

# ICTS Summer School on Numerical Relativity 2013 – Mathematical formulation

## Tutorial # 1

**1.1** Consider a flat, three-dimensional space in spherical polar coordinates,

$$ds^2 = g_{ab}dx^a dx^b = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2. \quad (1)$$

(a) Show that the only non-vanishing components of the connection are given by

$$\begin{aligned} \Gamma_{\theta\theta}^r &= -r & \Gamma_{\phi\phi}^r &= -r \sin^2 \theta \\ \Gamma_{\phi\phi}^\theta &= -\sin \theta \cos \theta & \Gamma_{r\theta}^\theta &= r^{-1} \\ \Gamma_{r\phi}^\phi &= r^{-1} & \Gamma_{\phi\theta}^\phi &= \cot \theta. \end{aligned} \quad (2)$$

and those that are related to the above by symmetry.

(b) Show that the Laplace operator acting on a scalar function  $f$ , defined as

$$\nabla^2 f \equiv \gamma^{ij} \nabla_i \nabla_j f, \quad (3)$$

reduces to the well-known expression

$$\nabla^2 f = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 f}{\partial \phi^2}. \quad (4)$$

(c) What is the Riemann tensor for this space?

**1.2** Consider a two-dimensional sphere of radius  $r$  with metric

$$ds^2 = g_{ab}dx^a dx^b = r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2. \quad (5)$$

Show that the Riemann tensor can be written as

$$R_{abcd} = \frac{1}{r^2} (g_{ac}g_{bd} - g_{ad}g_{bc}). \quad (6)$$

*Hint:* Confirm that the above expression has the correct symmetries, and think about the number of independent components of the Riemann tensor.

**1.3** In class we discussed the Schwarzschild metric

$$ds^2 = -f_0 dt^2 + f_0^{-1} dR^2 + R^2 d\Omega^2, \quad (7)$$

where  $f_0 = 1 - 2M/R$  and where  $R$  is the *areal* (or *circumferential*) radius. Consider a new radius  $r$  that satisfies

$$R = r \left( 1 + \frac{M}{2r} \right)^2 \quad (8)$$

(a) Show that under this transformation the Schwarzschild metric (7) takes the form

$$ds^2 = - \left( \frac{1 - M/(2r)}{1 + M/(2r)} \right)^2 dt^2 + \psi^4 (dr^2 + r^2 d\Omega^2) \quad (9)$$

and find the “conformal factor”  $\psi$ . This form of the metric is called *isotropic*.

(a) Find an expression for  $r$  in terms of  $R$ . Show that real solutions for  $r$  only exist  $R \geq 2M$ .

- 1.4 In the previous problem you considered a radial transformation of the Schwarzschild metric. Now consider a temporal transformation to a new time coordinate

$$t' = t + h(R), \tag{10}$$

where the *height function*  $h(R)$  describes by how much surfaces of constant new time  $t'$  “lift off” surfaces of constant Schwarzschild time  $t$ .

- (a) Show that in these new coordinates the Schwarzschild metric (7) takes the form

$$ds^2 = -f_0 dt'^2 + 2f_0 h' dt' dR + f_0^{-1} (1 + f_0^2 h'^2) dR^2 + R^2 d\Omega^2, \tag{11}$$

where  $h' \equiv dh/dR$ .

- (b) Find a suitable condition for  $h$  so that the Schwarzschild metric takes the *Painlevé-Gullstrand* form

$$ds^2 = -\left(1 - \frac{2M}{R}\right) dt'^2 + 2\sqrt{\frac{2M}{R}} dt' dR + dR^2 + R^2 d\Omega^2. \tag{12}$$

Painlevé-Gullstrand coordinates are remarkable in that they leave the spatial part of the metric flat.