

Quantum Kagome Spin Liquids

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Herbertsmithite $\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$ Antiferromagnet



Quantum Kagome Spin Liquids: The case of Herbertsmithite

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- J. C. Trombe, F. Duc, *CEMES, Toulouse*, P. Strobel, *Grenoble, France*
- M. de Vries, G. Nilsen, A. Harrison, *Edinburgh, UK*
- S. Nakamae, F. Ladieu, D. L'Hote, P. Bonville, *CEA Saclay, France*
- *A. Mahajan's group (IIT Mumbai)*

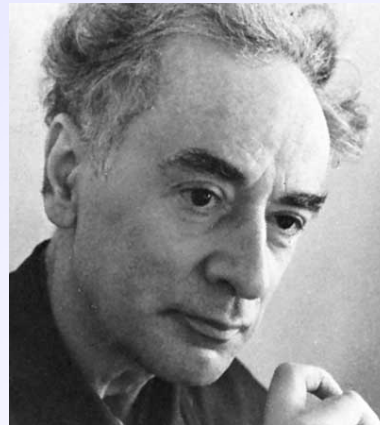
Quantum Kagome Spin Liquids

Part I

Novel states induced by
frustration

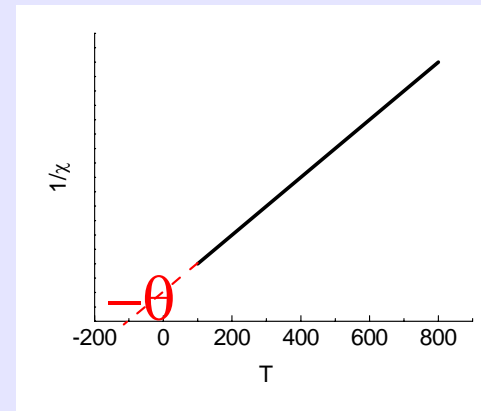
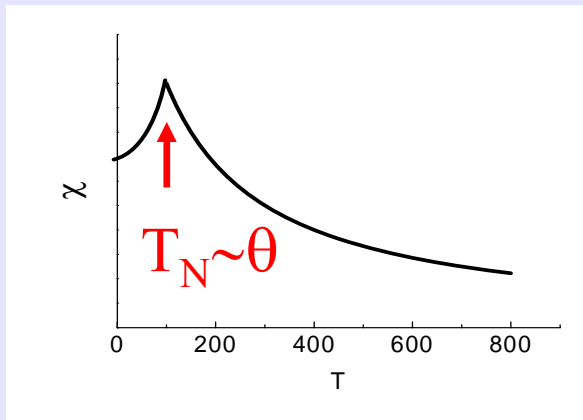
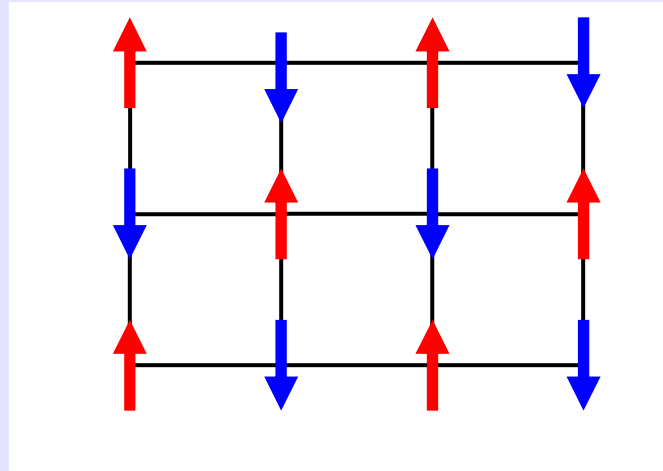
Towards spin liquids

Néel versus Anderson



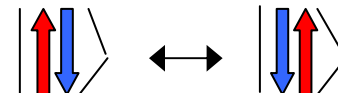
Néel State:
Antiferromagn.

$$\mathcal{H} = -J_{ij} \vec{S}_i \cdot \vec{S}_j, J_{ij} < 0$$



+

Quantum fluctuations for $S=1/2$



Antiferromagnetism, any alternative?



RESONATING VALENCE BONDS: A NEW KIND OF INSULATOR?*

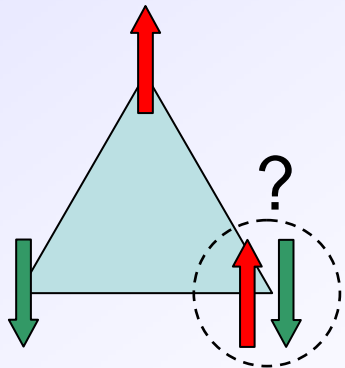
P. W. Anderson
Bell Laboratories, Murray Hill, New Jersey 07974
and
Cavendish Laboratory, Cambridge, England

(Received December 5, 1972; Invited**)

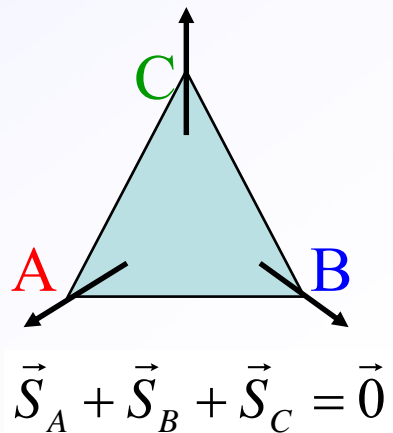
ABSTRACT

The possibility of a new kind of electronic state is pointed out, corresponding roughly to Pauling's idea of "resonating valence bonds" in metals. As observed by Pauling, a pure state of this type would be insulating; it would represent an alternative state to the Néel antiferromagnetic state for $S = 1/2$. An estimate of its energy is made in one case.

Geometrical Frustration of magnetic interactions



Class.
→



Antiferromagnetism, any alternative?



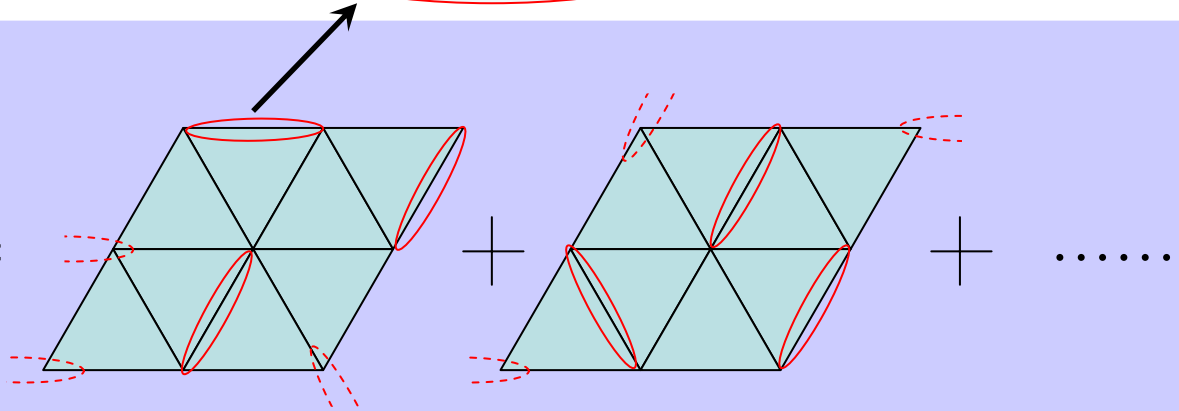
RESONATING VALENCE BONDS: A NEW KIND OF INSULATOR?*

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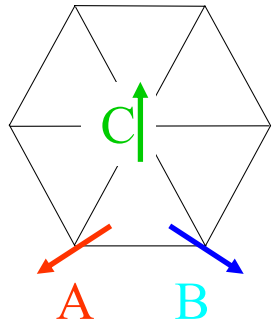
(Received December 5, 1972; Invited**)

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

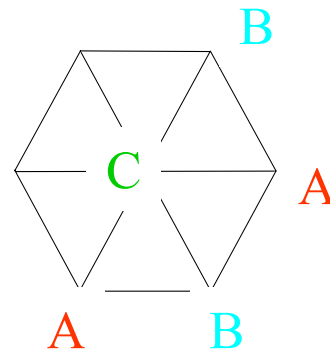
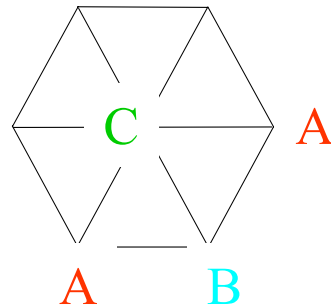
$$|\text{RVB}\rangle = \text{[Diagram 1]} + \text{[Diagram 2]} + \dots$$



Edge sharing: triangular lattice



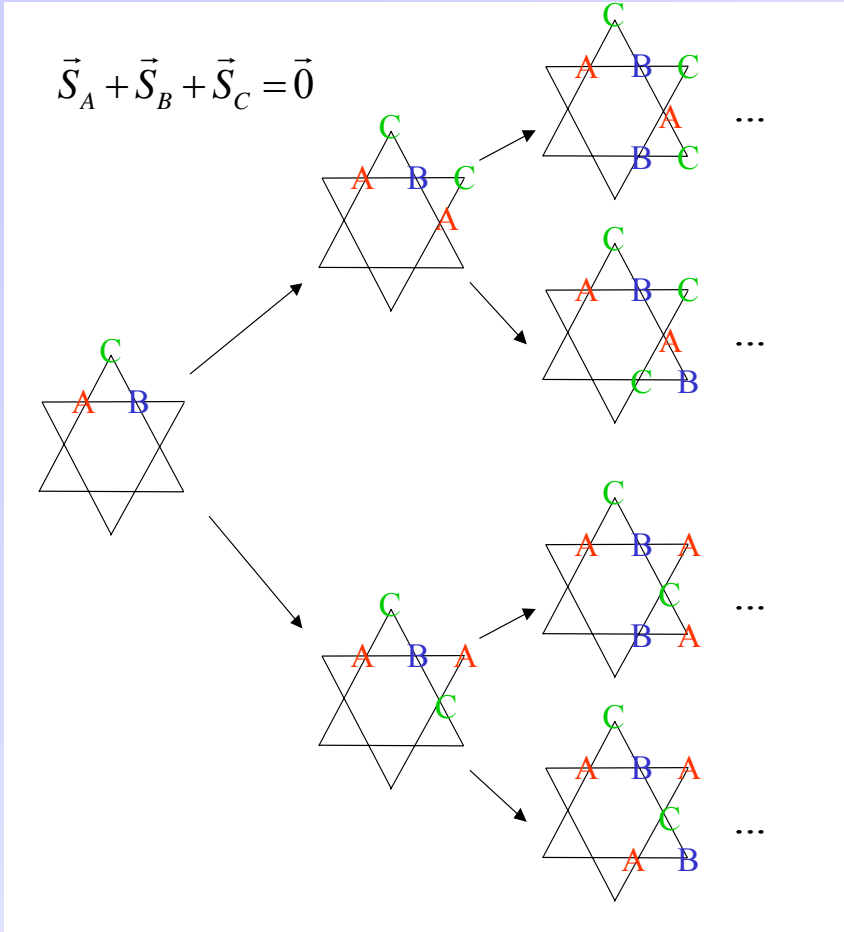
$$\vec{S}_A + \vec{S}_B + \vec{S}_C = \vec{0}$$



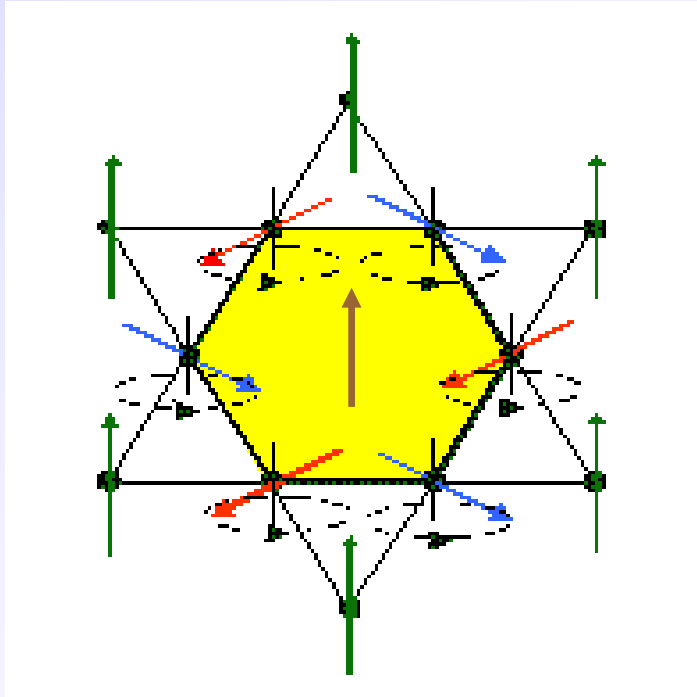
$$g = 2$$

(XY spins)

Corner sharing: classical kagomé lattice



Macroscopic degeneracy

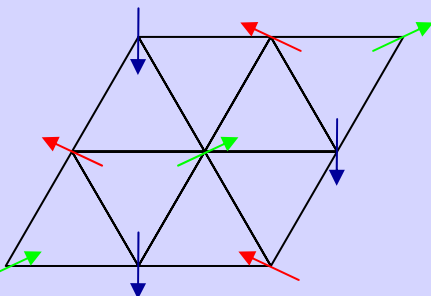


Soft modes

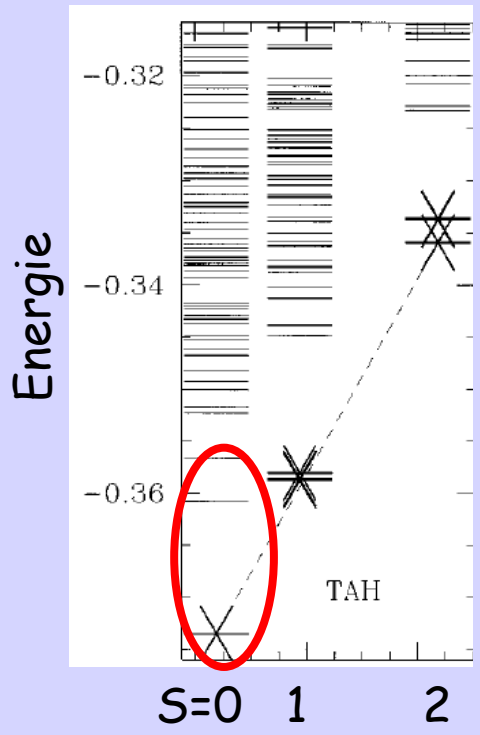
Antiferromagnetism, any alternative?

Exact diagonalizations
 Lecheminant, PRB **56**, 2521 (1997)
 Waldtmann, EPJB **2**, 501 (1998).

Triangular
 Edge sharing
 geometry

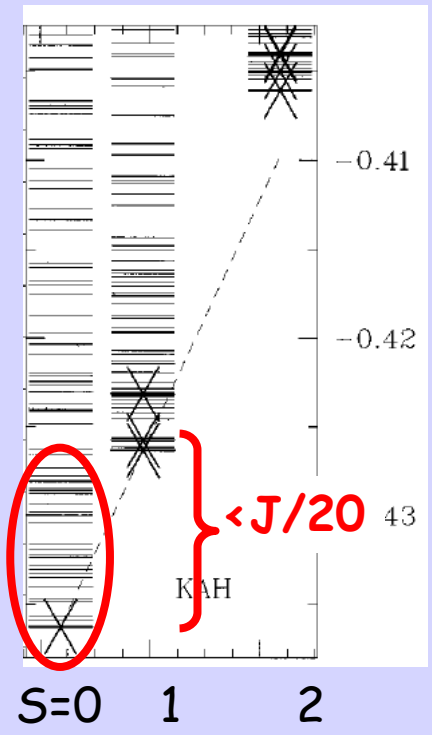
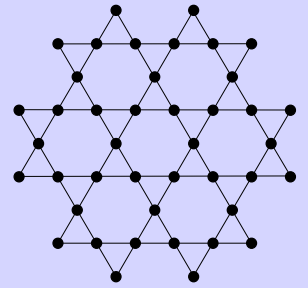


Fundamental Néel



Kagome

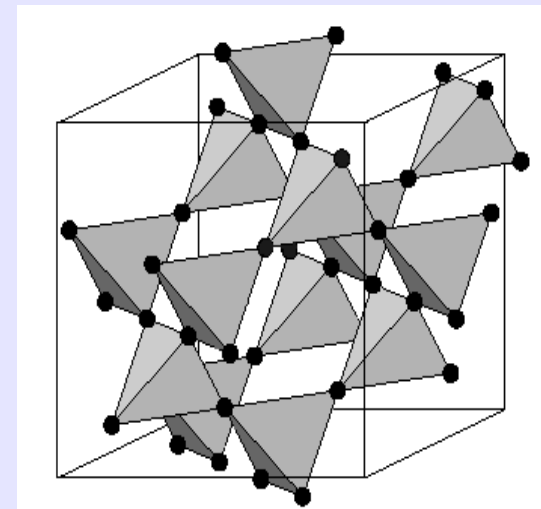
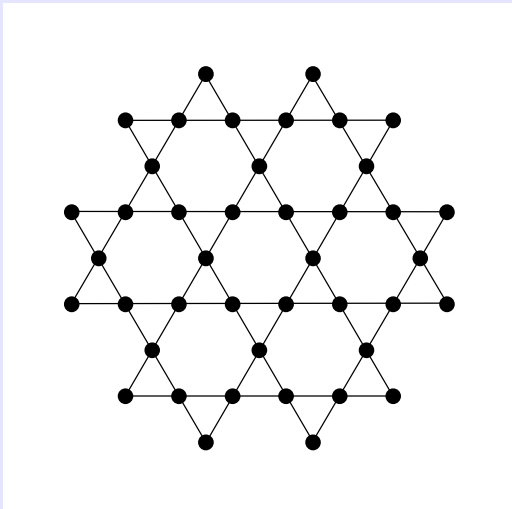
Corner sharing
 geometry



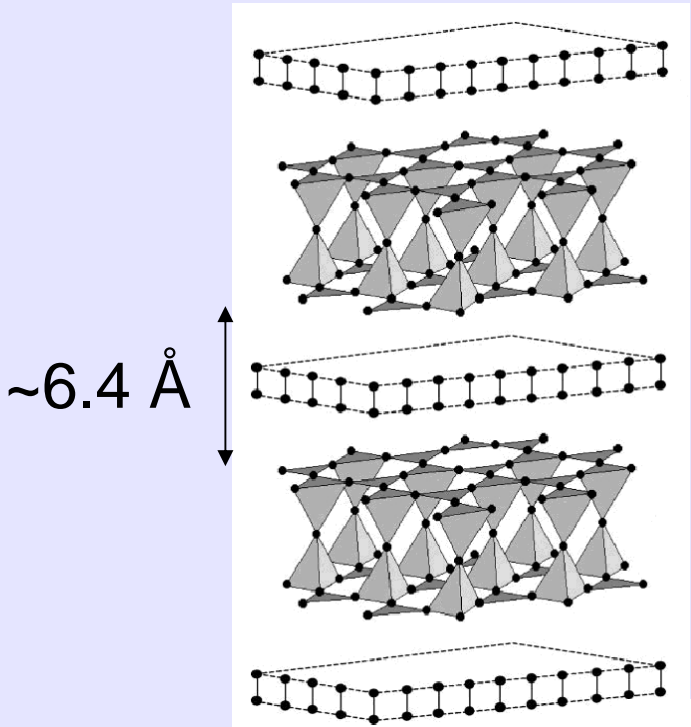
Fundamental 'spin liquid' ≠ Néel
 RVB ? Mila, PRL **81**, 2356 (2000)

IDEAL Highly Frustrated Magnet

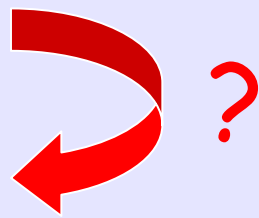
- Heisenberg spins
 - $S=1/2$ (quantum spins)
 - Corner sharing geometry:
 - Kagomé (2D) or pyrochlore (3D) lattice
- No « perturbation » (anisotropy, dipolar interaction, n.n. interactions, *dilution*)
- For kagomé: stacking should keep the 2D kagomé planes uncoupled



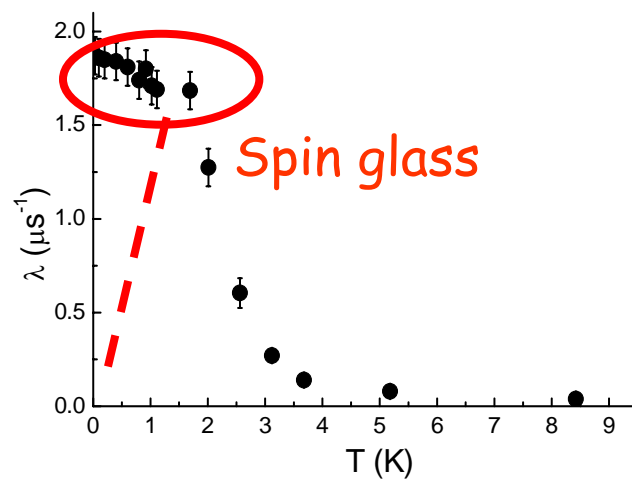
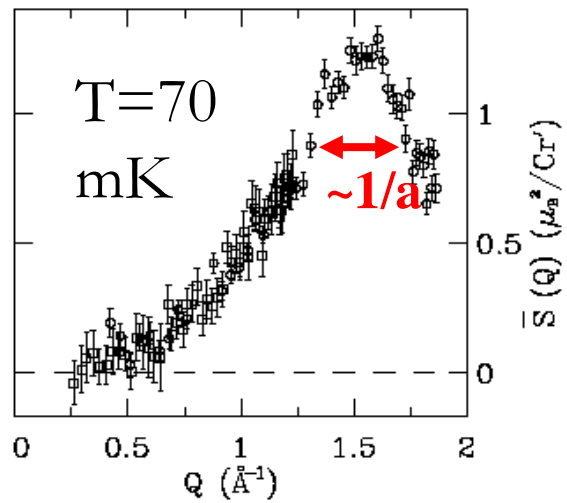
Kagomé bilayers (Cr^{3+} , $S = 3/2$)



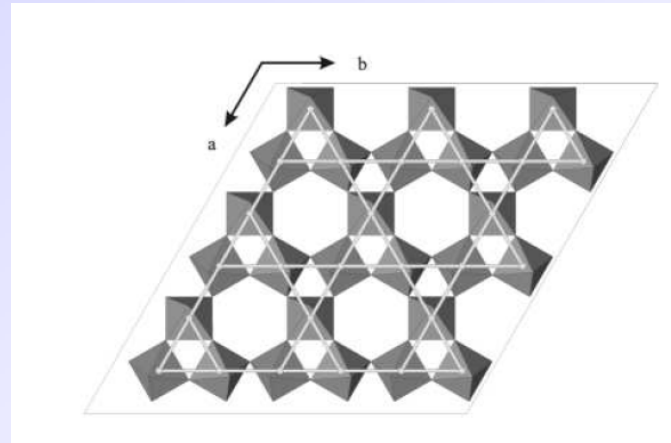
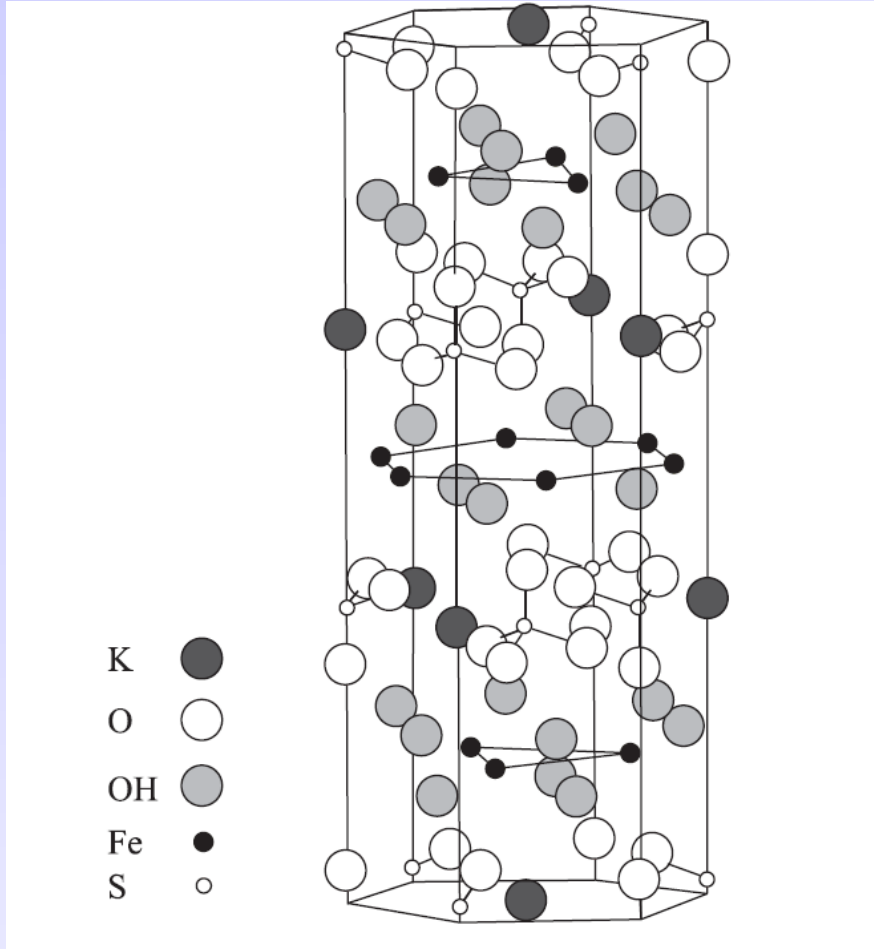
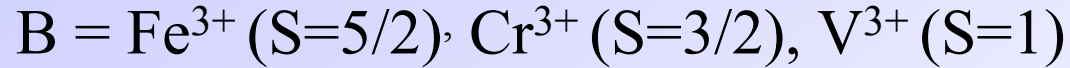
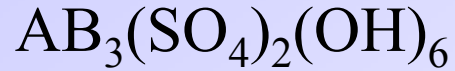
- Exotic spin glass at only low T
- High frustration ratio θ/T_g
- Fluctuations at low T
- Short correlation length
- Field independant C_v
- >3% spinless defects



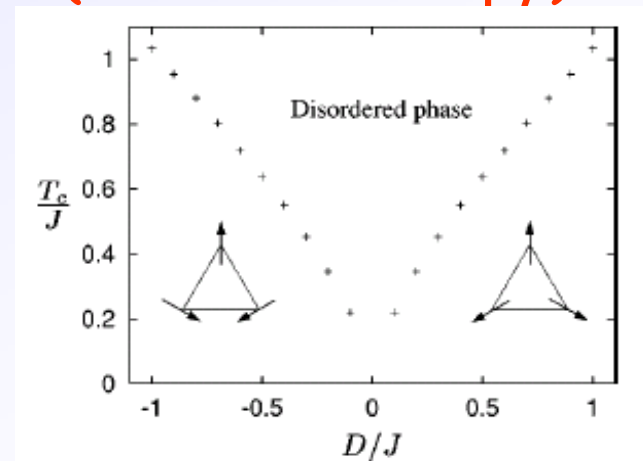
$\text{SrCr}_{9p}\text{Ga}_{12-9p}\text{O}_{19}$
 (SCGO)
 $\Theta \sim 500 \text{ K}$
 $T_g \sim 4 \text{ K}$



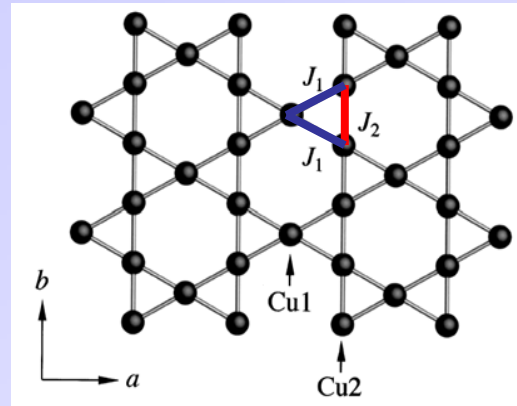
Jarosite Family



Most of Jarosites order!
(D.M. anisotropy)

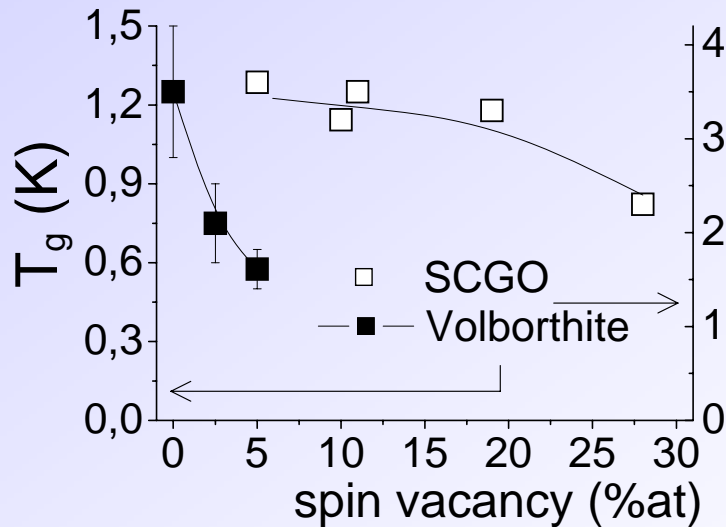


Volborthite, $\text{Cu}_3\text{V}_2\text{O}_7(\text{OH})_2 \cdot 2\text{H}_2\text{O}$



Z. Hiroi et al, J. Phys. Soc. Jpn 70 (2001) 3377

Another class of imperfections: J_1 and J_2



$S = 1/2$
High purity

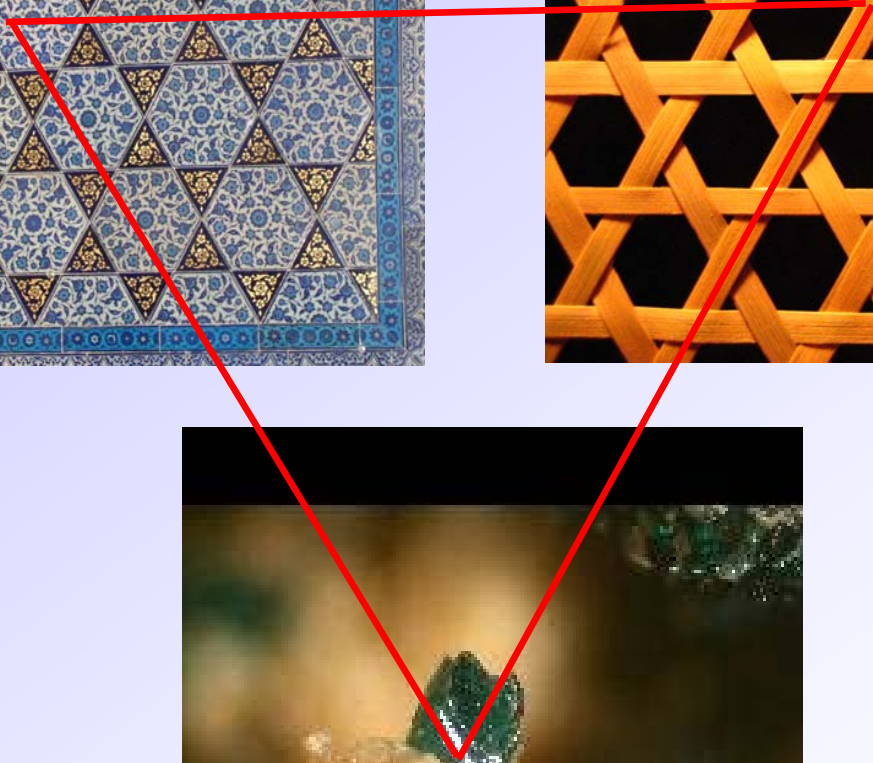
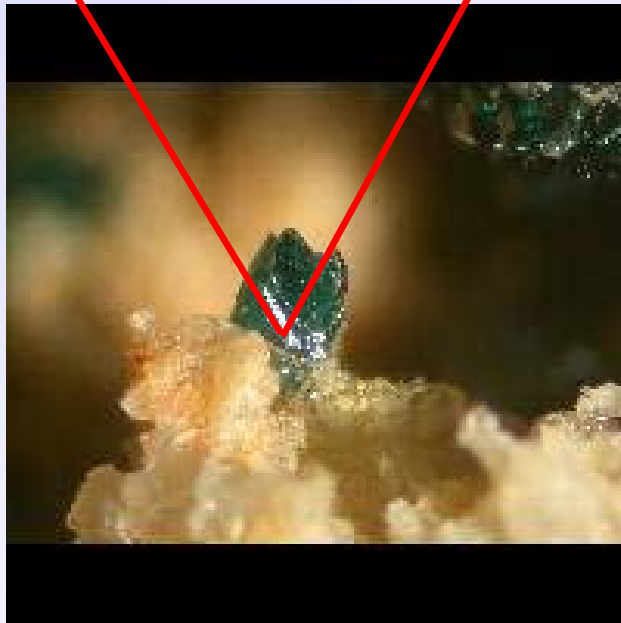
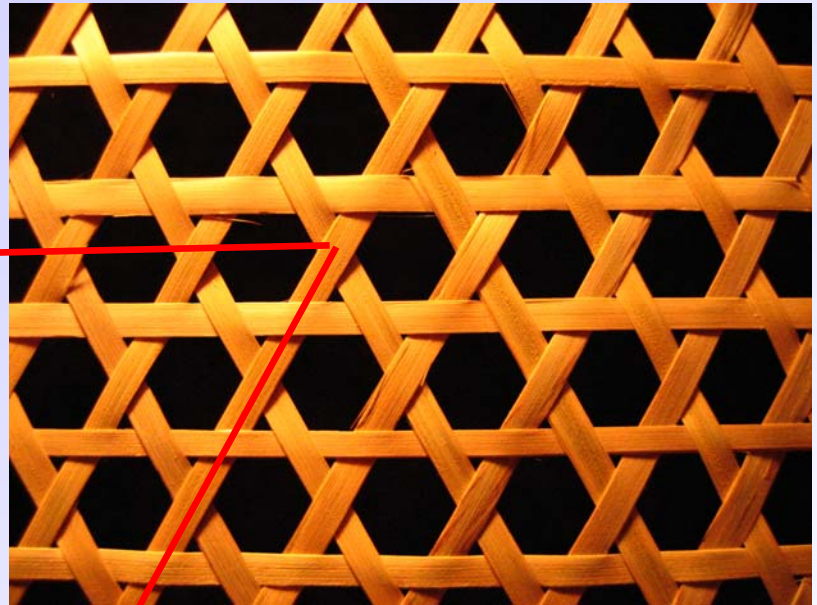
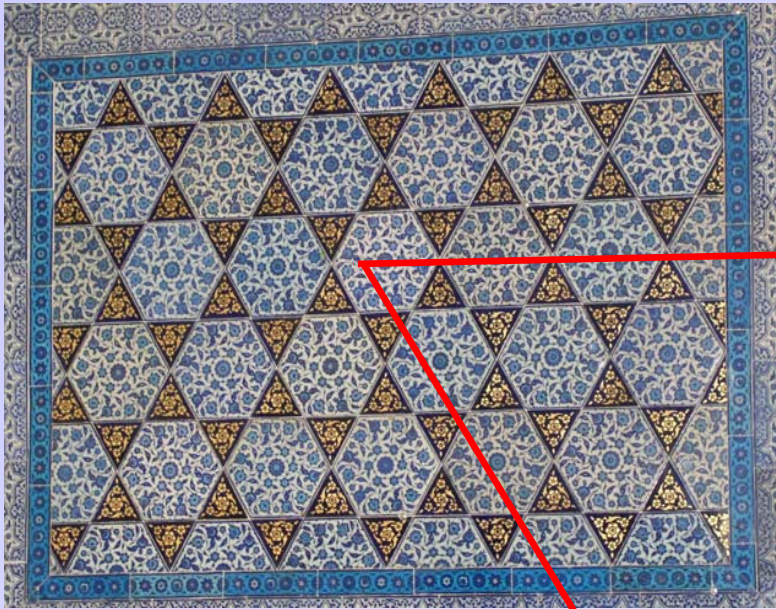
But a transition

Quantum Kagome Spin Liquids

Part II

Novel states induced by
frustration

Herbertsmithite: a true spin
liquid?



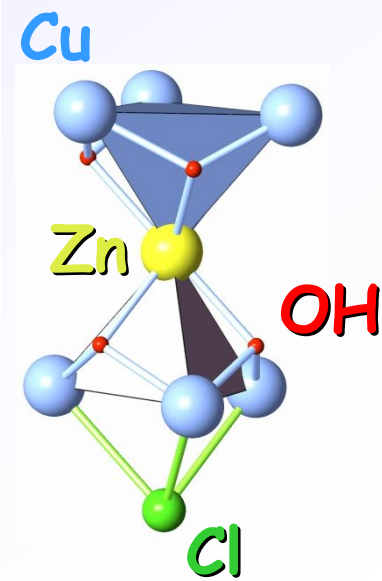
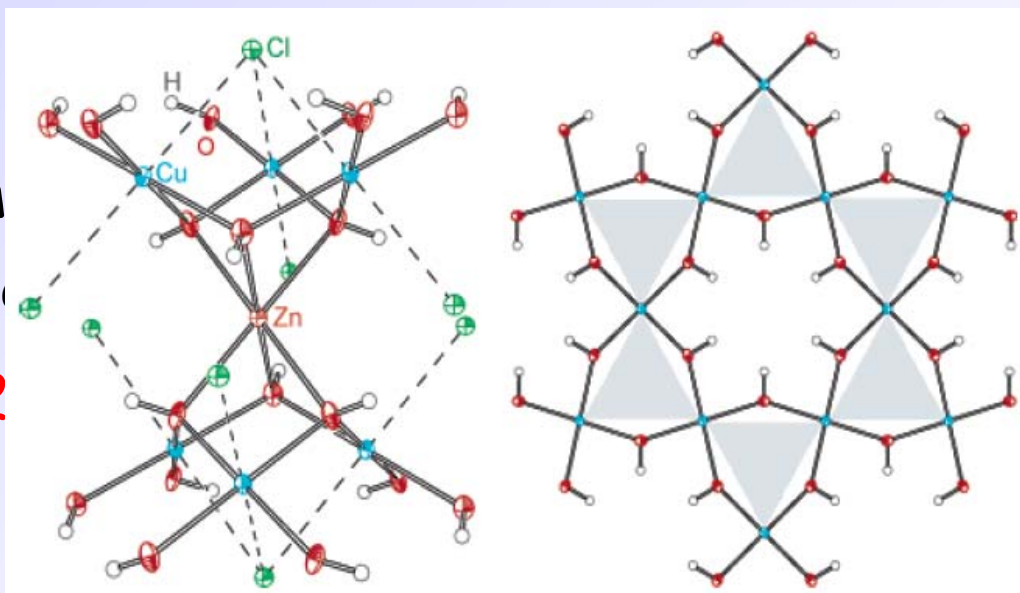
A Structurally Perfect $S = 1/2$ Kagomé Antiferromagnet

Matthew P. Shores, Emily A. Nytko, Bart M. Bartlett, and Daniel G. Nocera*

*Department of Chemistry, 6-335, Massachusetts Institute of Technology, 77 Massachusetts Avenue,
Cambridge, Massachusetts 02139-4307*

Received June 13, 2005; E-mail: nocera@mit.edu

Herbertsm
 $\text{ZnCu}_3(\text{OH})_4\text{Cl}_2$
 Cu^{2+} , $S=1/2$



Herbertsmithite is the first example of a quantum kagome antiferromagnet

- perfect kagomé lattice
- no freezing at least down to $J/4000$

-> renewal of the search for scenarios for the kagome ground state

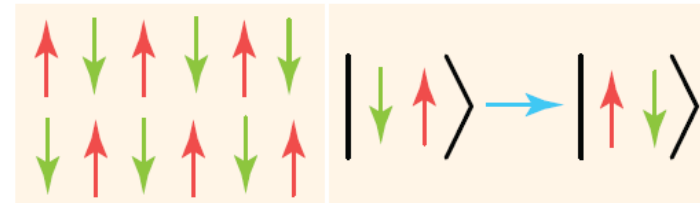
VBC : R.R.P. Singh and D.A. Huse, PRB (2007, 2008)

Dirac spin liquid : Y. Ran et al, PRL (2007), PRB (2008)

An End to the Drought of Quantum Spin Liquids

Patrick A. Lee

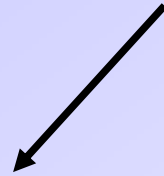
After decades of searching, several promising examples of a new quantum state of matter have now emerged.



Sciences, perspectives sept 2008

1- $\text{Zn}_x\text{Cu}_{4-x}(\text{OH})_6\text{Cl}_2$: paratacamite family $x < 1$

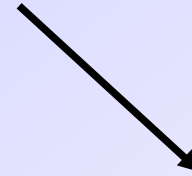
2- $\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$: Herbertsmithite: ideal kagome ?



Gap?

Measurements of χ_{local}

Dynamical measurements

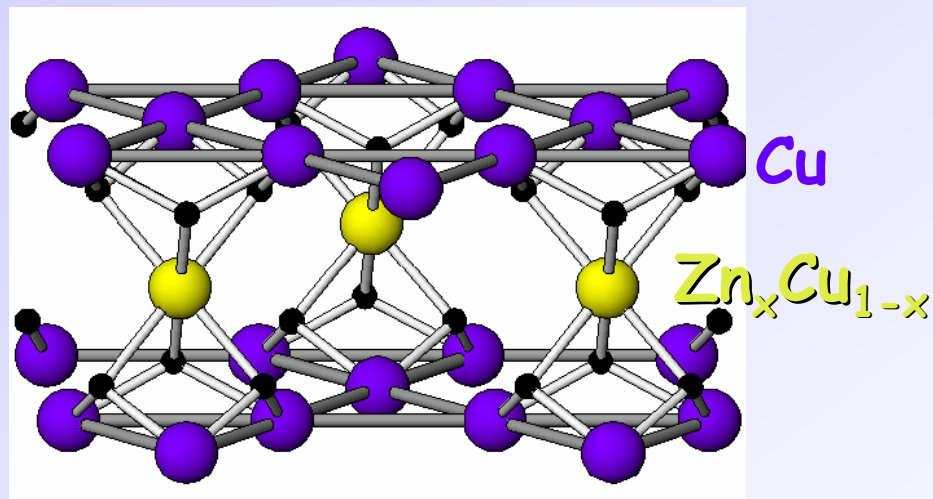


Non-magnetic defects

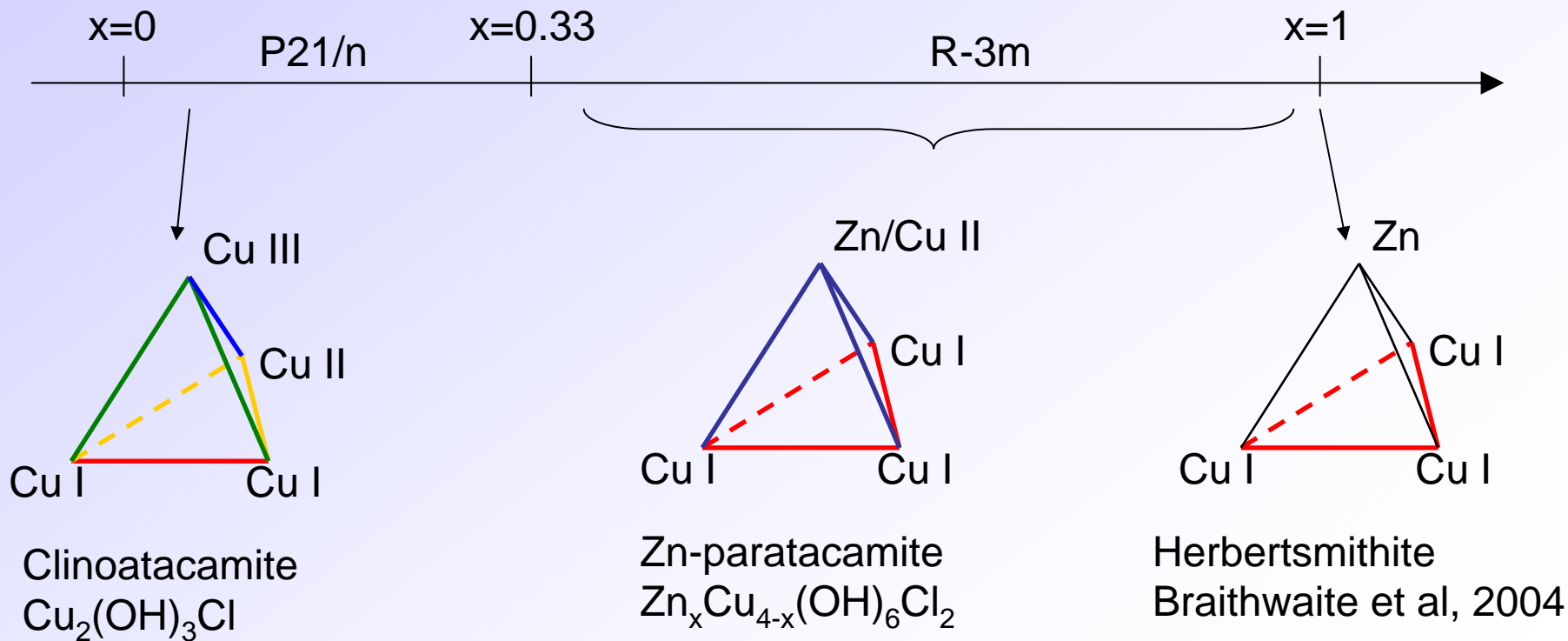
Cu/Zn

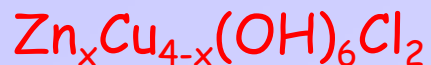
« Perturb to reveal »

$Zn_xCu_{4-x}(OH)_6Cl_2$
atacamite family

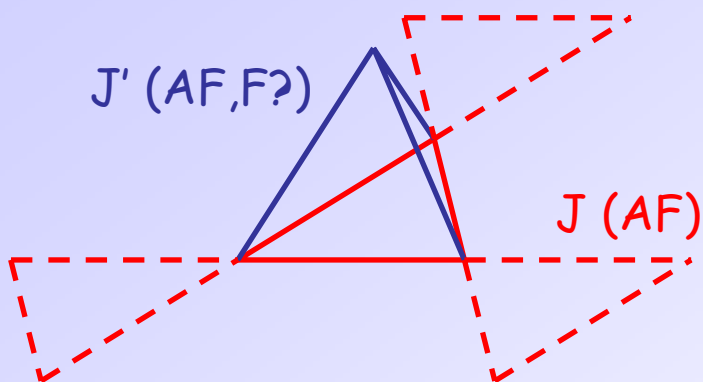


Zn/Cu substitution rate



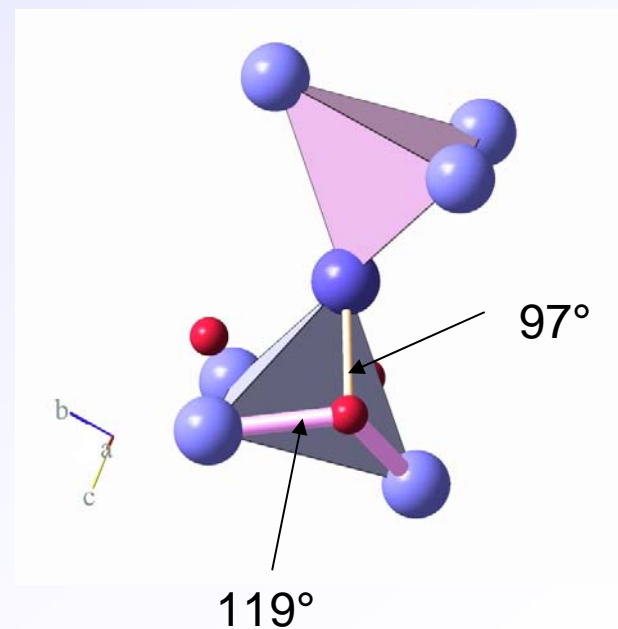
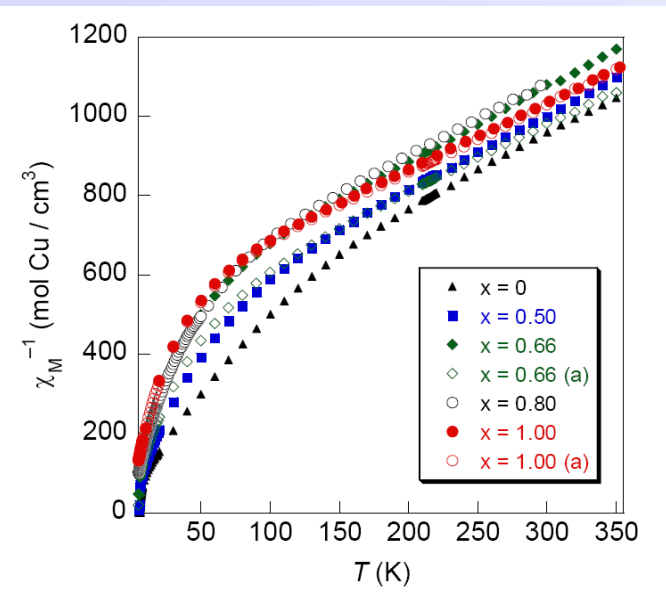


Curie Weiss behavior for all x
 → antiferromagnetic correlations



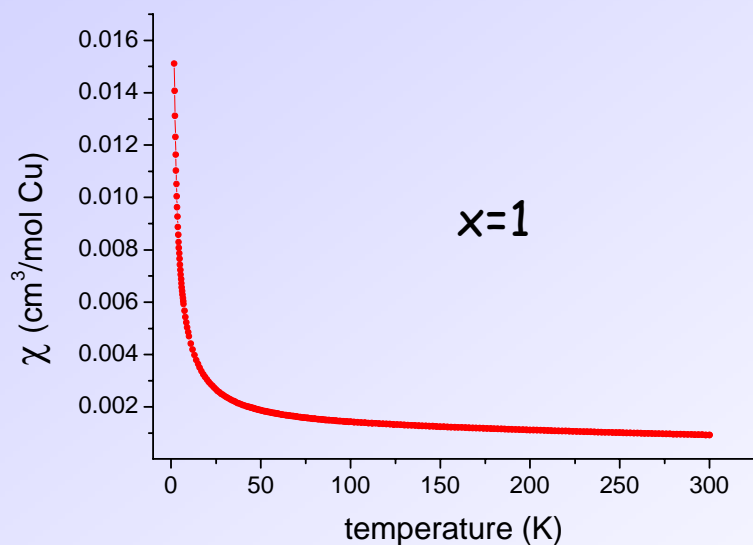
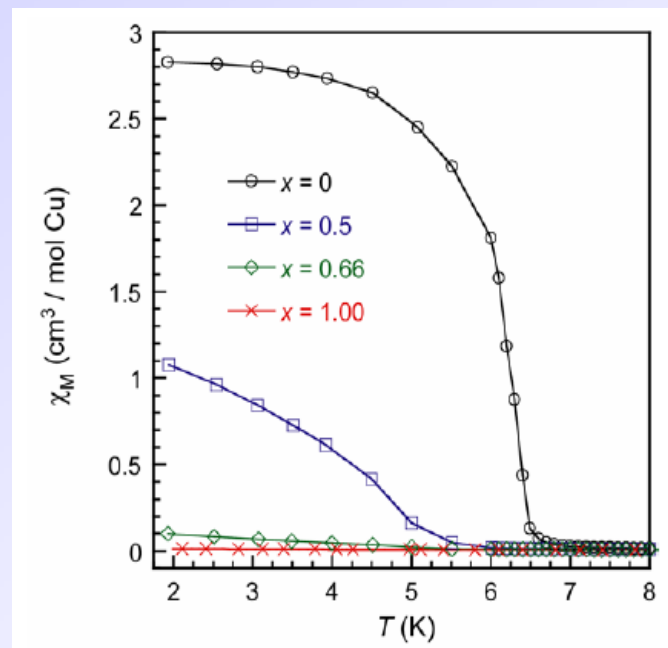
J : in-plane coupling AF ~ 175 K
 J' : inter-plane coupling small, maybe Ferro

no transition for $T \ll J$
 → highly frustrated antiferromagnets





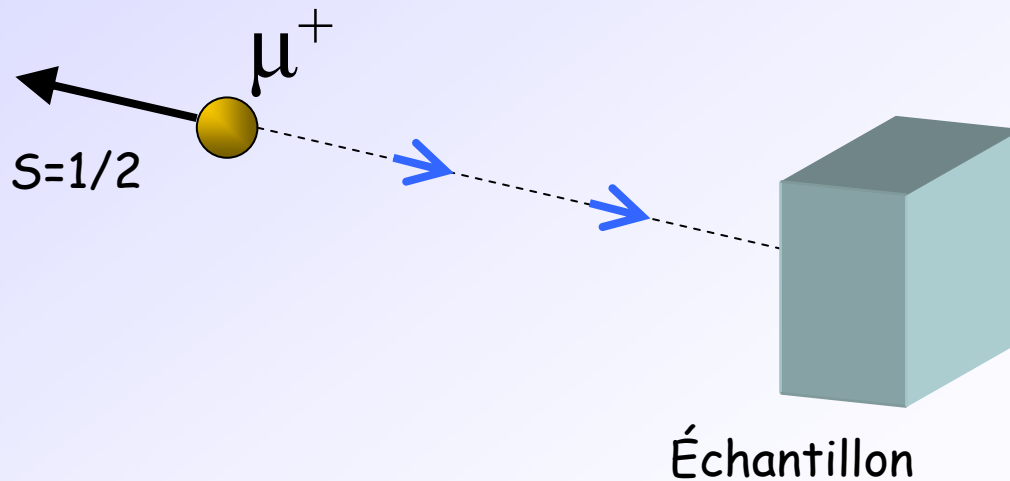
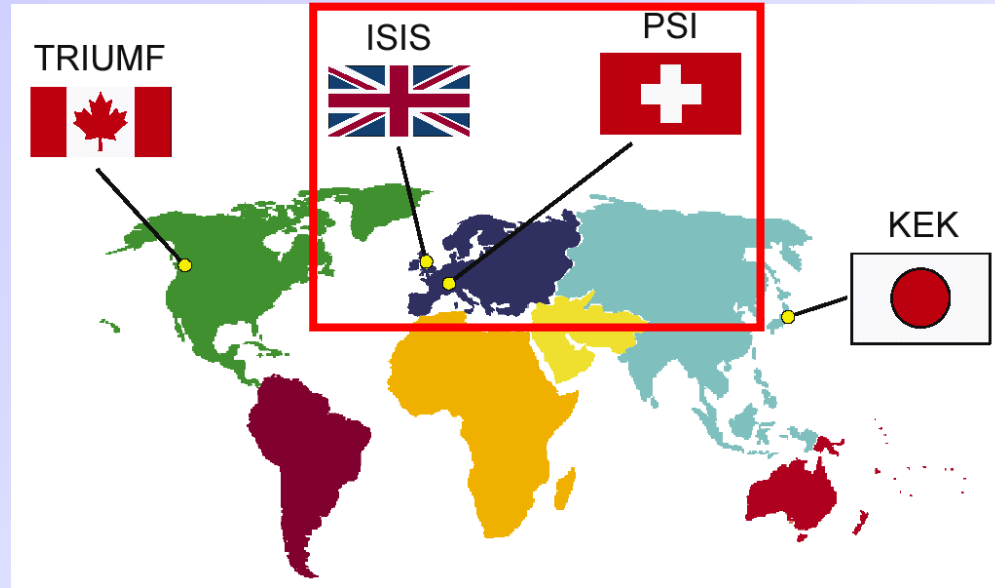
- Ferromagnetic-like transition at $T_N \sim 6\text{K}$
- Vanishes for $x \rightarrow 1$

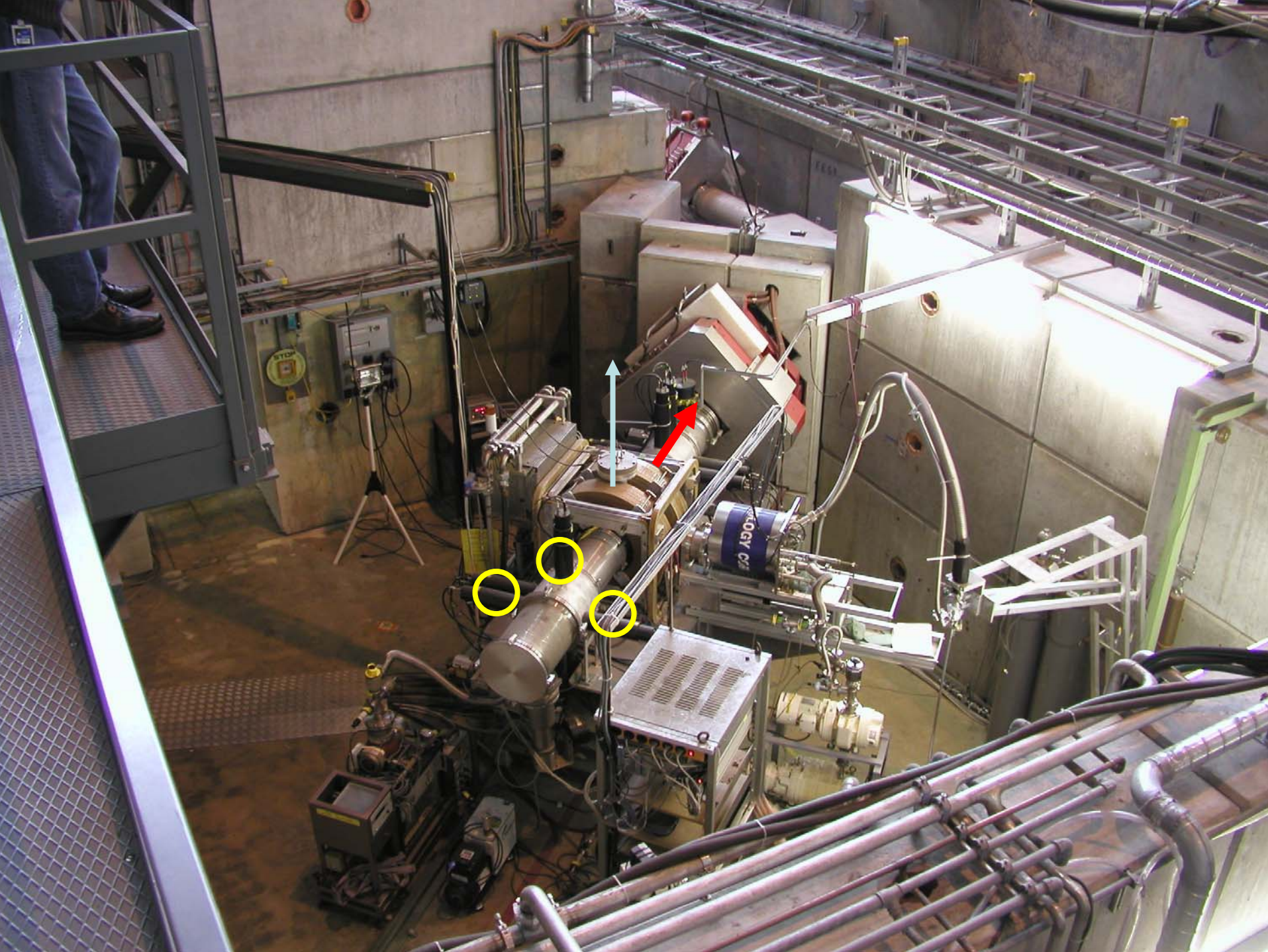


Herbersmithite $x=1$

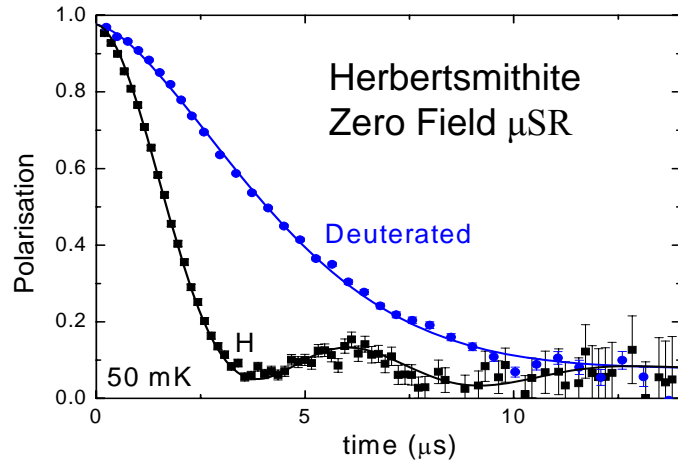
- No sign of transition for $T > 2\text{K}$
- Low T Curie like upturn for $T < 50\text{K}$

μ SR - Muon spin resonance, relaxation





$\mu\text{SR} : \text{ZnCu}_3(\text{OH})_6\text{Cl}_2, x=1$



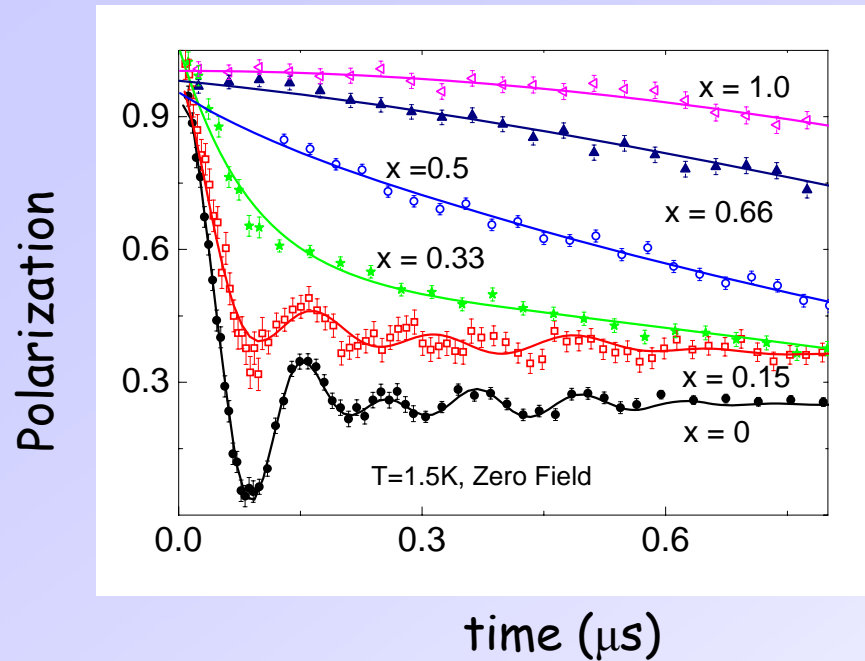
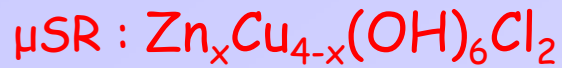
P. Mendels et al, PRL 98, 077204 (2007)

Also: ac- χ Helton et al, PRL 98 107204 (2007)
 μSR O. Ofer et al, cond-mat/0610540

No order or frozen disorder down to 50 mK despite $J=175$ K !

At 50mK, no sign of ordering relaxation arises from small “static” nuclear fields.

upper limit of a frozen moment for Cu^{2+} , if any : $6 \times 10^{-4} \mu_{\text{B}}$

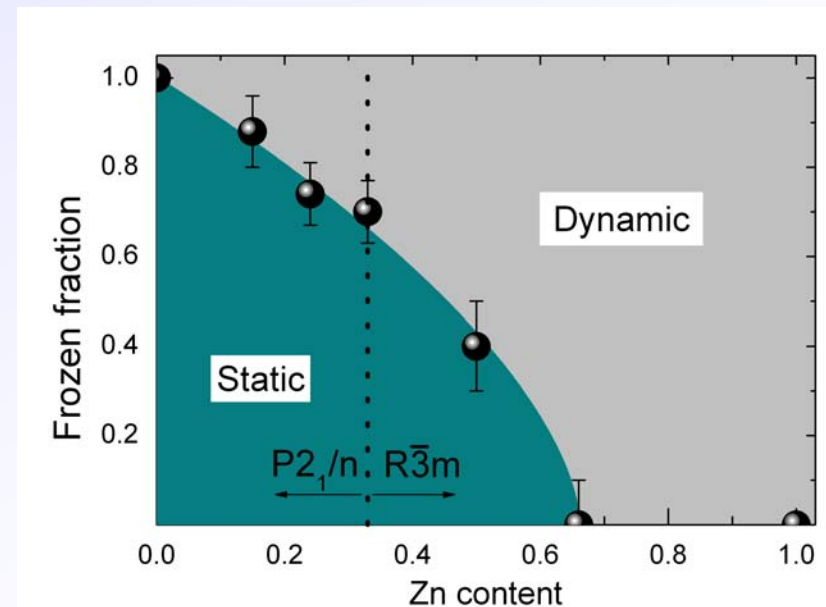


-x=0 : fully ordered below ~18K
 X.G. Zheng et al, PRL 95, 057201 (2005)

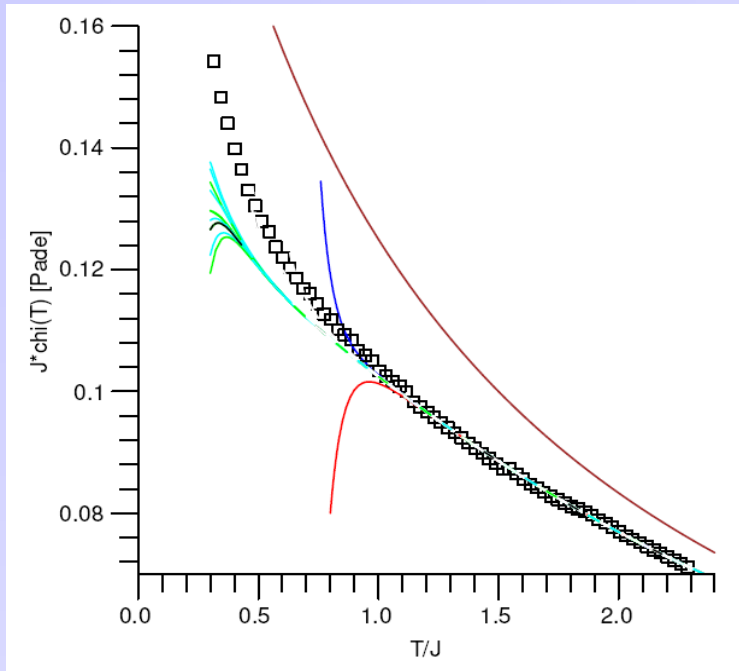
When x increases from 0 to 1 :
 -Oscillations are smeared out
 -A paramagnetic (x=1 type) component emerges at the expense of the frozen one

Large domain of stability $0.66 < x < 1$ of a dynamical ground state
 -> surprisingly small influence of interlayer coupling

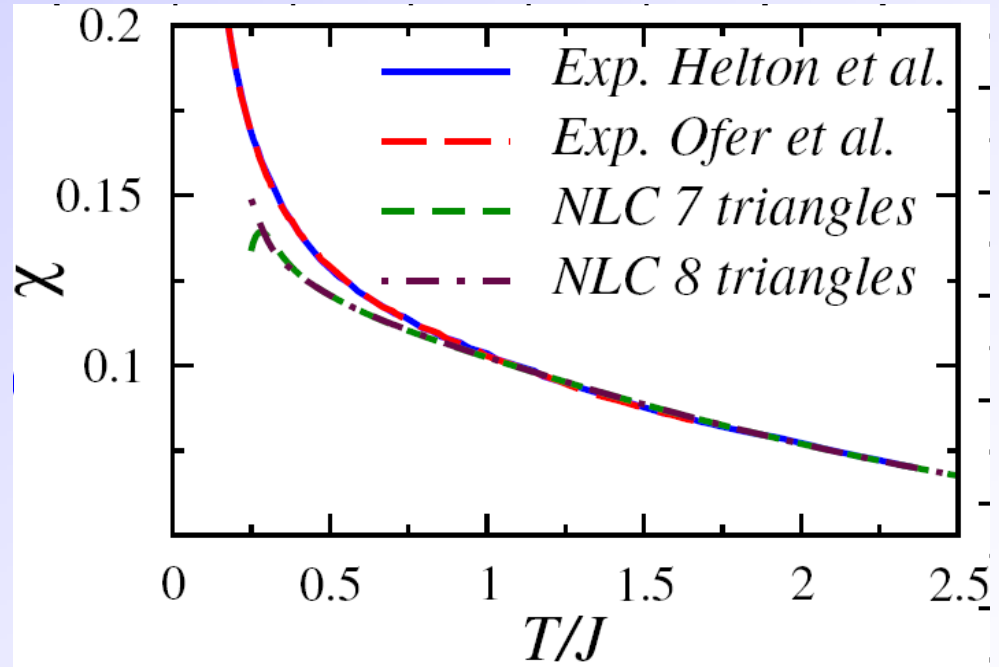
P. Mendels et al, PRL 98, 077204 (2007)



Herbertsmithite macroscopic susceptibility



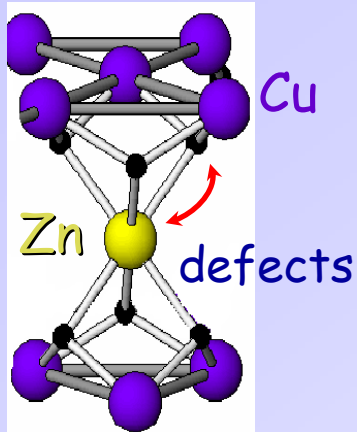
High temperature series expansion
(G. Misguish, C. Lhuillier)



M. Rigol and R. P. Singh, PRL **98**, 207204 (2007)

-> need for additional terms to the KAF Hamiltonian to account for Herbertsmithite macroscopic susceptibility

Magnetic defects : Zn/Cu intersite mixing



Cu on the Zn site
Nearly free $\frac{1}{2}$ spins

Zn in the kagome plane

-> magnetic vacancy -> staggered magnetization
-> effective paramagnetic defects (small moment?)

P. Mendels et al., J. Phys.: Condens. Matter 19, 145224 (2007)

Neutron: structure refinement + H

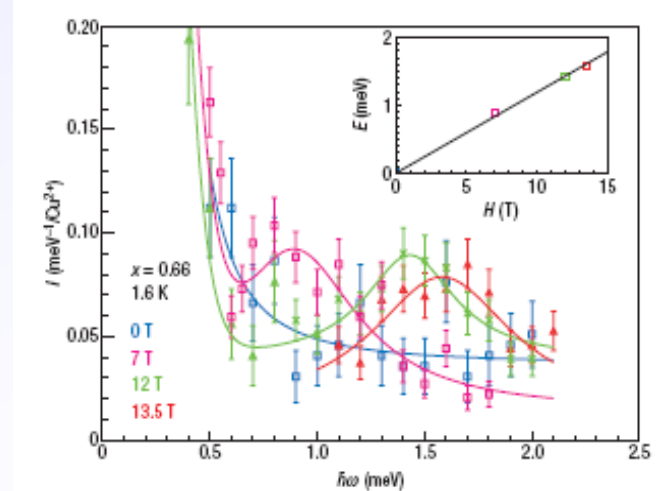
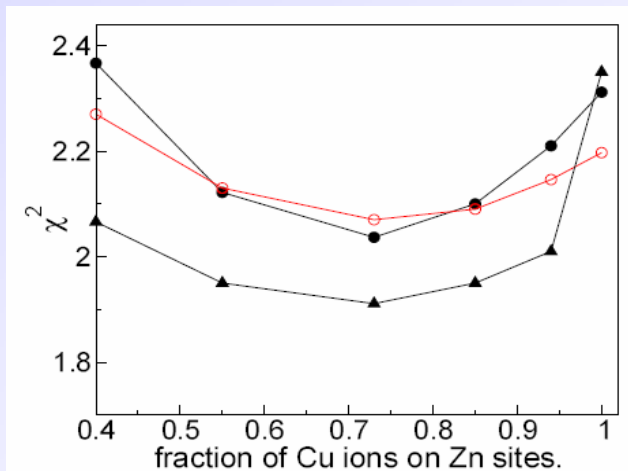
S.H. Lee et al, Nature Materials (2007)

M.A. de Vries et al, PRL 100, 157205 (2008)

dilution

-> ~10%

-> 9(2)%

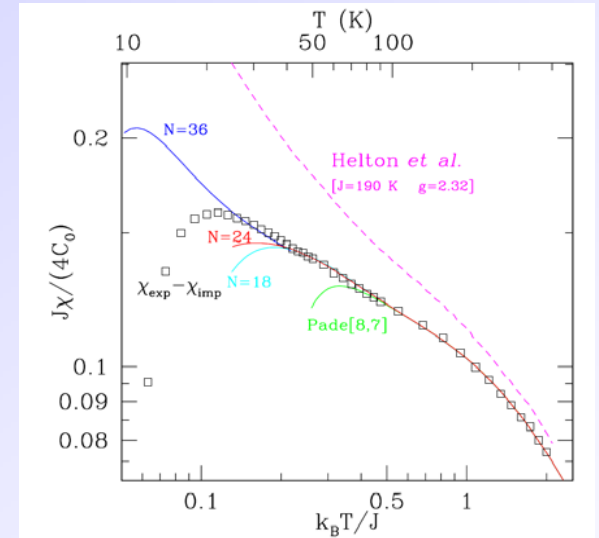


Magnetic defects : Zn/Cu intersite mixing

- **Susceptibility fit** → ~5% dilution

exact diagonalization+
5% weakly interacting $S=1/2$ defects

G. Misguich and P. Sindzingre, Eur. Phys. J. B 59, 305 (2007)



- **Low T, High Field Magnetization** → ~7% dilution

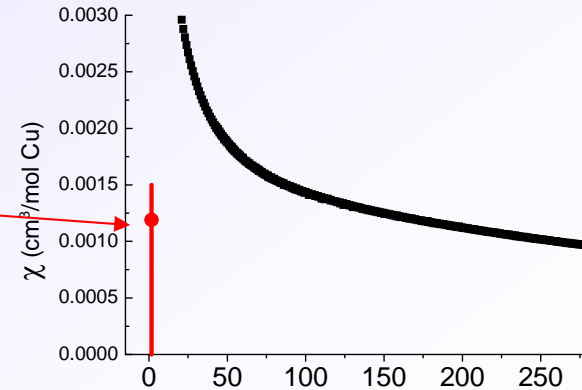
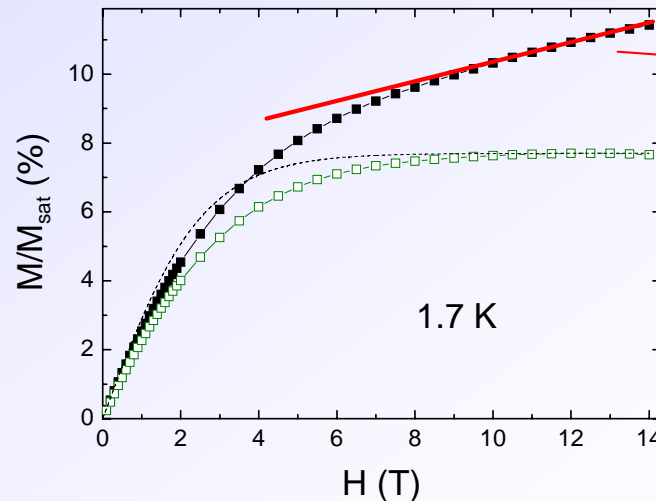
F. Bert et al, PRB 76, 132411 (2007)

$$M_{\text{tot}} = M_{\text{defect}} + \chi_i H$$

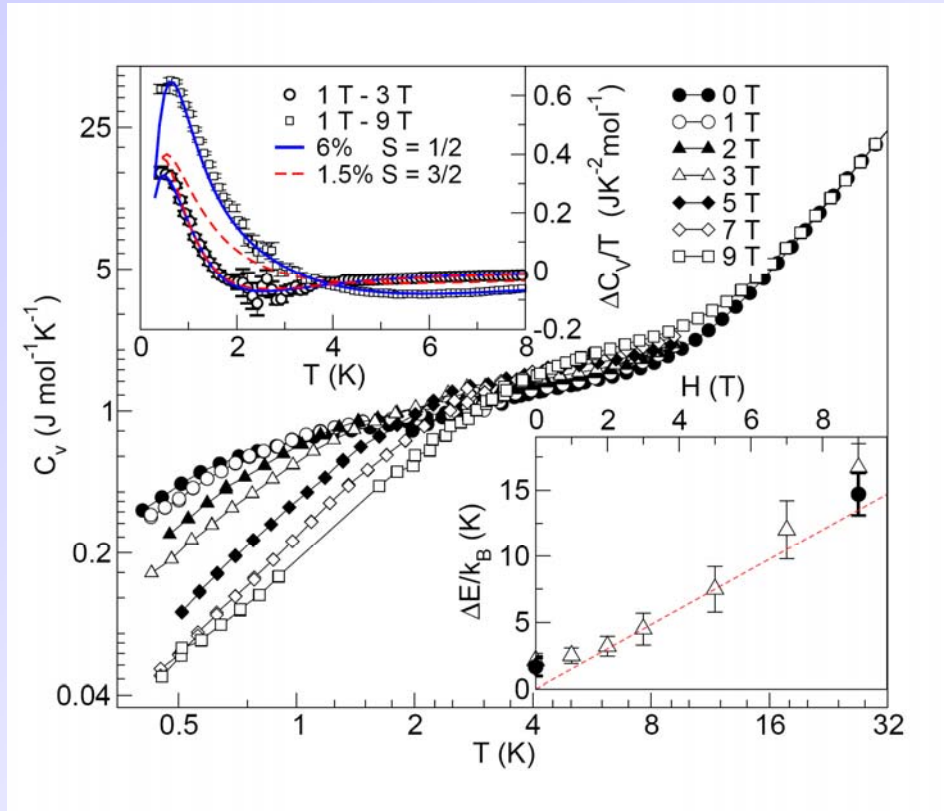
$$M_{\text{defect}} \sim \text{Brillouin}(H/T)$$

- ~7% of interlayer Cu^{2+}

- $\chi_i \ll \chi_{\text{macro}}$ at low T



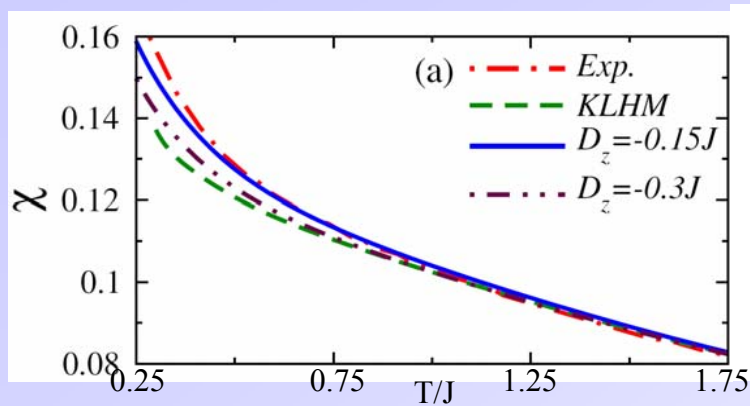
○ Schottky like anomaly in heat capacity



-> Kagome lattice dilution 6.5+/-1%

M.A. de Vries et al, PRL **100**, 157205 (2008)

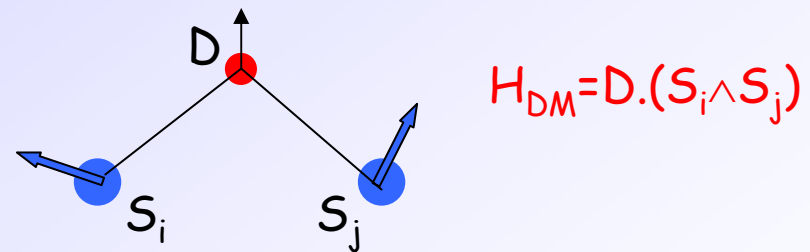
Dzyaloshinskii-Moriya interactions



Direct fit of susceptibility (no defect contribution):

$$|D_z| = 0.15J, |D_p| = 0.3J$$

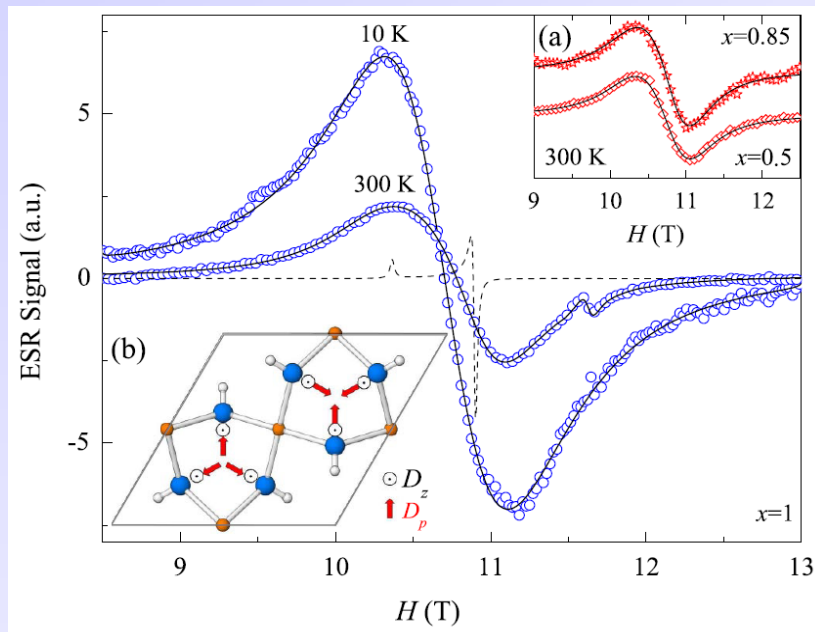
M. Rigol and R. P. Singh, PRL **98**, 207204 (2007)



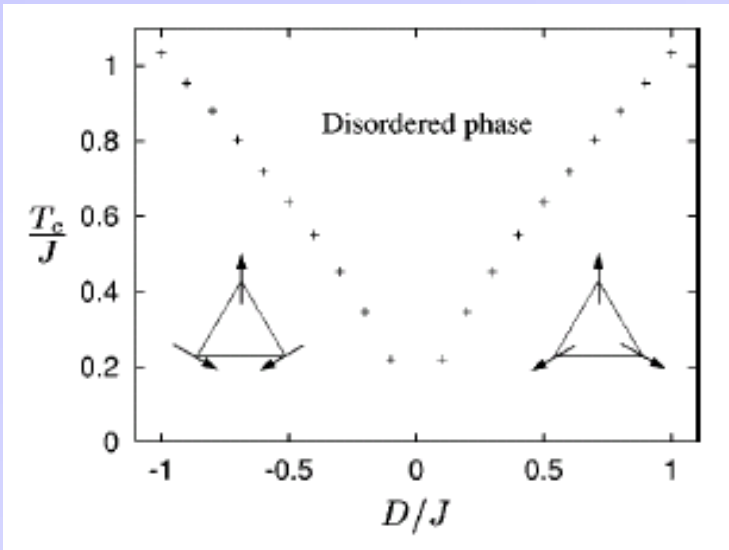
Broad room T ESR line \leftarrow magnetic anisotropy from DM

$$|D_z| = 0.08J, |D_p| \sim 0.01J$$

A. Zorko et al, PRL **101**, 026405 (2008)



Dzyaloshinskii-Moriya interactions

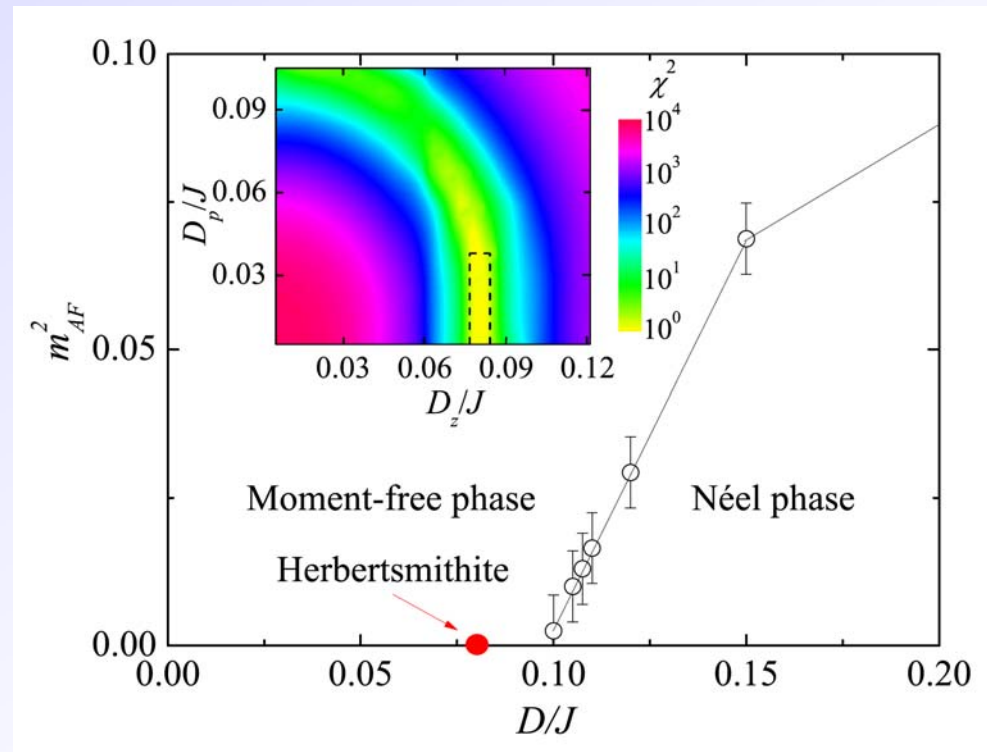


For classical spins, DM stabilizes ordered phases (cf jarosites)

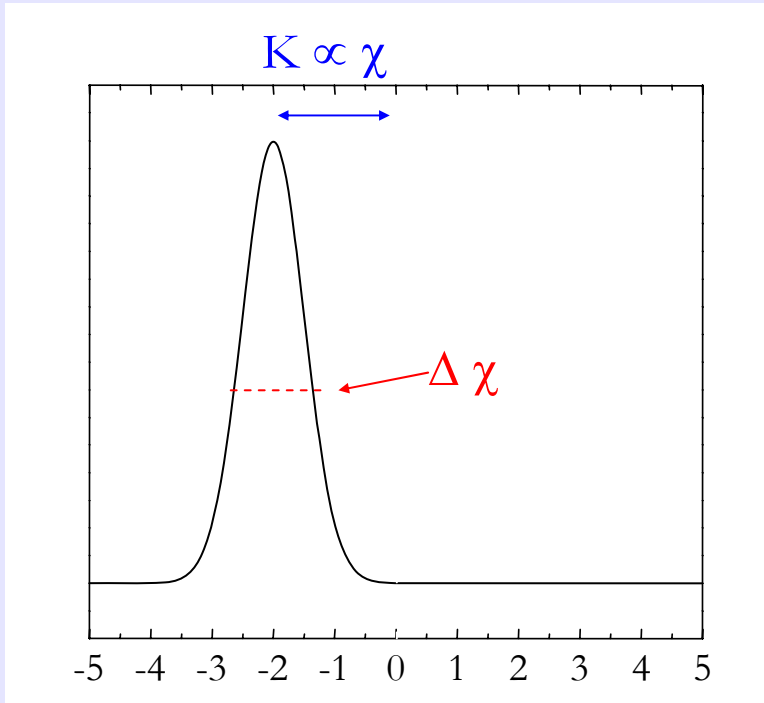
M. Elhajal et al, PRB **66**, 014422 (2002)

In the quantum case,
a moment free phase
survives up to $D/J \sim 0.1$

O. Cepas et al, PRB **78**, 140405 (R) (2008)



NMR: principles



Line shift K
susceptibility χ_{frustr}

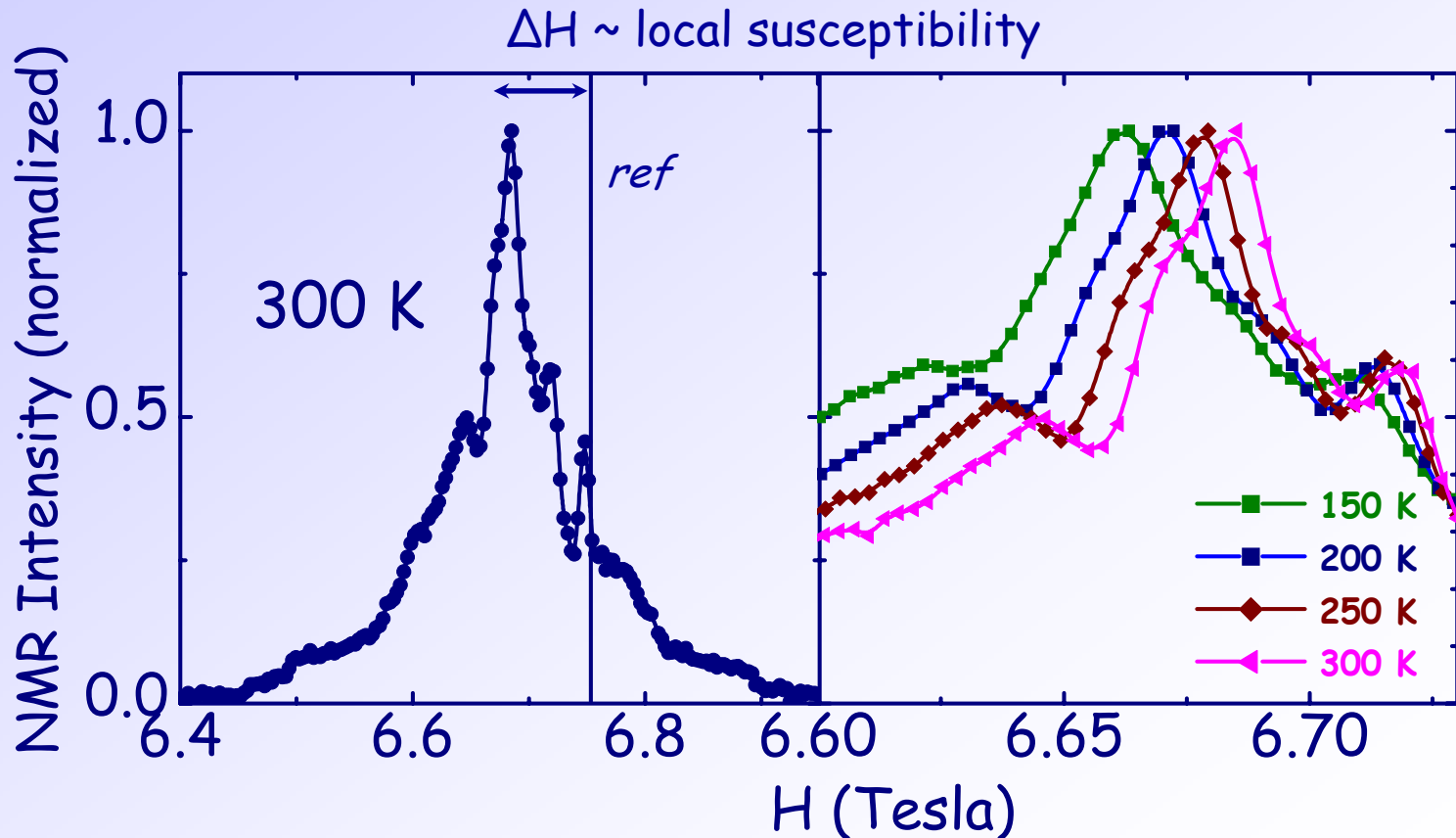
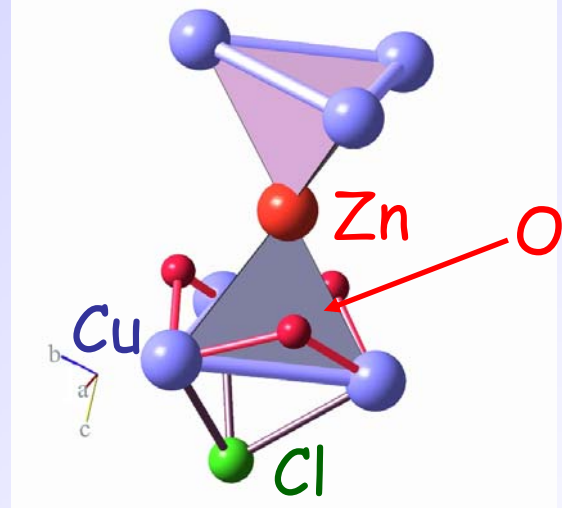
$$H = H_Z + A \vec{I} \cdot \vec{S}$$

Linewidth ΔH
spatially inhomogeneous
susceptibility (dilution)

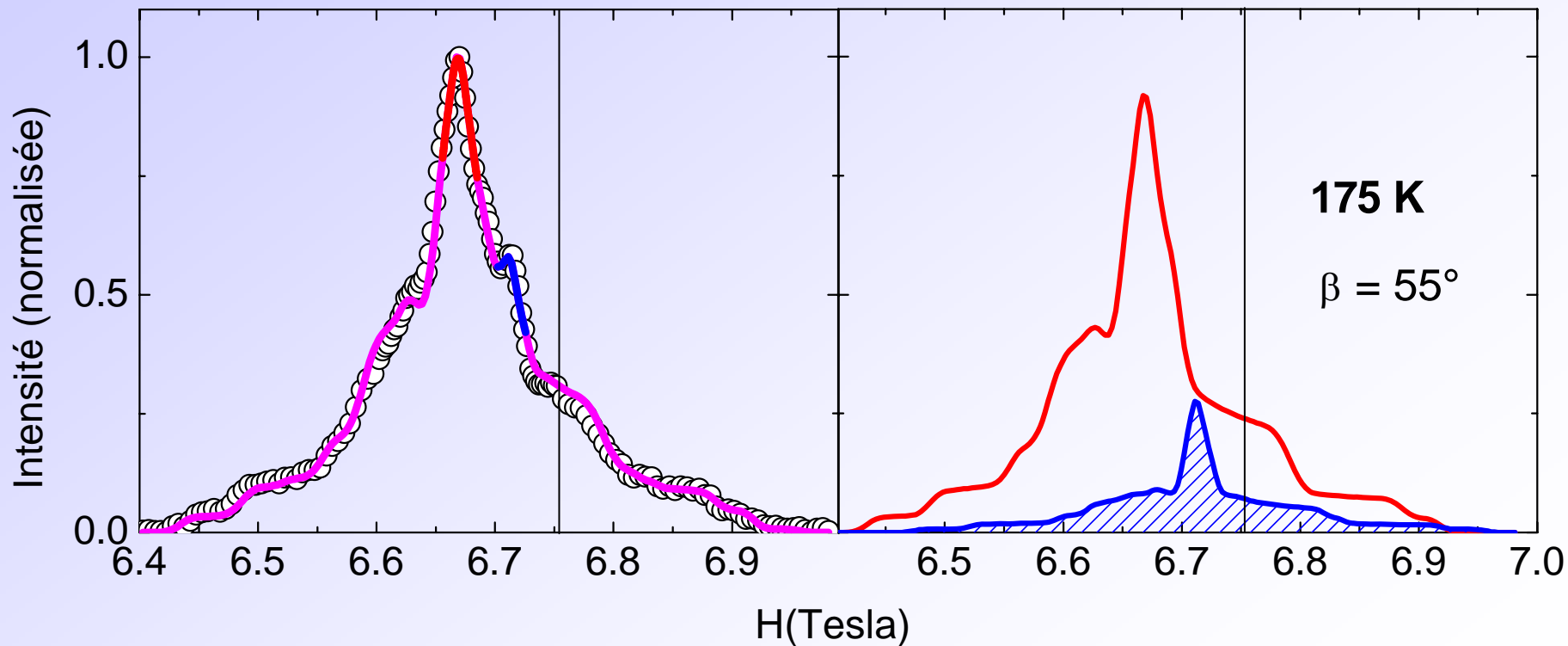
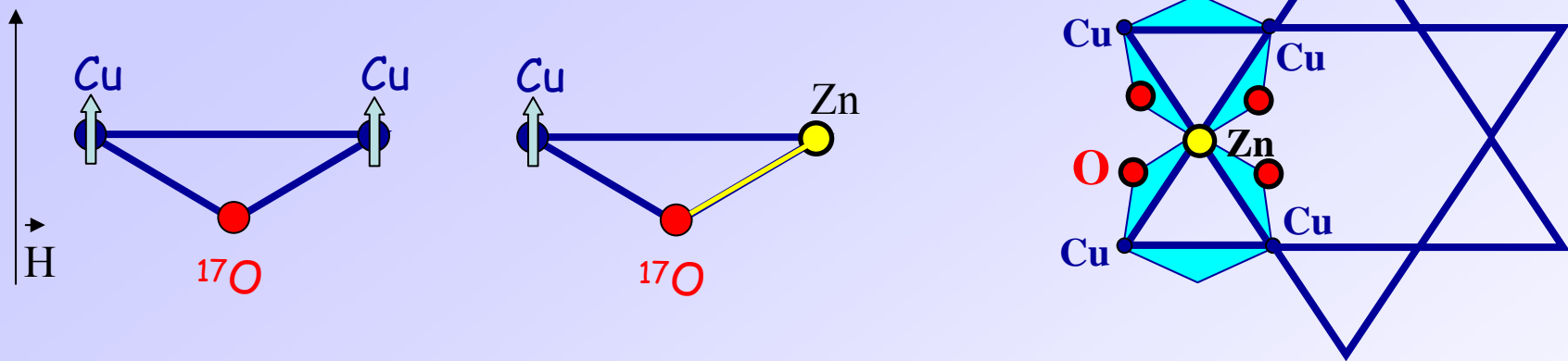
^{17}O NMR, local susceptibility

^{17}O : coupled to two Cu of the kagome plane

$I = 5/2 \rightarrow$ quadrupolar effects

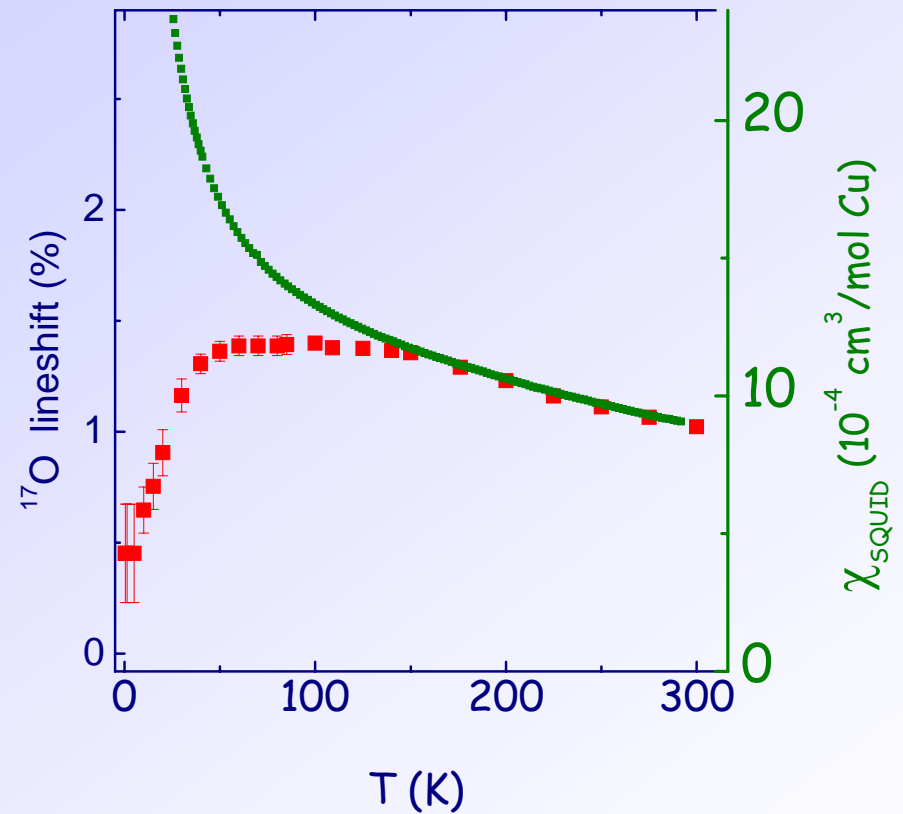
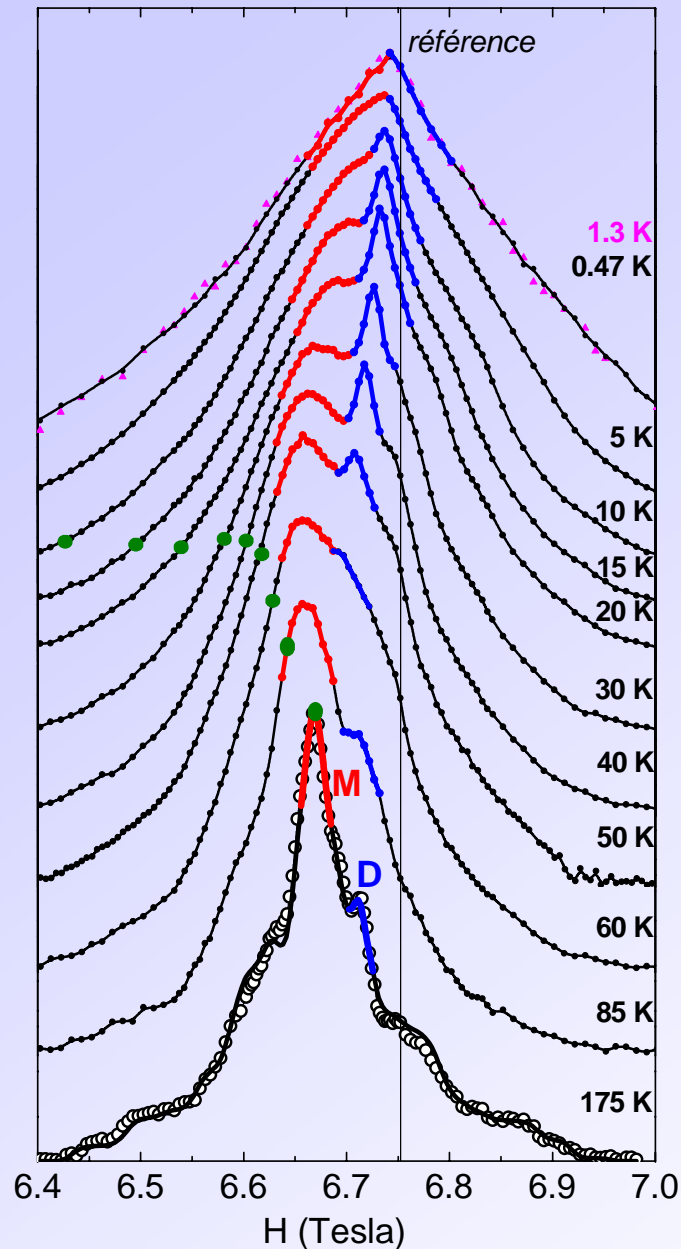


two magnetic site
O next to a Zn defect in the kagome plane



~ 20% intensity \rightarrow ~ 5% Zn/Cu defects in kagome planes

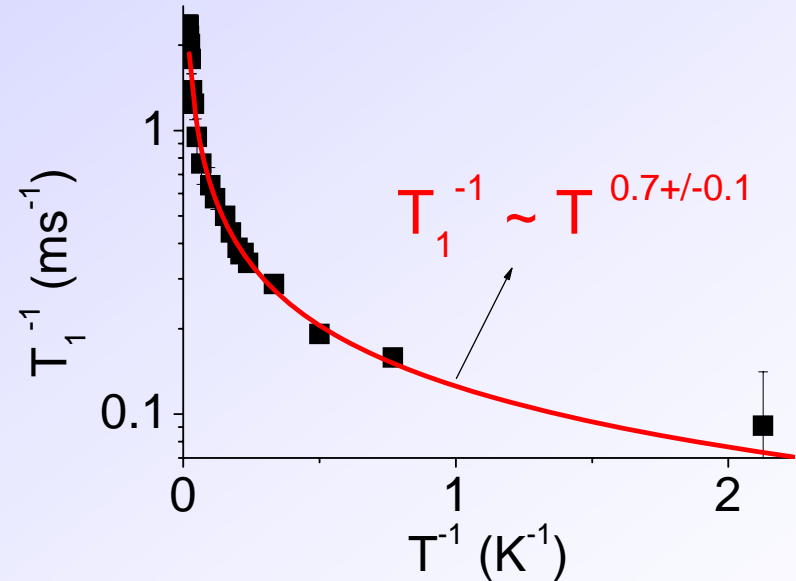
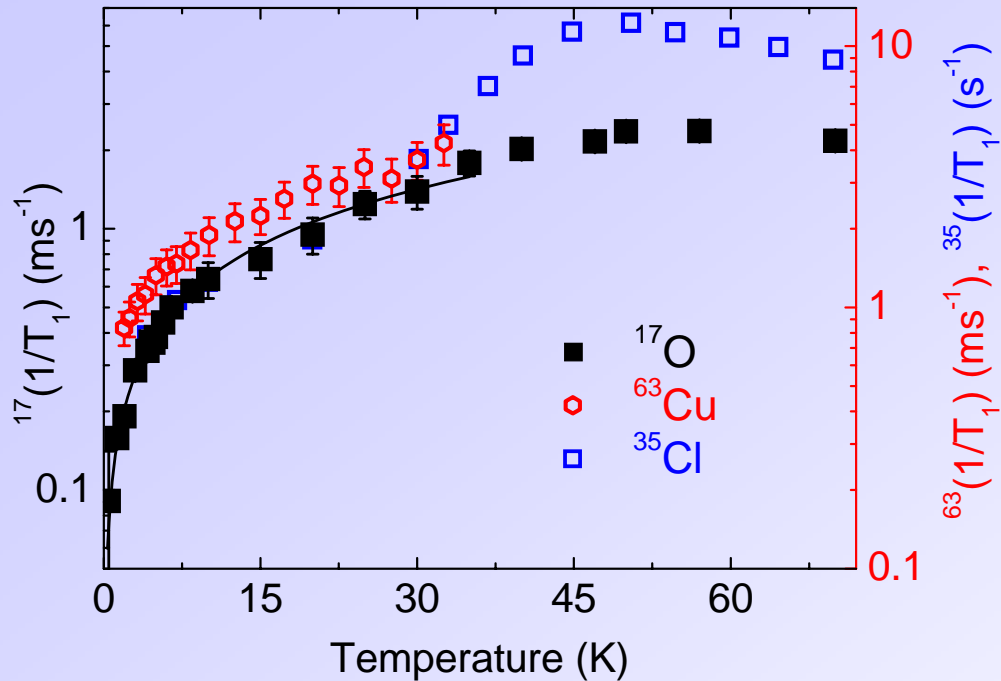
Main line, susceptibility of the kagome planes



-Susceptibility decreases below 50 K
-> enhancement of short range correlations
-> new energy scale

-No gap and Finite $T \rightarrow 0$ susceptibility
intrinsic or field effect, DM ?

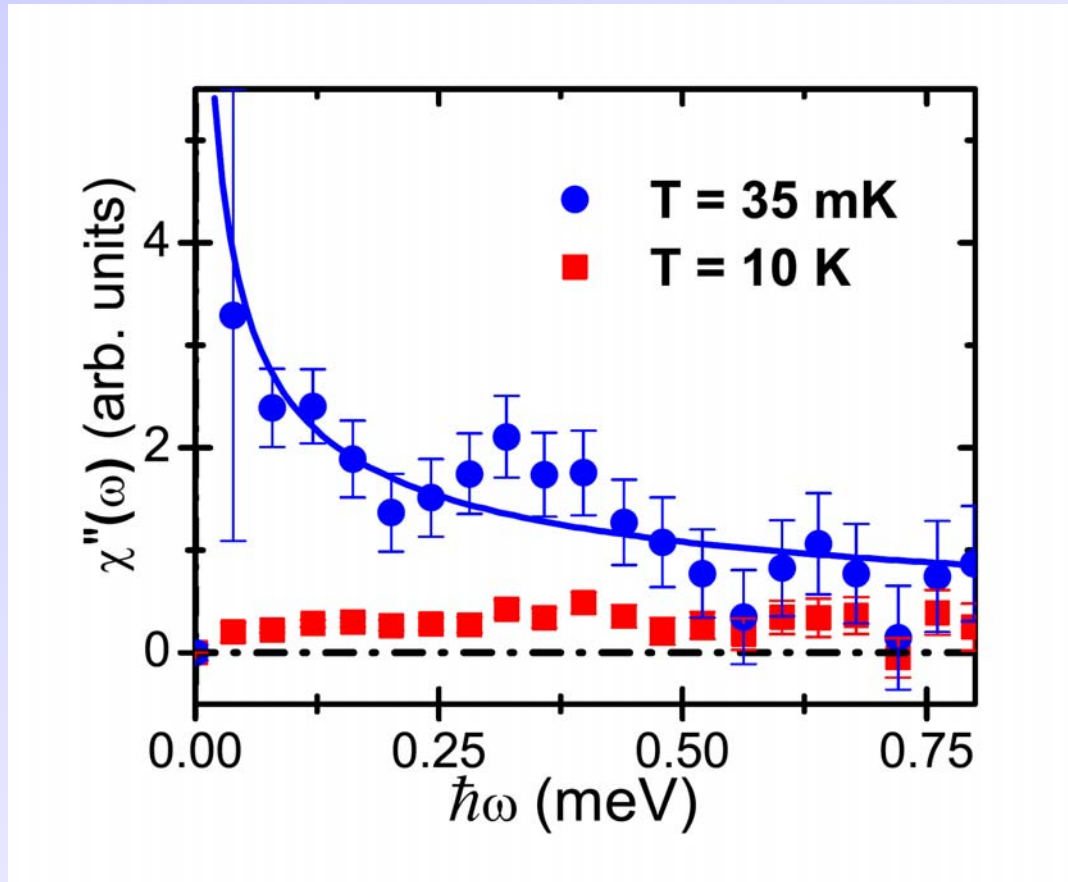
Low T Dynamics (NMR: T_1)



^{63}Cu and ^{35}Cl NMR from T. Imai et al, PRL 100, 077203 (2008)

- No spin gap behavior
 - > in agreement with absence of a gap $> 0.1 \text{ meV}$ in INS : Helton et al, PRL 98 107204 (2007)
- original sub-linear T dependence

Low T Dynamics (neutrons: χ'')



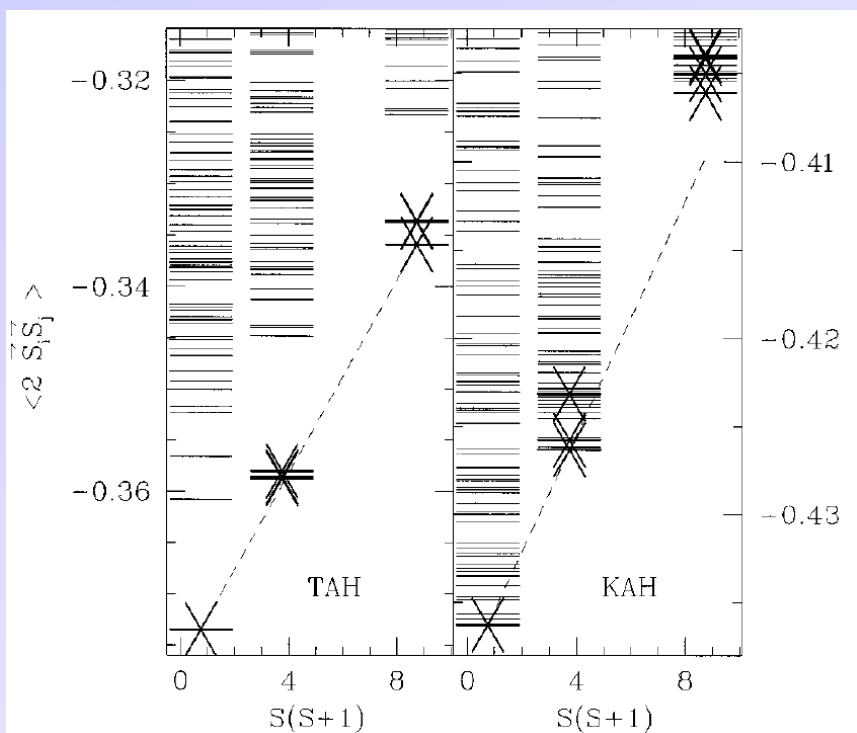
$$\chi''(\omega) \sim \omega^{-0.7(3)}$$

Helton et al, PRL 98 107204 (2007)

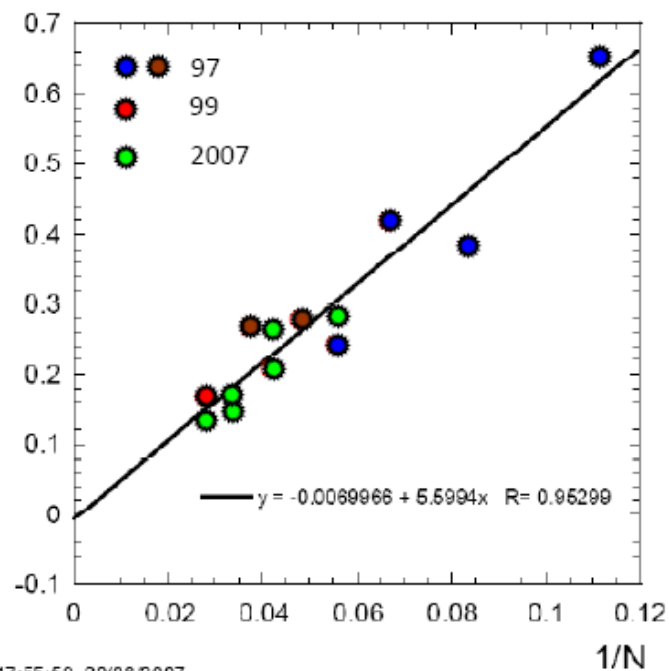
Exact Diagonalization

- Lecheminant, PRB **56**, 2521 (1997)
- Waldtmann *et al.*, EPJB **2**, 501 (1998).

- $\Delta < J/20$ gap between singlet ground state and 1st triplet state if any...



Gap KAH (all data) Ph. Sindzingre et C.L. 2007



Data 3 17:55:59 20/03/2007

$$\text{Gap} = 0 \pm 0.027$$

- no gap in the singlet sector.

No gap for Herbertsmithite!

- No gap (?) in exact diagonalizations (ED)
- $\chi(T \rightarrow 0)$ well reproduced in ED

Could be an intrinsic property

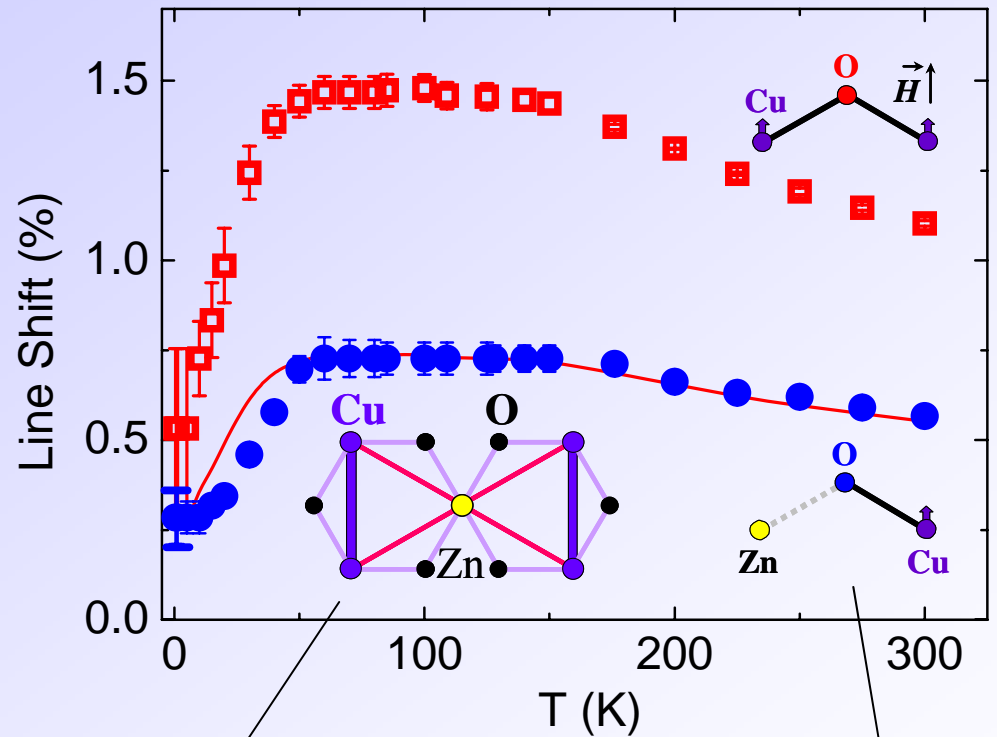
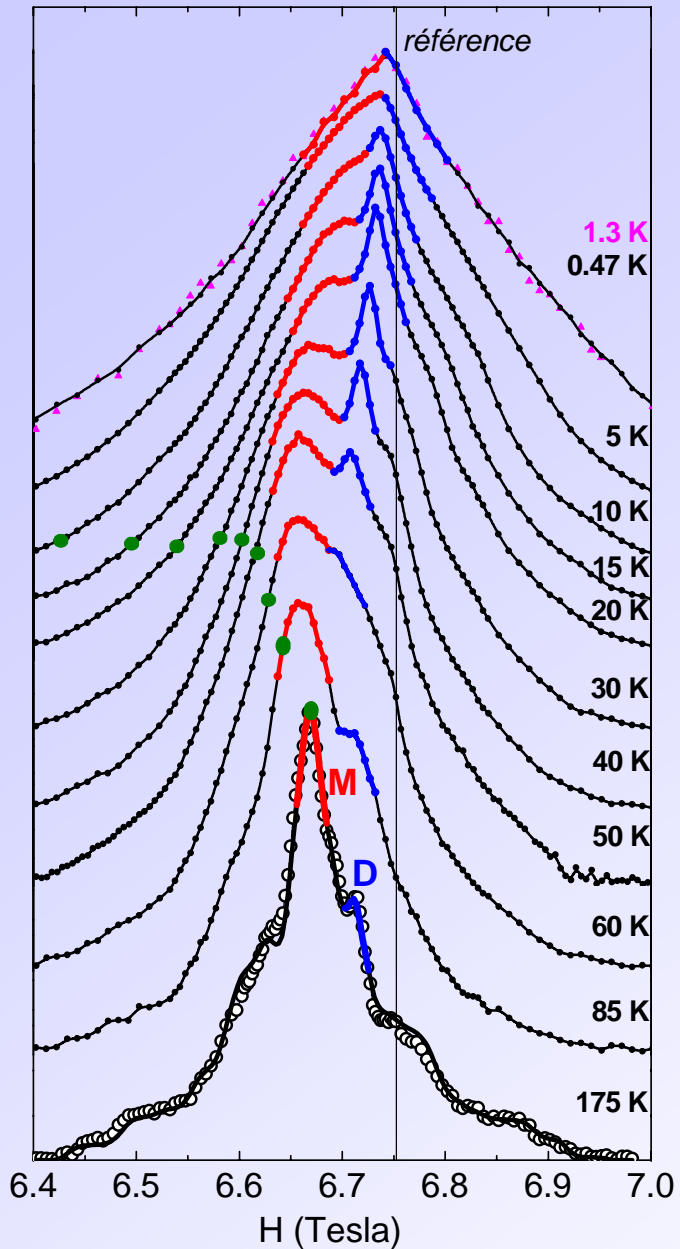
- DM interaction mixes singlet and gapped triplet and could restore a susceptibility

Could be an extrinsic property but one should also explain dynamical properties

Two classes of models

- Valence Bond Crystals with fluctuations of quantum dimers but forming an ordered pattern: a small singlet - triplet gap
- RVB-type spin liquids: fractional excitations (spinons), e.g. algebraic spin liquids
 - . Algebraic decrease of correlations
 - . no gap
 - . Spinons are bound (singlet - spinon interact)

Defect line



~30% oxygen sites
dimer localization around
a spin vacancy

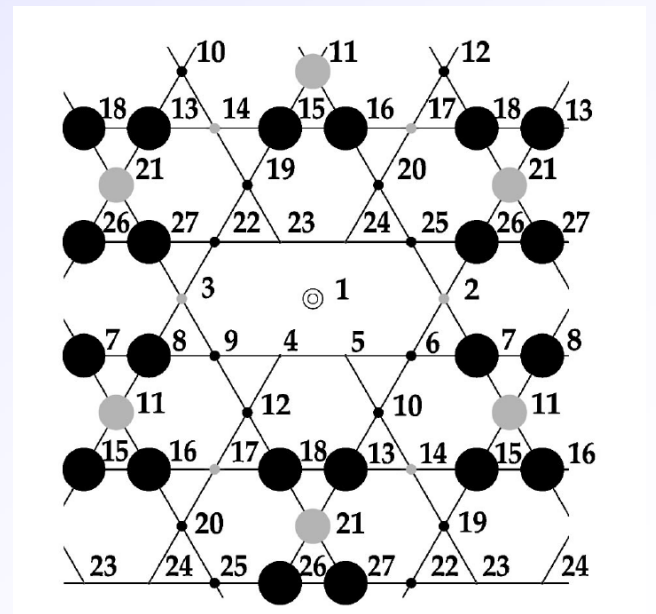
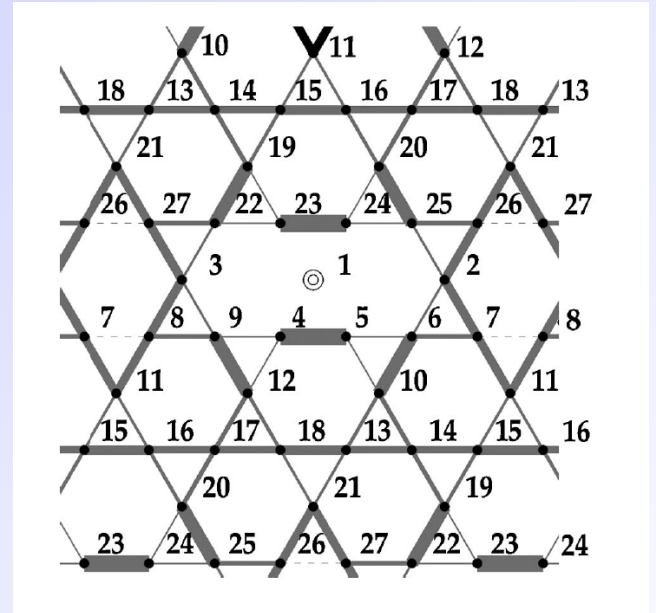
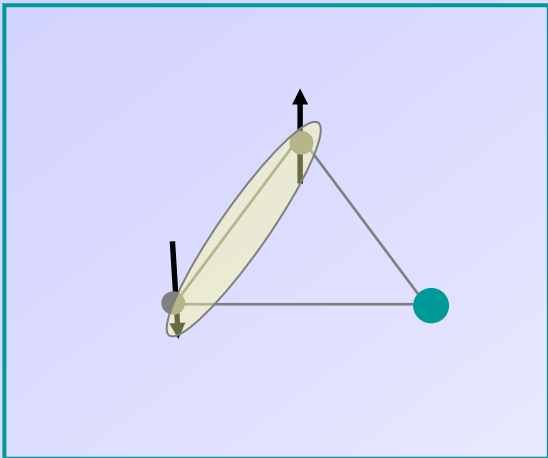
S. Dommange et al,
PRB 68 (2003) 224416

~20% oxygen sites
half shifted

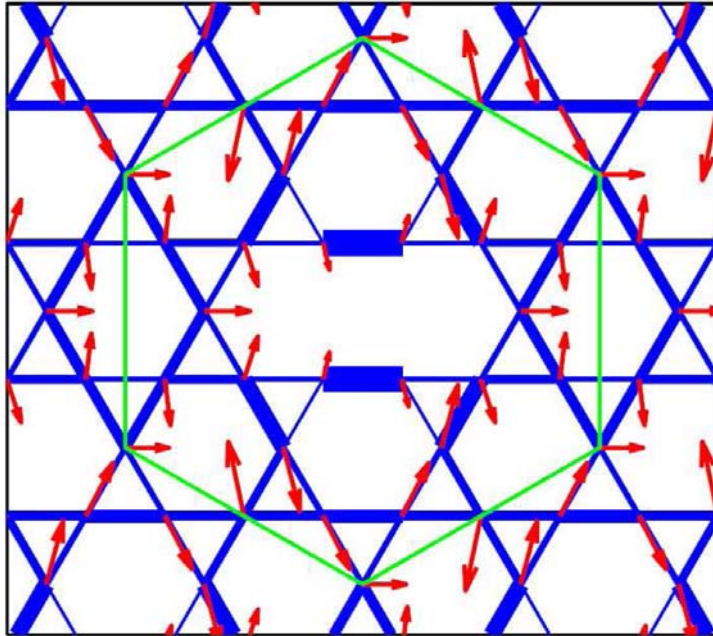
I. Rousochatzakis et al,
Phys.Rev.B 79, 214415 (2009)

-> ~5% spin vacancies in the kagome planes

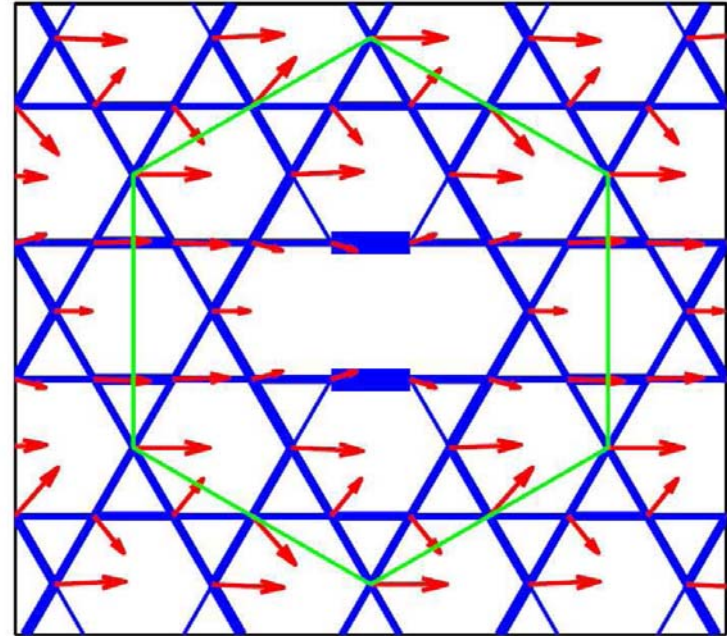
ED: one impurity



(a) $D/J=0.05$



(b) $D/J=0.1$



-Dimers next to the impurity survive up to $D \sim J$

-DM interaction removes the singlet nature of the ground state

Conclusions on Herbertsmithite

- first $S=1/2$ antiferromagnet with perfect kagome lattice (3 fold symmetry)
- no order or frozen disorder down to 50mK despite $J=175K$ and perturbations

Quantum spins on a perfect kagome lattice

- > allows close comparison with theory
- > renewed interest in understanding the GS

VBC : R.R.P. Singh and D.A. Huse, PRB (2007, 2008)

Dirac spin liquid : Y. Ran et al, PRL (2007), PRB (2008)

- magnetic defects (rather complex, in plane and out of plane) which impact the low T thermodynamic measurements
 - > local probe techniques
 - > **second sample generation : controlled (or no) defects**
- sizeable DM interaction : $D/J \sim 0.1$: **probe criticality: on-going**
- Local susceptibility determined from ^{17}O NMR:
 - ground state appears to be gap-less with a finite susceptibility
 - intrinsic to Heisenberg kagome model or **DM interaction closes the gap ?**
 - (no field effect)

PHFM 2010
Perspectives in Highly
Frustrated Magnetism

Dresden , 19 - 23 April

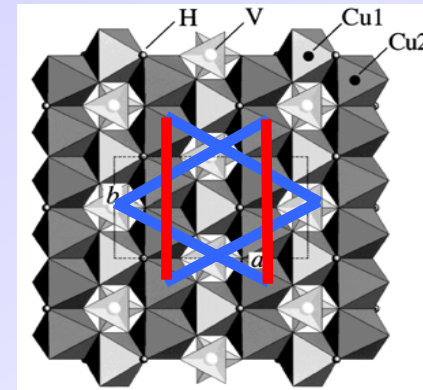
The logo of the European Science Foundation is displayed within a white rectangular box. It features the words "EUROPEAN", "SCIENCE", and "FOUNDATION" stacked vertically in a bold, sans-serif font. The first letter of each word (E, S, F) is colored green, while the remaining letters are grey. Horizontal lines are positioned below the words "EUROPEAN" and "SCIENCE".

EUROPEAN
SCIENCE
FOUNDATION

Frustration and original ground states - Collaborations

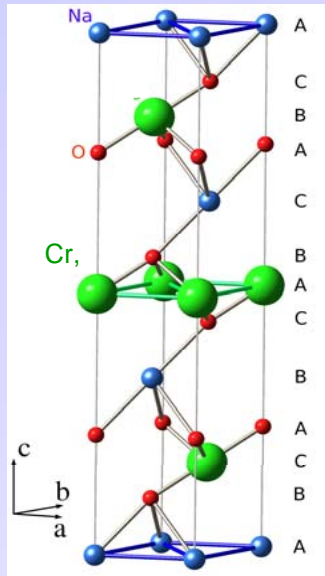
Herbertsmithite

F. Ladieux, S. Nakamae, P. Bonville
SPEC, CEA Saclay
MA de Vries, A. Harrison, Edinburgh
F. Duc, JC Trombe CEMES, Toulouse
P. Strobel, Institut Néel, Grenoble



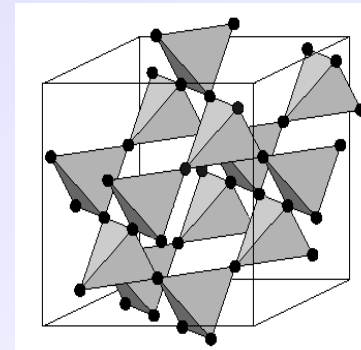
Kagome

Volborthite,
CEMES, Toulouse
Bert et al, PRL 2005



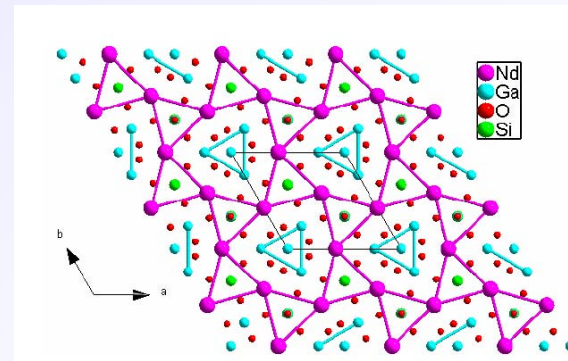
Triangular

NaCrO₂
R. Cava, Princeton
Olariu et al, PRL 2006



Pyrochlore

Tb₂Sn₂O₇
LLB, Saclay
Bert et al, PRL 2006



Kagome+

Spin anisotropy
Langasites
Institut Néel,
Grenoble
Zorko et al, PRL (2008)

Thank you!

Herbertsmithite:
 $\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$

Cu^{2+} , $S=1/2$

