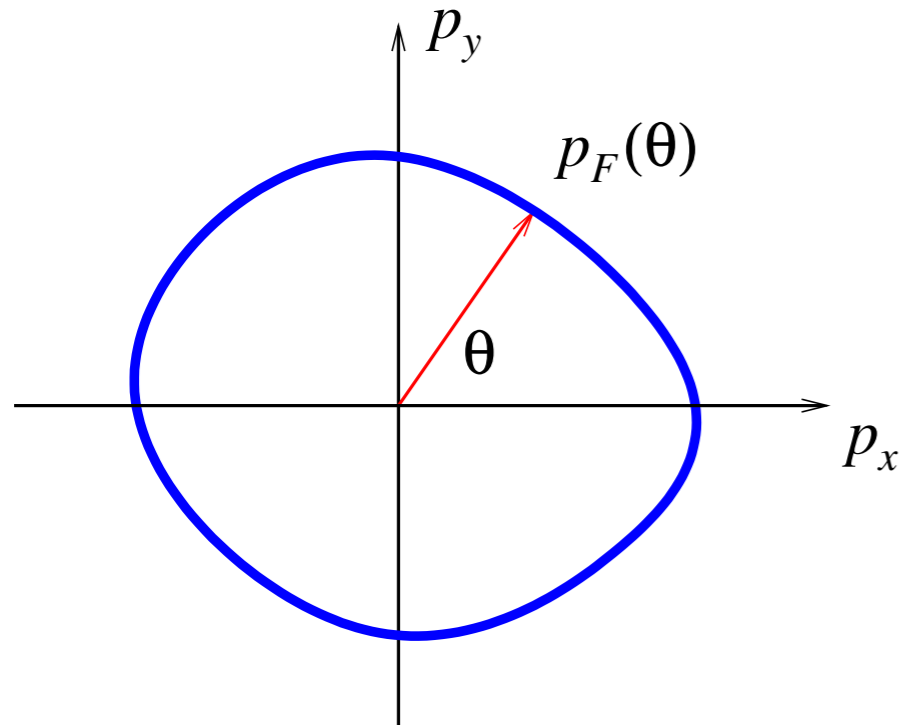
At small frequencies, a fluid with Hall viscosity

$$\eta^H = \frac{sn}{2}$$

s = average “orbital spin”

What is the relevance
to FQHE?

CF Fermi surface



Low-energy, long-wavelength excitations: fluctuations of the shape of the Fermi surface

$$p_F(t, \mathbf{x}, \theta) = p_F^0 + \sum_{n=-\infty}^{\infty} u_n(t, \mathbf{x}) e^{-in\theta}.$$

One scalar field per spin

At low momenta we can limit ourselves to a few lowest modes

$$v_F q \ll \omega$$

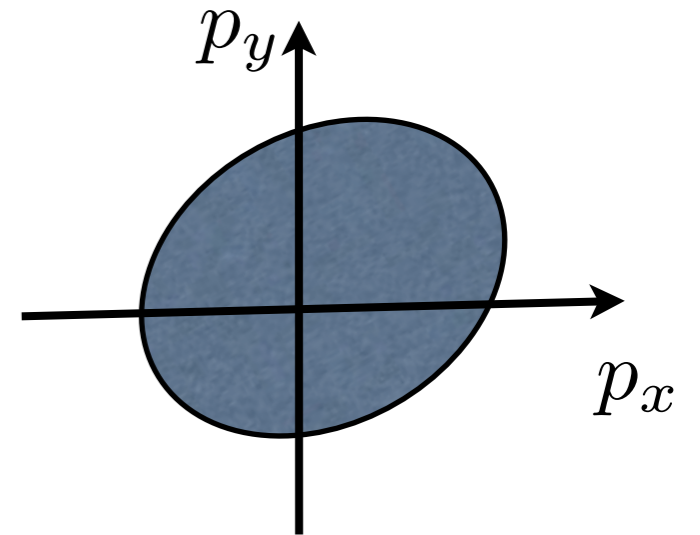
G_{ij} parametrizes an elliptical Fermi surface

“Nematic” hydrodynamics

- Degrees of freedom:

- density

$$n(\mathbf{x}) = \int \frac{d\mathbf{p}}{(2\pi)^2} f(\mathbf{x}, \mathbf{p})$$



- momentum density

$$\pi_i(\mathbf{x}) = \int \frac{d\mathbf{p}}{(2\pi)^2} p_i f(\mathbf{x}, \mathbf{p})$$

- effective metric

$$\int \frac{d\mathbf{p}}{(2\pi)^2} p_i p_j f(\mathbf{x}, \mathbf{p}) = \frac{\pi_i \pi_j}{n} + \pi n(\mathbf{x}) G_{ij}(\mathbf{x})$$

$$\sqrt{\det G} = n$$

Chiral metric hydro

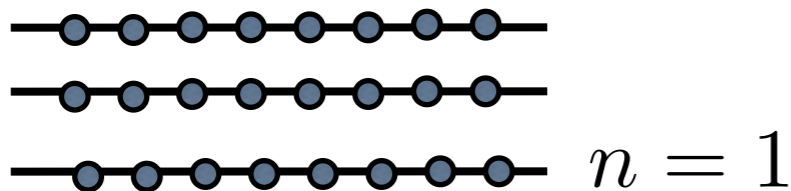
- The commutation relations can be “derived” from semiclassical arguments

s related from the Hall viscosity

= average “orbital spin” of composite fermion

$$\nu = \frac{N}{2N + 1}$$

$$s = \frac{1}{N + \frac{1}{2}} \left(\frac{1}{2} \cdot 0 + 1 + 2 + \dots + N \right) = \frac{N(N + 1)}{2N + 1}.$$



Dipoles

- On the LLL, the CF are electric dipole
- dipole moment proportional to and perpendicular to momentum
- dipole density

$$\frac{\epsilon^{ij} \pi_j}{B}$$

Electron density

$$\rho = \frac{B - b}{4\pi} - \epsilon^{ij} \partial_i \left(\frac{\pi_j}{B} \right) \leftarrow \text{dipole contribution}$$

$$= n - \frac{b + \omega}{4\pi}$$

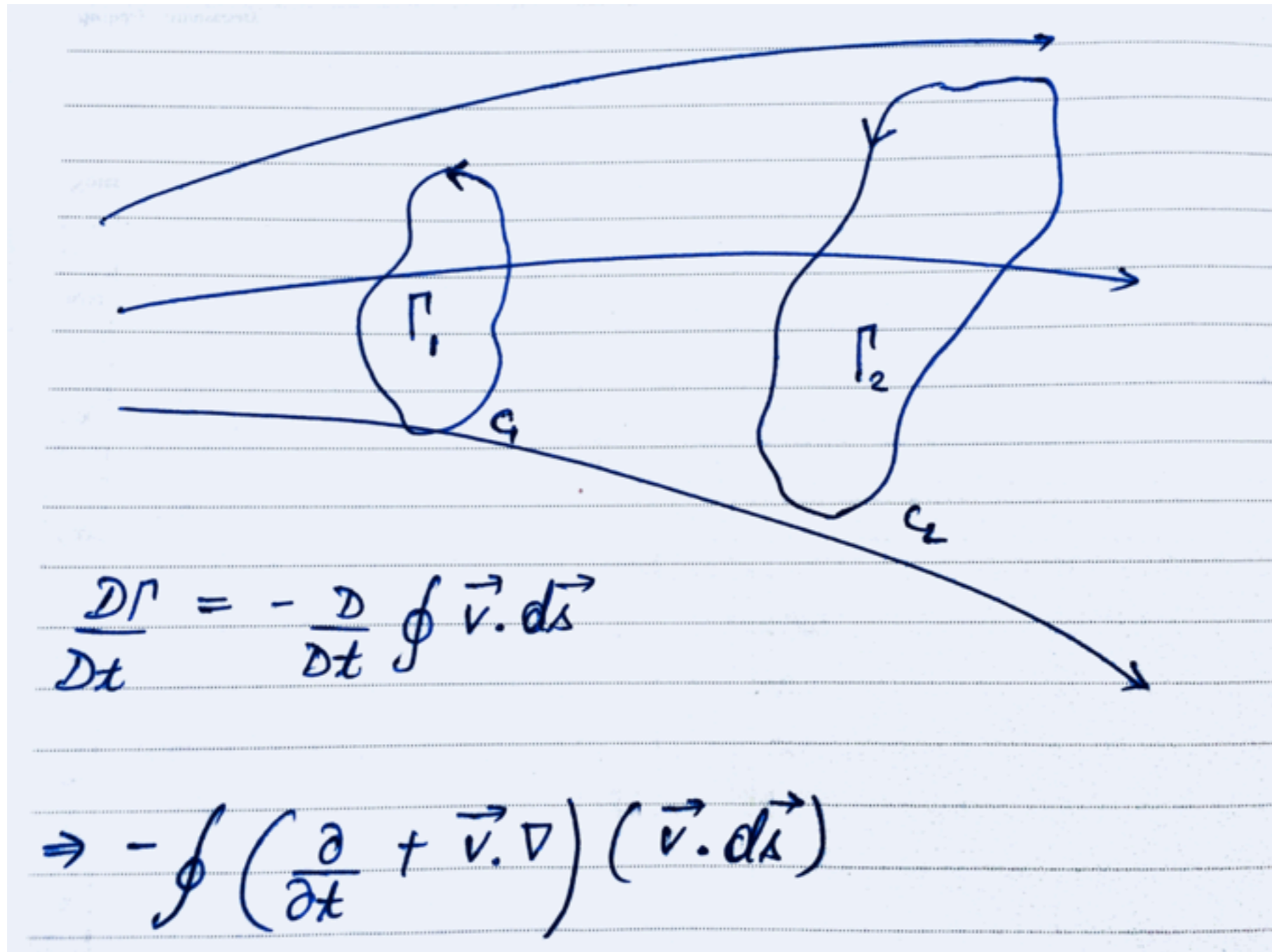
$$\omega = \vec{\nabla} \times \begin{pmatrix} \vec{\pi} \\ n \end{pmatrix}$$

“vorticity”

$$H = H_0[n, \pi_i, G_{ij}] + \int d\mathbf{x} \left(-a_0 n + \frac{\epsilon^{ij} E_j}{B} \pi_i \right)$$

Kelvin's circulation theorem

1869



- In ideal hydrodynamics vorticity is carried with the flow

$$\dot{\omega} + \vec{\nabla} \cdot (\omega \vec{v}) = 0$$

vorticity

$$\omega = \vec{\nabla} \times \begin{pmatrix} \vec{\pi} \\ n \end{pmatrix}$$

- Leads to an infinite number of conserved quantities (Casimirs of the Poisson algebra)

Kelvin's circulation theorem

- In the presence of magnetic field and metric degree of freedom, Kelvin's theorem is modified

$$\Omega = b + \omega + \frac{s}{2} \sqrt{G} R[G] \quad \dot{\Omega} + \vec{\nabla} \cdot (\Omega \vec{v}) = 0$$

$$\rho_e = \frac{B}{4\pi} - \frac{b + \omega}{4\pi}$$

- $\Omega = \text{constant}$

$$\delta \rho_e = \frac{s}{8\pi} \sqrt{G} R[G].$$

An immediate consequence

$$\delta\rho_e = \frac{s}{8\pi} \sqrt{GR[G]} \sim \partial_i \partial_j G_{ij}$$

$$\rightarrow \langle \delta\rho_e \delta\rho_e \rangle_{\omega, q} \sim q^4$$

Property of the lowest Landau level

In fact, numerical coefficient can be found for $\nu = \frac{N}{2N+1}$

$$\langle \delta\rho_e \delta\rho_e \rangle_q = \frac{N(N+1)}{2N+1} \frac{q^4}{16\pi B}.$$

$N=1$: matches exactly with the Laughlin wave function

Conclusion

- Low- q regime of FQH liquid: described by a fluid with internal metric degree of freedom, coupled to a gauge field
- Electron density \sim curvature of dynamic metric
- Static structure factor: algebraic calculation

Thank you